Search for ultralight dark matter with laser interferometric gravitational wave detectors

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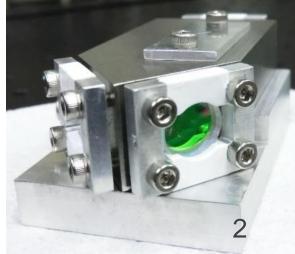
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Slides are available at https://tinyurl.com/YM20200713

Self Introduction

- Yuta Michimura (道村唯太)
 Department of Physics, University of Tokyo
- Laser interferometric gravitational wave detectors
 - KAGRA
 - DECIGO
- Search for new physics with laser interferometry
 - Lorentz violation
 - Macroscopic quantum mechanics
 - Dark matter searches etc ...





Plan of the Talk

Basics of laser interferometry

- Michelson interferometer and optical cavity
- Laser interferometric gravitational wave detectors
- Key aspects

