

[arXiv:2303.03594](https://arxiv.org/abs/2303.03594) submitted on March 7

First Results of Axion Dark Matter Search with DANCE

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Overview

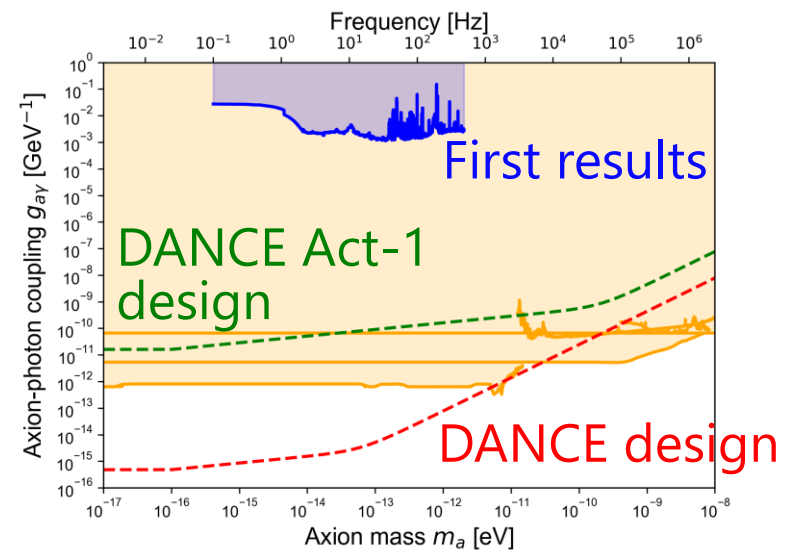
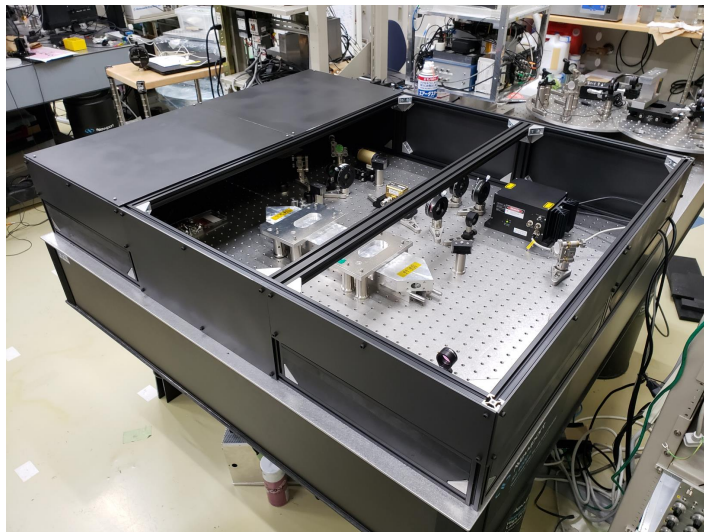
- New experimental project to search for axion DM with an optical cavity

I. Obata, T. Fujita, Y. Michimura
[PRL 121, 161301 \(2018\)](#)

DANCE: Dark matter **A**xion search
with ri**N**g **C**avity **E**xperiment

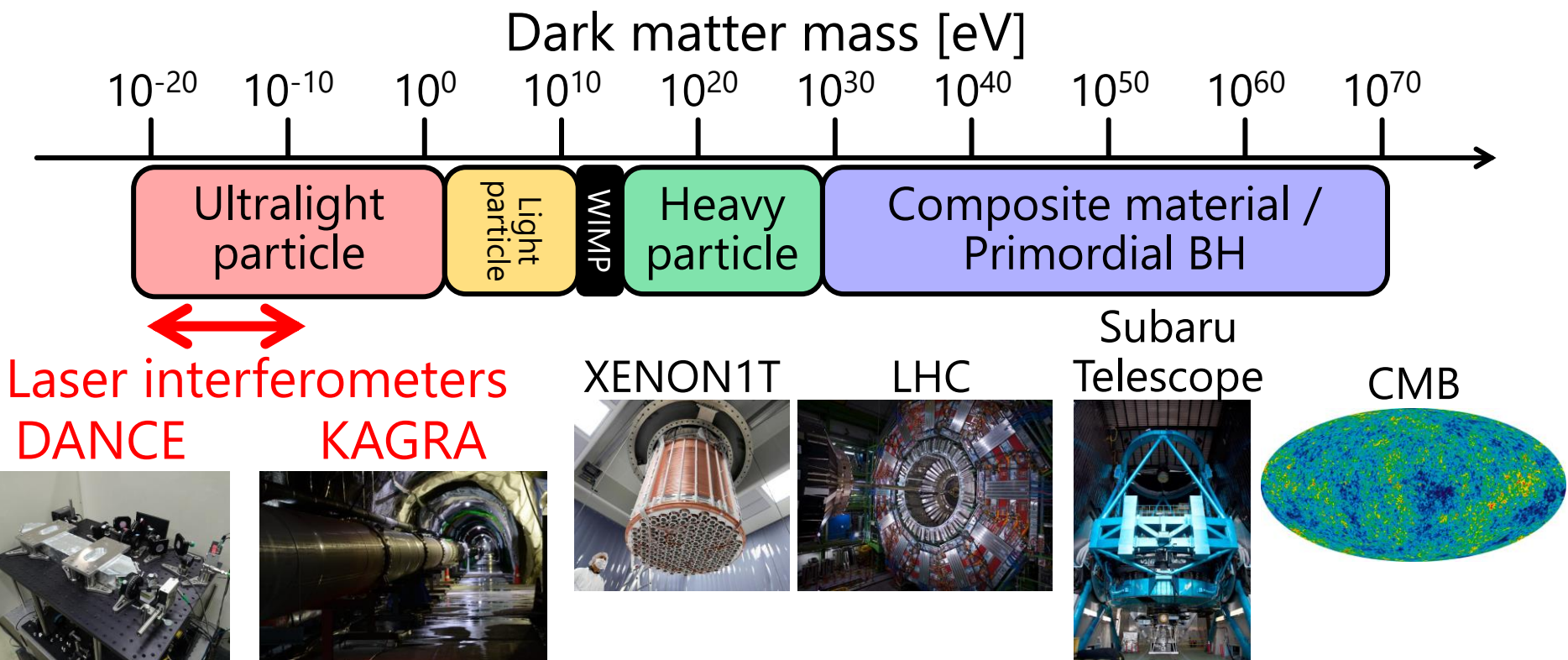
- First results of prototype experiment **DANCE Act-1** from long-term observation

YO, H. Fujimoto+, [arXiv:2303.03594](#)



Axion search with laser interferometers

- We need to search for DM in a wider mass range
- Laser interferometers are useful to search for ultralight DM
- DANCE focuses on axion DM



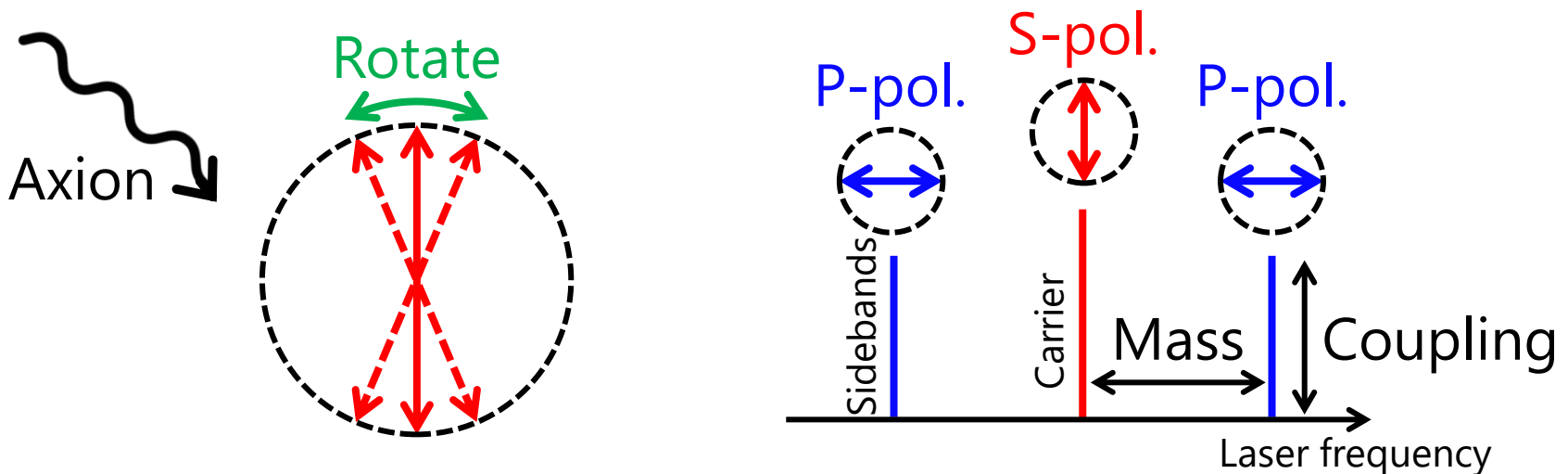
Polarization rotation from axions

- Axion-photon coupling causes phase velocity difference between left- and right-handed photons

$$c_{L/R} = \sqrt{1 \pm \frac{g_{a\gamma} a_0 m_a}{k} \sin(m_a t + \delta_\tau)}$$

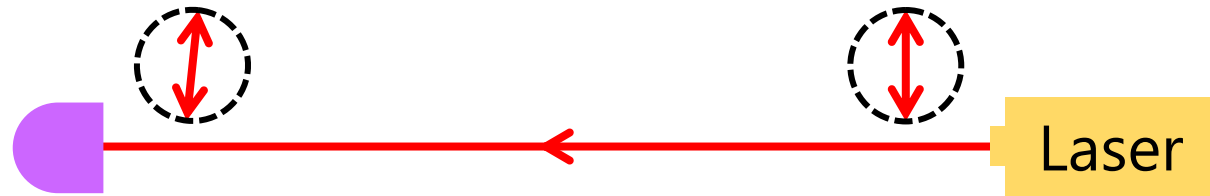
Coupling constant Axion field Axion mass

- Phase velocity difference of circular polarizations makes linear polarization rotate and oscillate

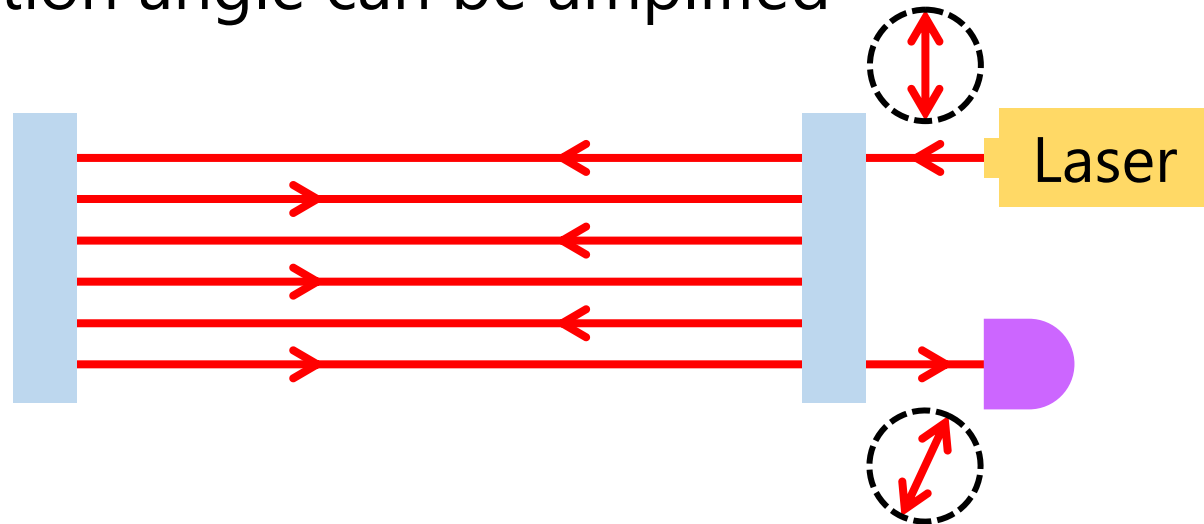


Signal amplification with cavities

- Rotation angle is too small to be observed without a cavity

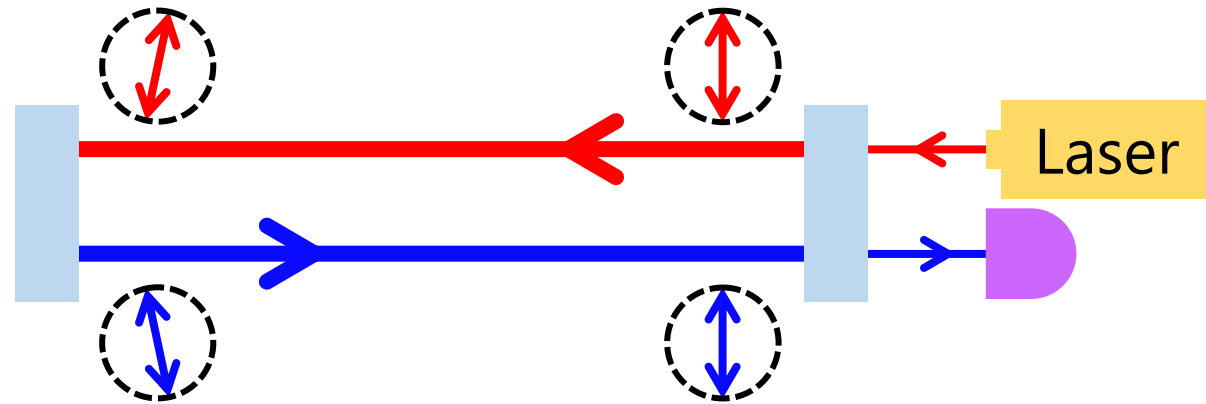


- Laser light runs many times between mirrors in an optical cavity
→ Rotation angle can be amplified

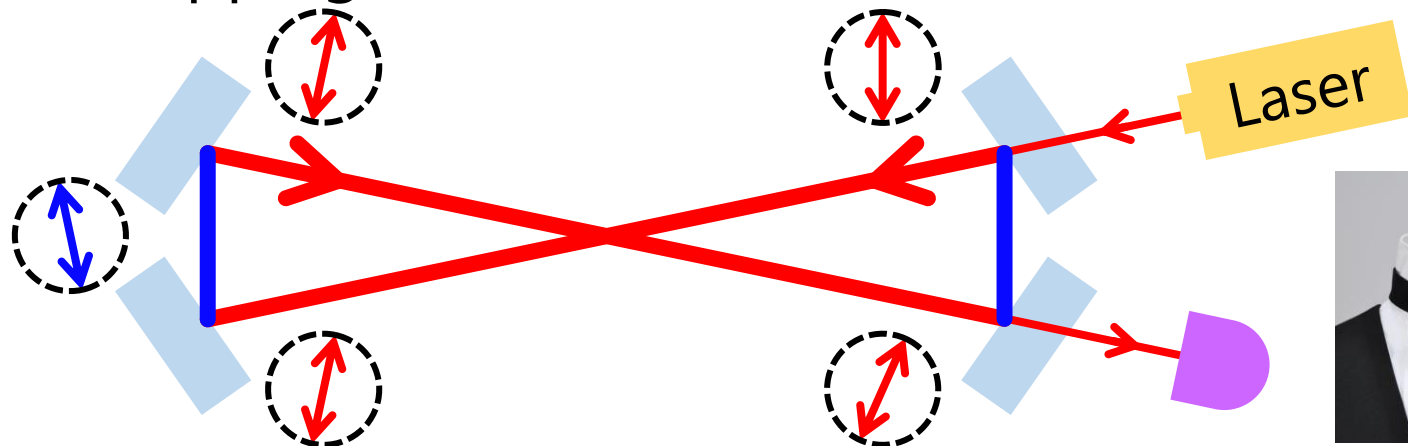


Bow-tie ring cavity

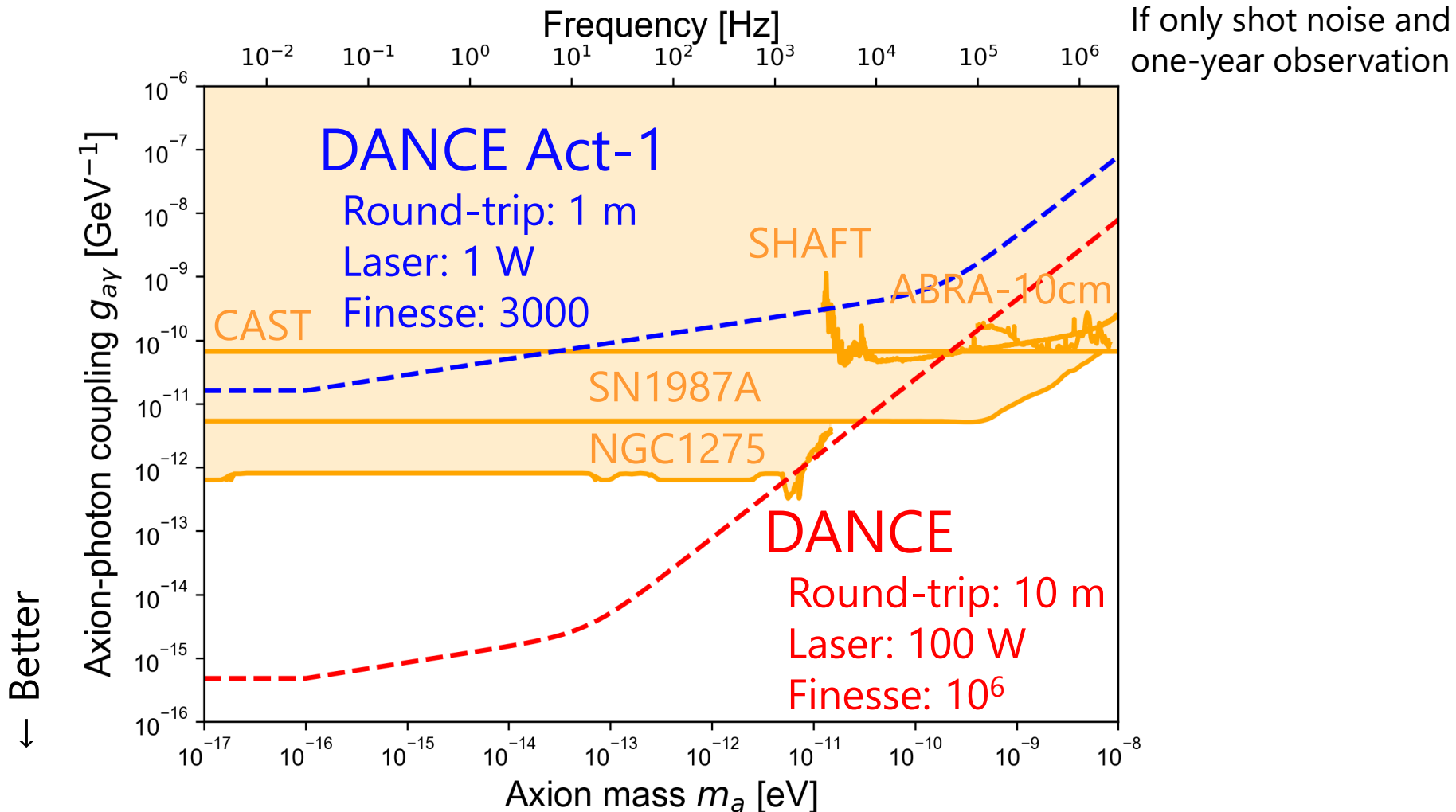
- Rotated direction is inverted by reflection on mirrors
→ Rotation effect is canceled out



- A bow-tie ring cavity prevents linear polarization from flipping

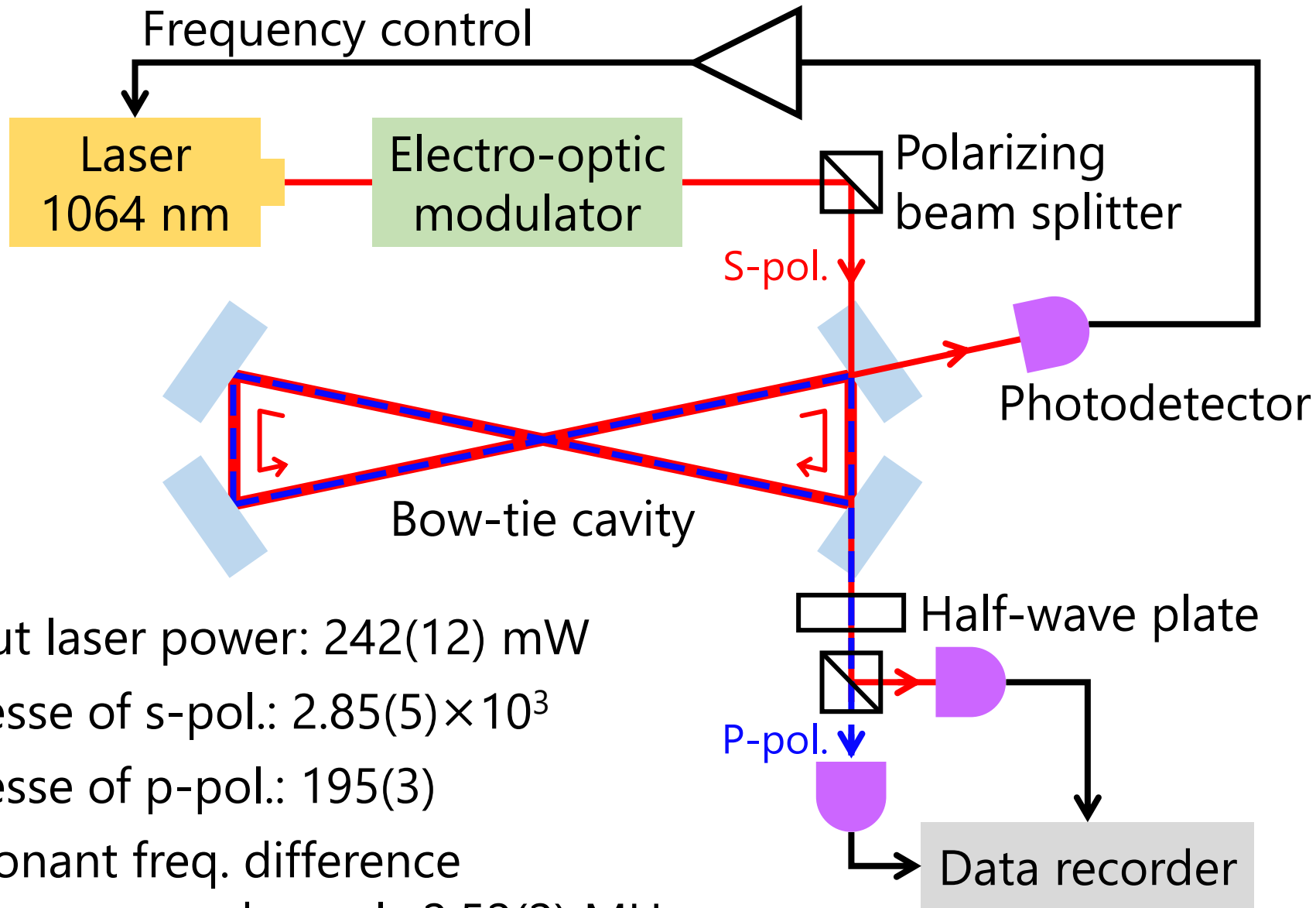


Design sensitivity of DANCE



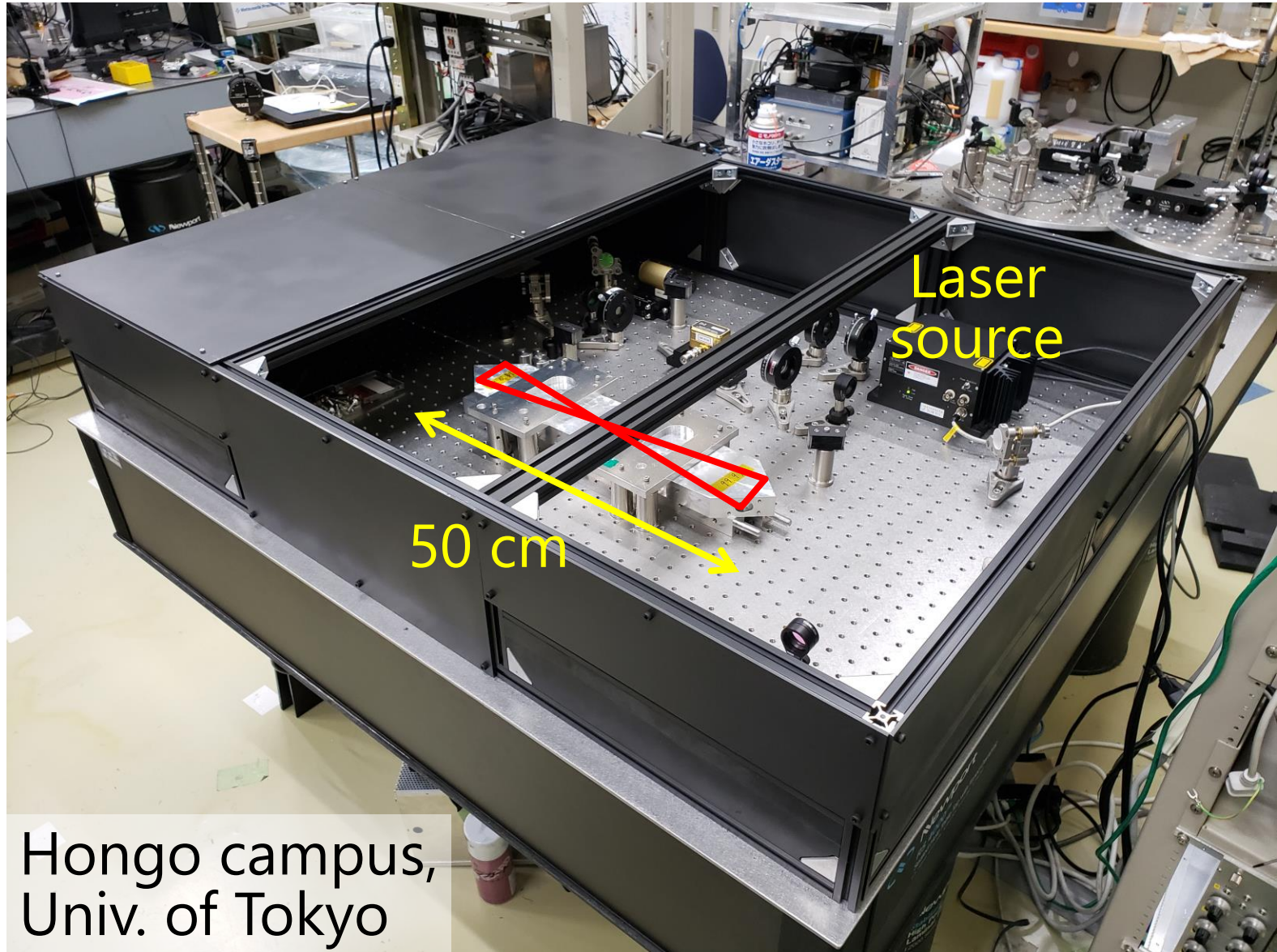
- Shot noise is caused by fluctuations in photons' number
- Need to minimize the other noise sources

Setup of DANCE Act-1



- Input laser power: 242(12) mW
- Finesse of s-pol.: $2.85(5) \times 10^3$
- Finesse of p-pol.: 195(3)
- Resonant freq. difference between s- and p-pol.: 2.52(2) MHz

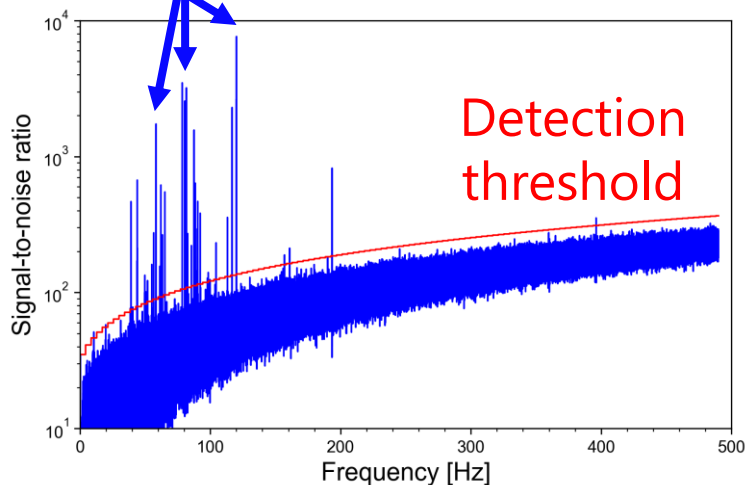
Picture of DANCE Act-1



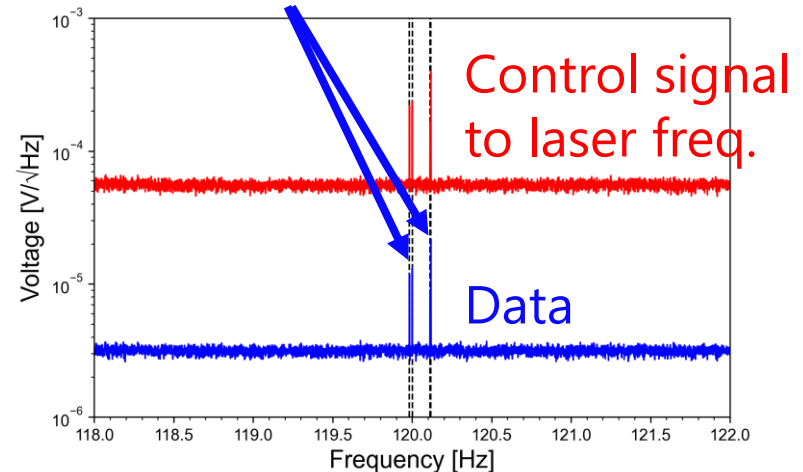
Data acquisition and analysis

- Recorded the data in May 18-30, 2021
- The first 86,400-second (24-hour) data was selected
- After passing the data through the analysis pipeline, 551 points exceeded the threshold
- All candidate peaks were rejected by 3 veto procedures
 - Consistency veto: 551 \rightarrow 271
 - Linewidth veto: 271 \rightarrow 7
 - Control signal veto: 7 \rightarrow 0

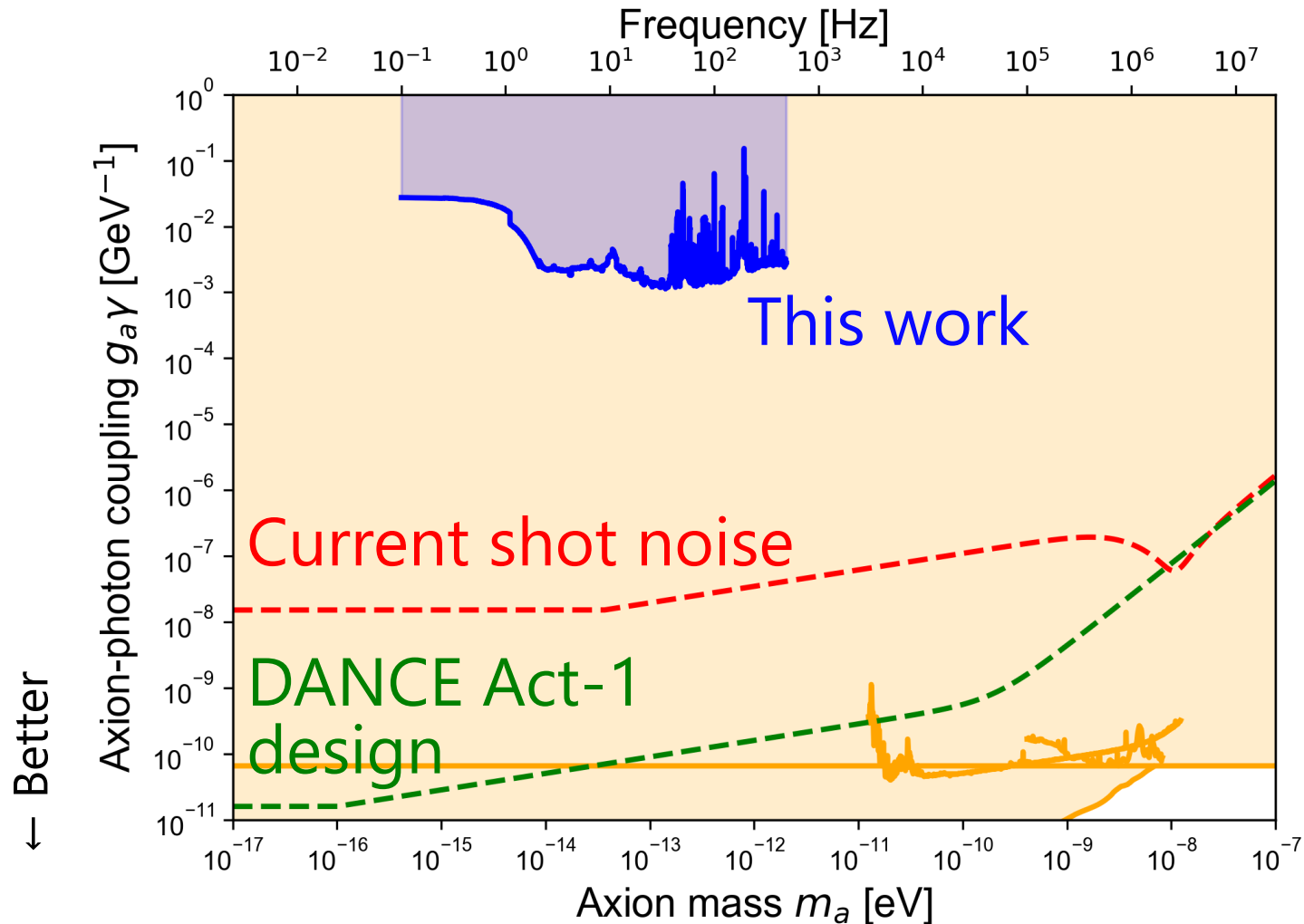
Candidate peaks



Probably due to mechanical resonance

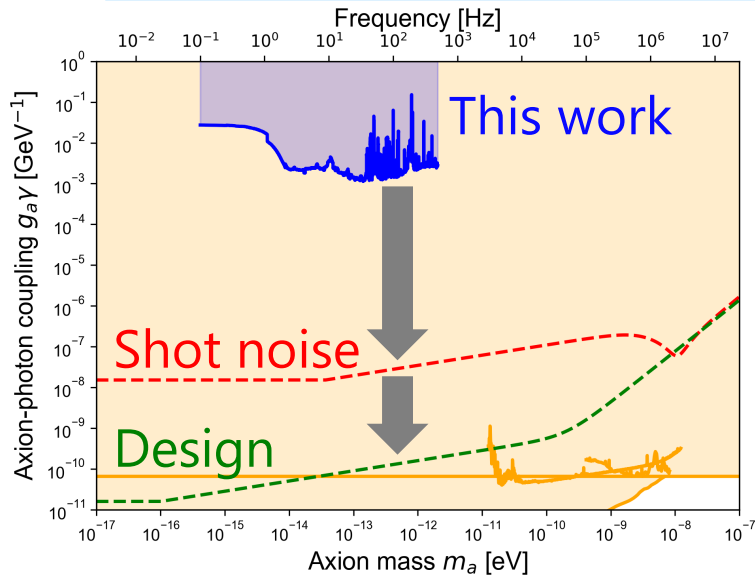


Results



- Worse than design sensitivity by 7 orders of magnitude
- First results of axion DM search with an optical cavity

Discussion to improve the sensitivity



1. This work to shot noise

- We need to reduce classical noises
 - Laser intensity noise
 - Laser frequency noise
 - Mechanical vibration

2. Shot noise to design sensitivity

- We need to improve the parameters
 - Input laser power: 0.2 W \rightarrow 1 W
 - Observation time: 24 hours \rightarrow 1 year
 - Resonant freq. difference between s- and p-pol.: 3 MHz \rightarrow 0 Hz (simultaneous resonance)

We are installing an auxiliary cavity (Hiroki's talk)

Summary

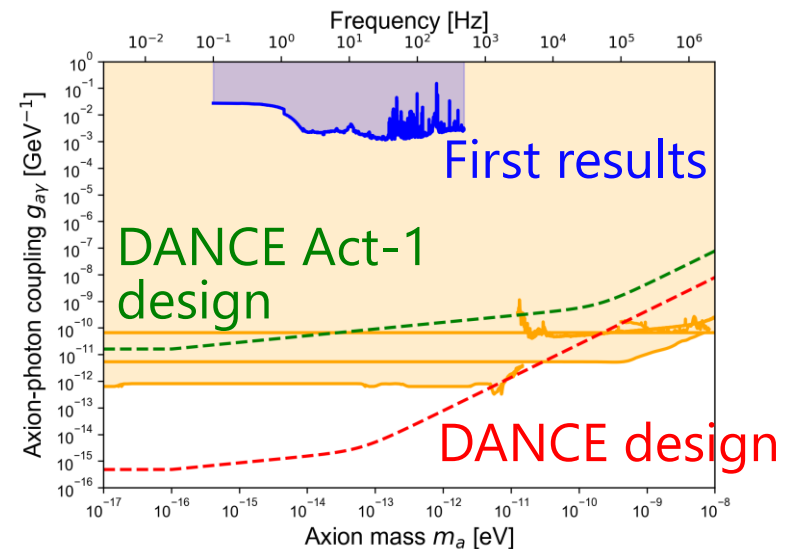
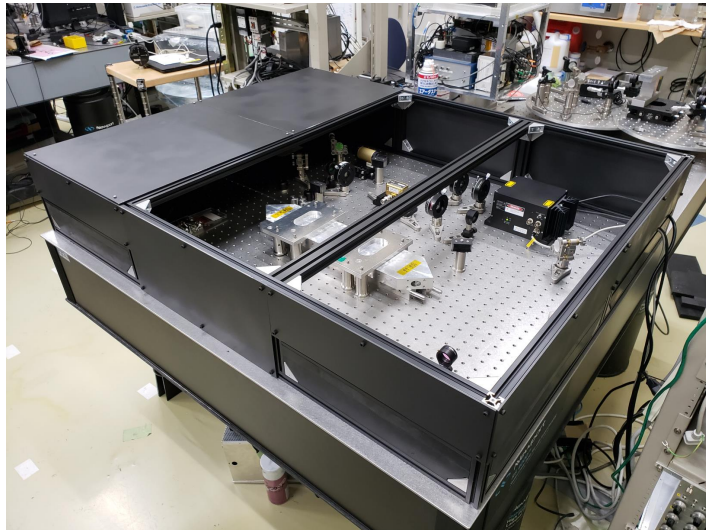
- New experimental project to search for axion DM with a bow-tie cavity: **DANCE**

I. Obata, T. Fujita, Y. Michimura
[PRL 121, 161301 \(2018\)](#)

- Prototype experiment **DANCE Act-1** is ongoing
 - Long-term observation in May 2021
 - **The first upper bounds on $g_{a\gamma}$ with an optical cavity**
 - We continue to improve the sensitivity

YO, H. Fujimoto+, [arXiv:2303.03594](#)

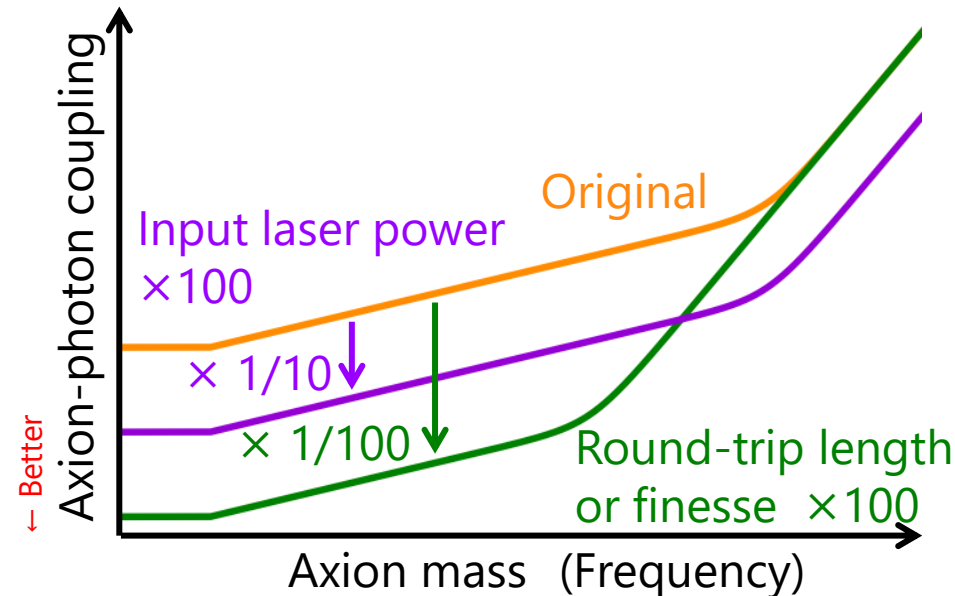
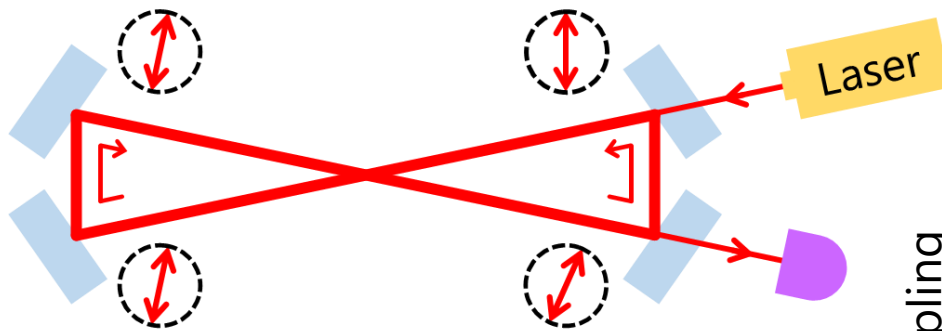
H. Fujimoto, YO+, [JPCS 2156, 012182 \(2021\)](#)



Extra Slides

Important parameters (1)

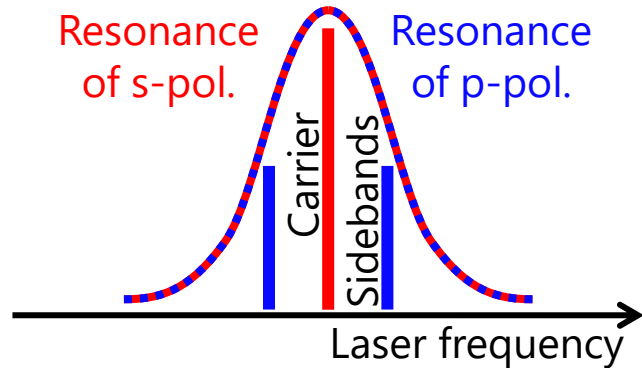
- Input laser power
 - ... Shot noise
- Round-trip length
 - ... Optical length
 - ... Number of round trip
- Finesse



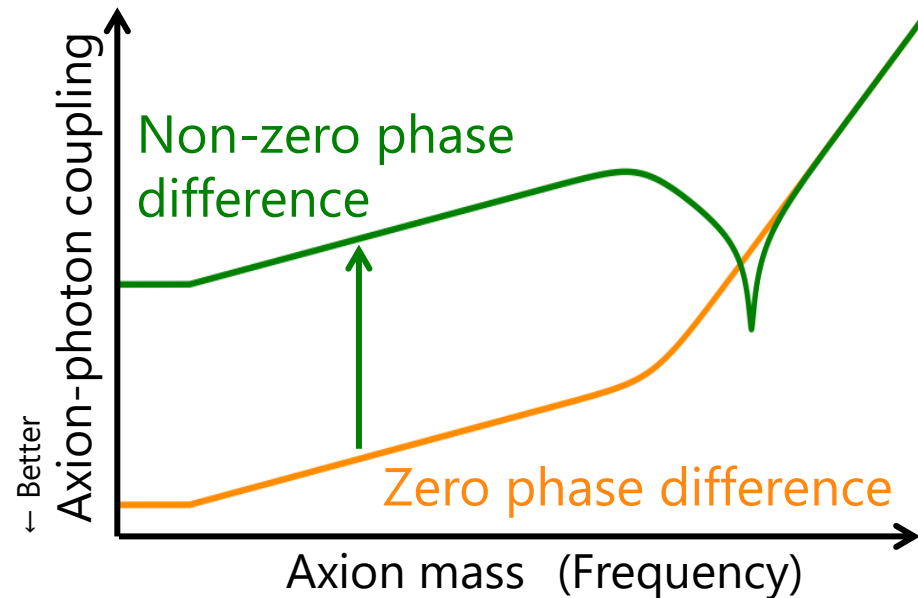
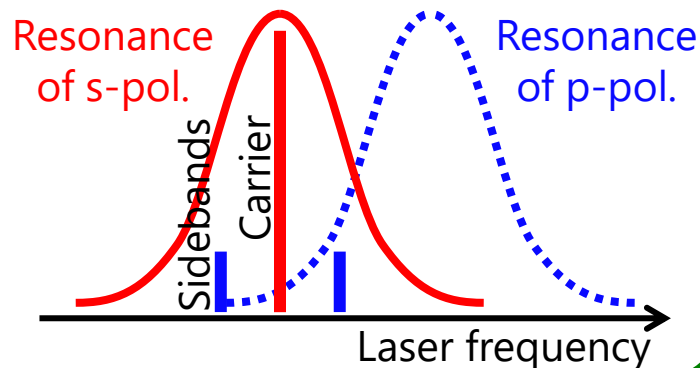
Important parameters (2)

- Reflective phase difference between s- and p-pol.
 - From different phase shifts by reflections on mirrors

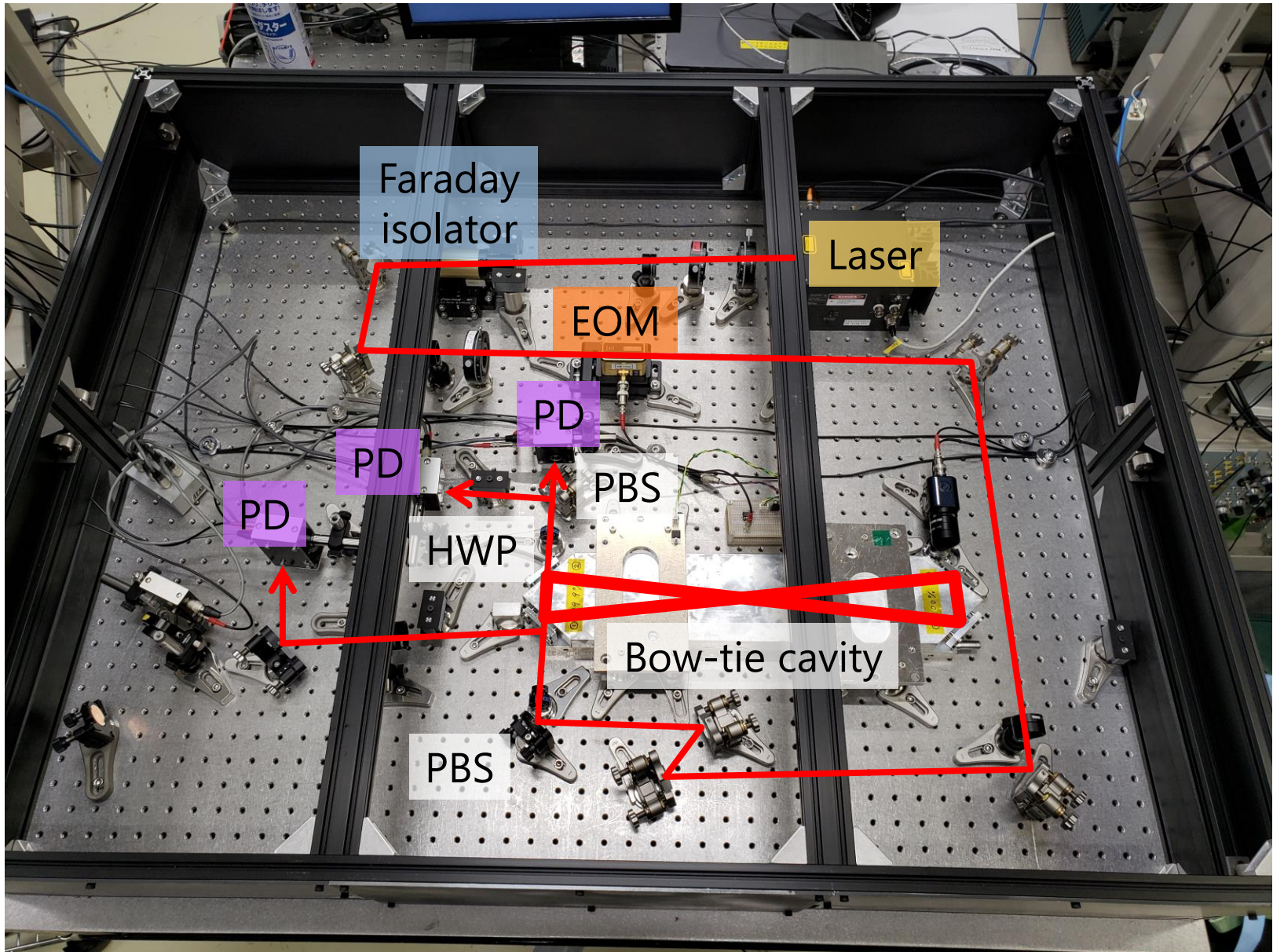
Zero phase difference



Non-zero phase difference

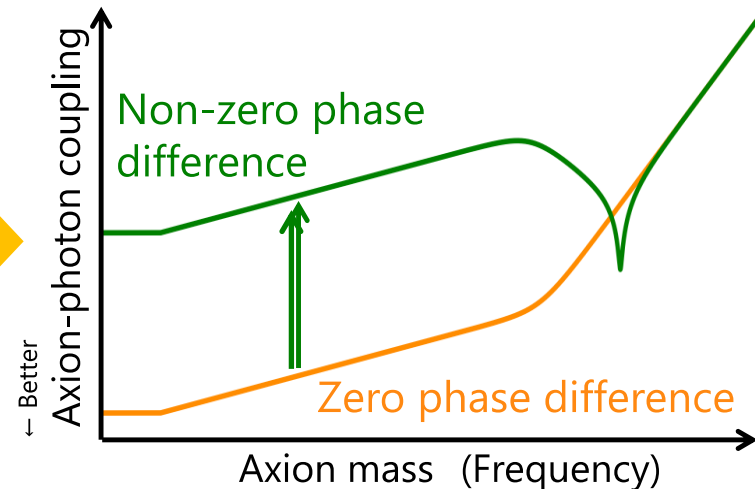
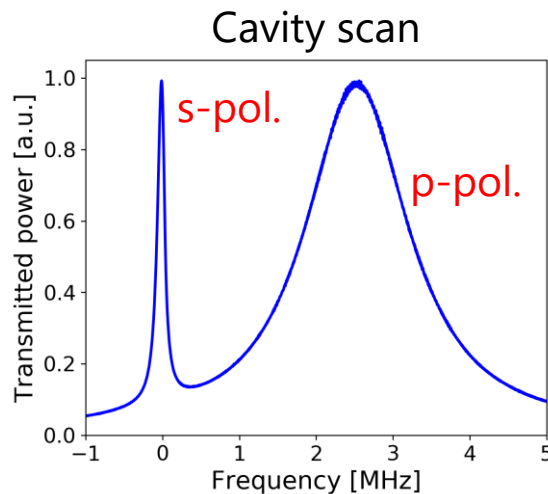


Picture of experimental setups

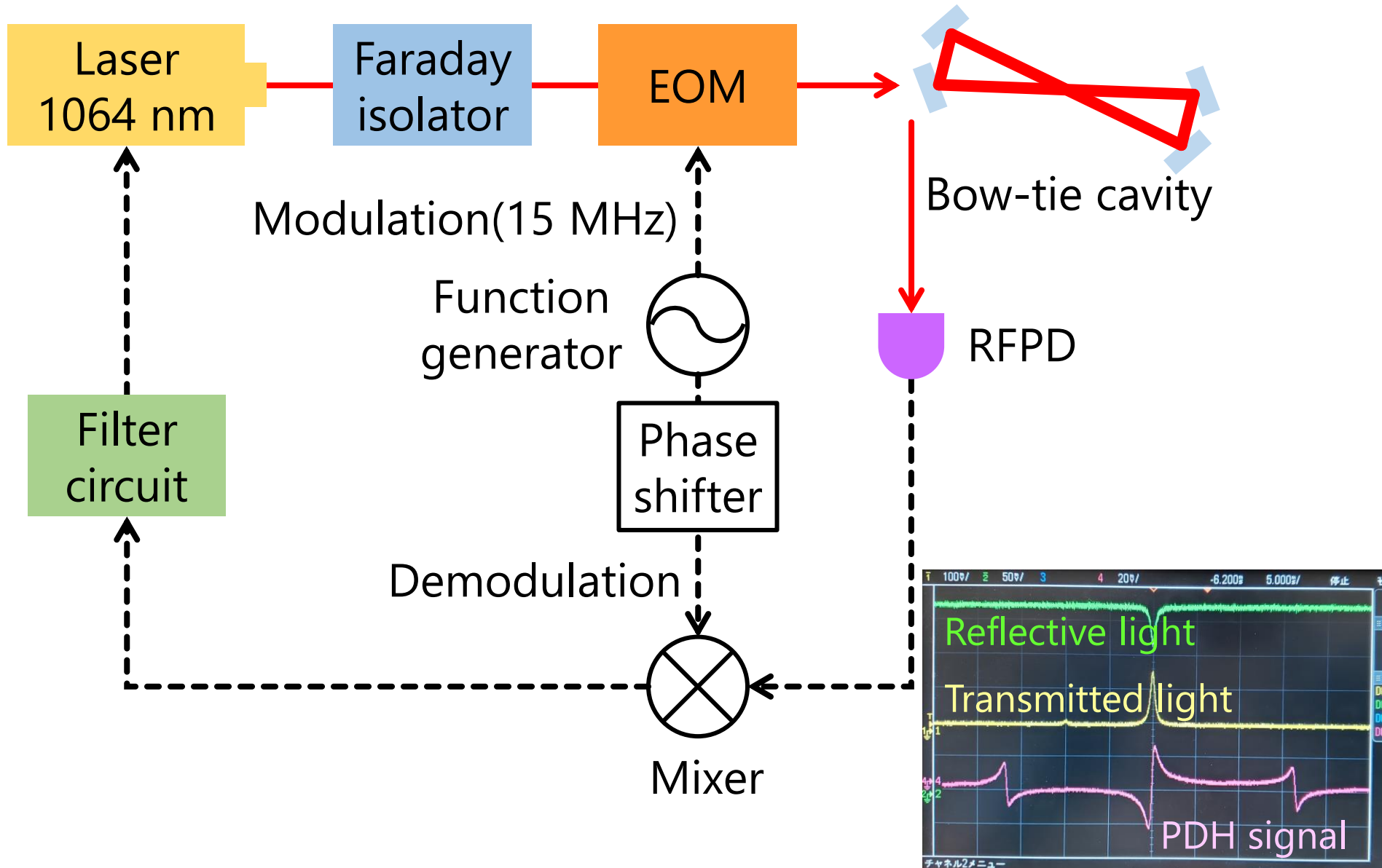


Performance of the cavity

	Designed values	Measured values
Input laser power	1 W	242(12) mW
Transmitted laser power	1 W	153(8) mW
Finesse of s-pol.	3×10^3	$2.85(5) \times 10^3$
Finesse for p-pol.	3×10^3	195(3)
Reflective phase difference between s- and p-pol.	<1 mrad	52.4(4) mrad

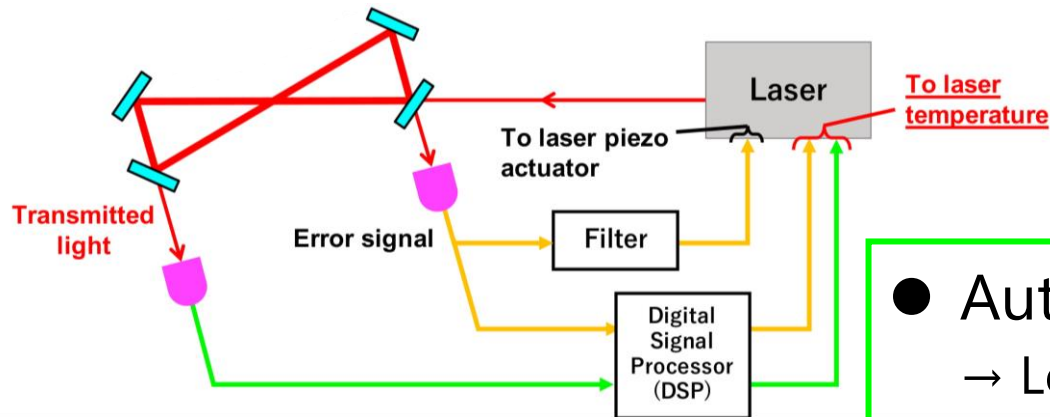


Frequency servo by PDH technique



Long-term frequency control

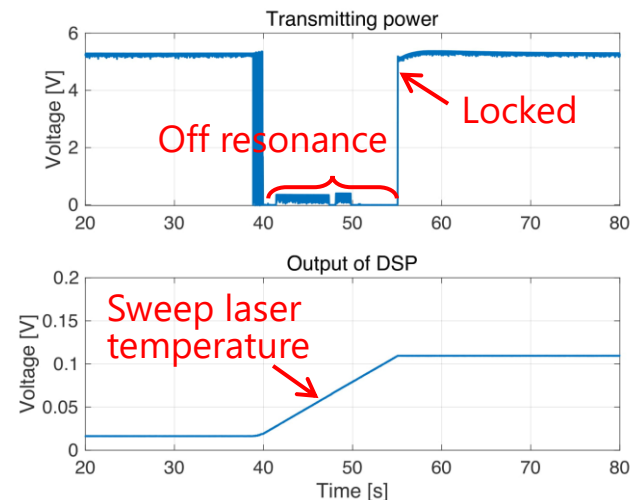
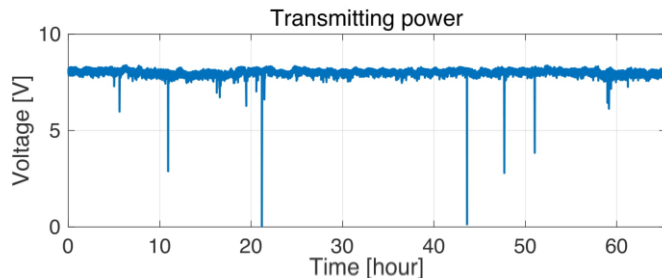
- SNR can be improved with observation time
 - Need to observe for a long time (e.g. First test operation: 12 days, Final goal: 1 year)



Experiment by H. Fujimoto
[arXiv:2105.08347 \(2021\)](https://arxiv.org/abs/2105.08347)

- Automated cavity locking system
 - Locked automatically by identifying whether on or off resonance

- Double-loop control
 - Kept resonance longer than 60 hours

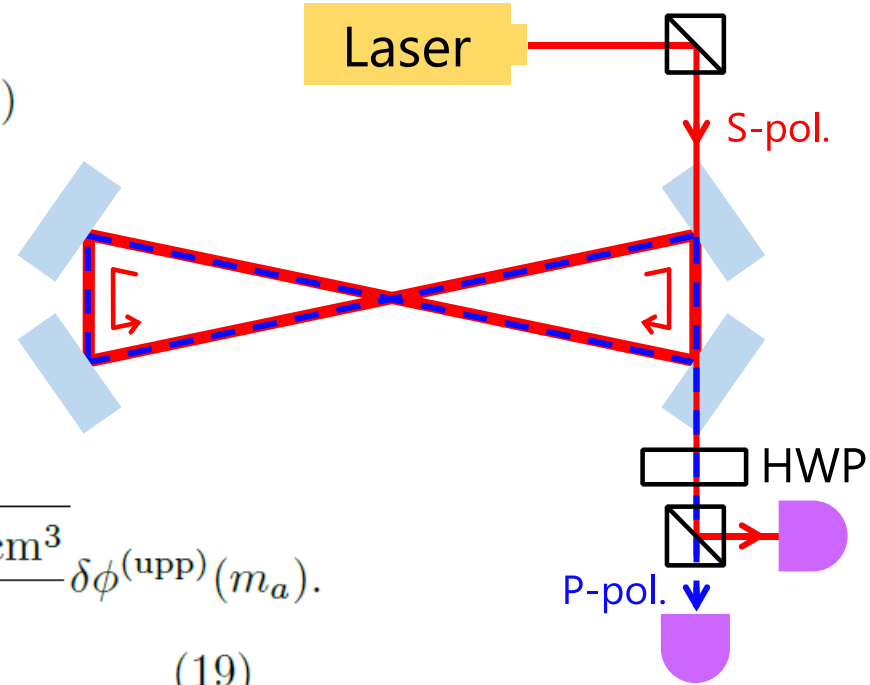


Signal calibration

$$\phi(t) = \sqrt{\frac{P_p(t)}{P_s(t) + P_p(t)}} - 2\theta_{\text{HWP}}, \quad (16)$$

$$g_{a\gamma}(m_a) = \frac{2k_0 \delta\phi^{(\text{uPP})}(m_a)}{\sqrt{2\rho_a} |H'_a(m_a)|}$$

$$= 5.10 \times 10^{11} \text{ GeV}^{-1} \frac{\text{eV}^{-1}}{|H'_a(m_a)|/k_0} \sqrt{\frac{\text{GeV}/\text{cm}^3}{\rho_a}} \delta\phi^{(\text{uPP})}(m_a). \quad (19)$$

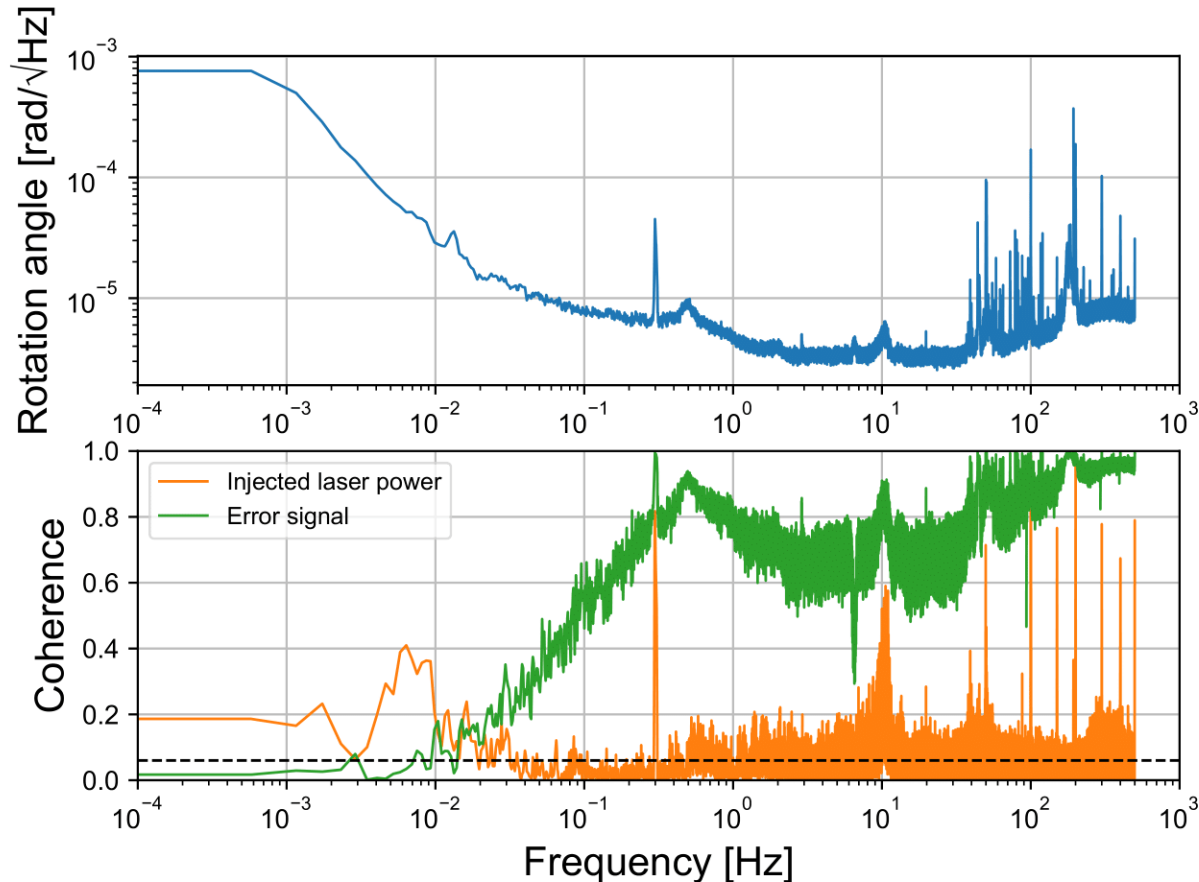


$$H_a(\omega) \equiv k_0 \sqrt{\frac{1 - r_{1p}^2}{1 - r_{1s}^2}} \frac{1}{i\omega(1 - r_{1p}^2 r_{2p}^2) e^{-i\omega(2l_1 + 2l_2)}}$$

$$\times \left[(1 - e^{-i\omega l_2}) \left(r_{1s} r_{2s} r_{1p} r_{2p} e^{-i\omega l_1} + r_{1p}^2 r_{2p}^2 e^{-i\omega(2l_1 + l_2)} \right) \right.$$

$$\left. - (1 - e^{-i\omega l_1}) \left(r_{1s} r_{2s}^2 r_{1p} + r_{1s} r_{1p} r_{2p}^2 e^{-i\omega(l_1 + l_2)} \right) \right]. \quad (13)$$

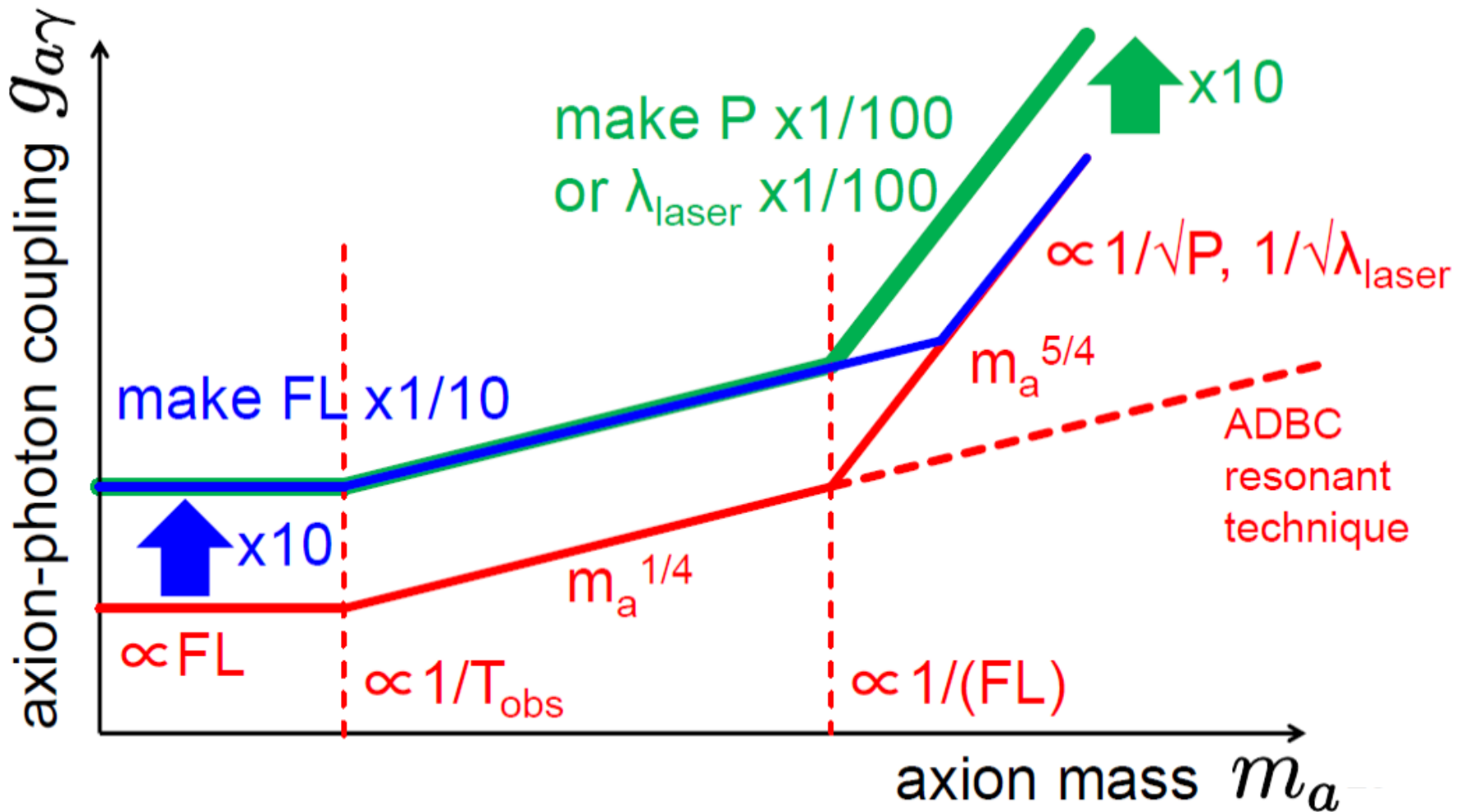
Discussion for noise sources



- -0.01 Hz: Suggested to be limited by laser intensity noise
- 0.01 Hz-: Suggested to be limited by mechanical vibration and laser freq. noise

Sensitivity Design

- Brute force necessary, you cannot win for free



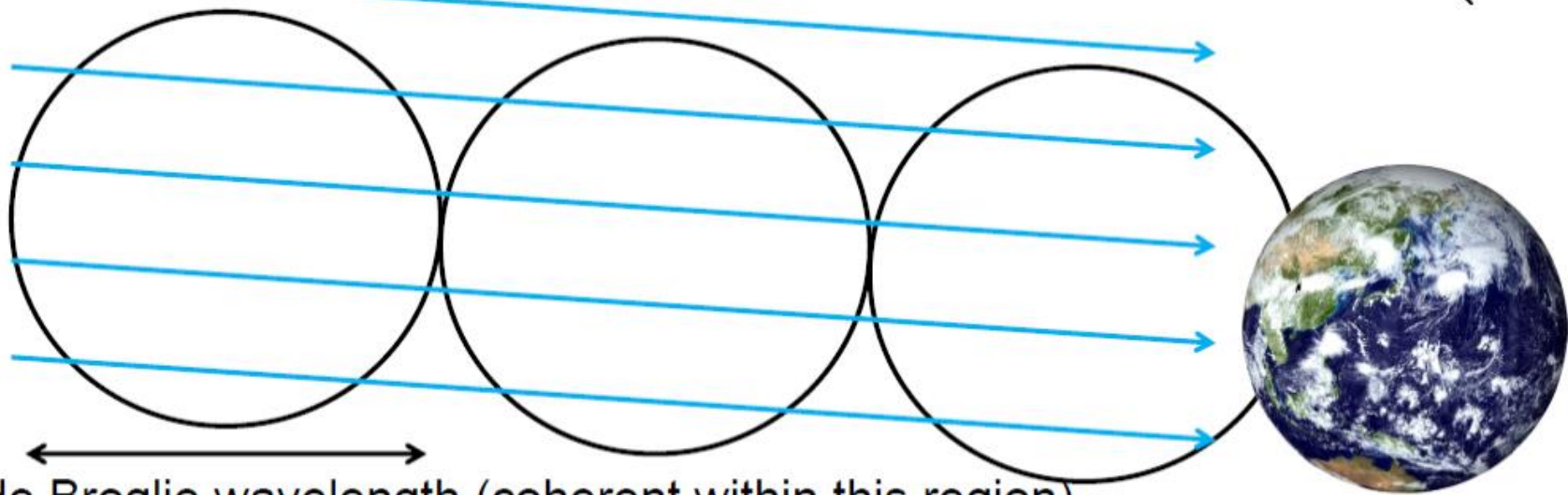
Coherent Time Scale

- SNR grows with $\sqrt{T_{\text{obs}}}$ if integration time is shorter than coherent time scale
- SNR grows with $(T_{\text{obs}})^{1/4}$ if integration time is longer

$$\text{SNR} = \begin{cases} \frac{\sqrt{T_{\text{obs}}}}{2\sqrt{S_{\text{noise}}(f)}} \frac{\delta c}{c} & (T_{\text{obs}} \lesssim \tau) \\ \frac{(T_{\text{obs}}\tau)^{1/4}}{2\sqrt{S_{\text{noise}}(f)}} \frac{\delta c}{c} & (T_{\text{obs}} \gtrsim \tau) \end{cases}$$

$$\tau \simeq 1 \text{ year} \left(\frac{10^{-16} \text{ eV}}{m_a} \right)$$

axion wind



de Broglie wavelength (coherent within this region)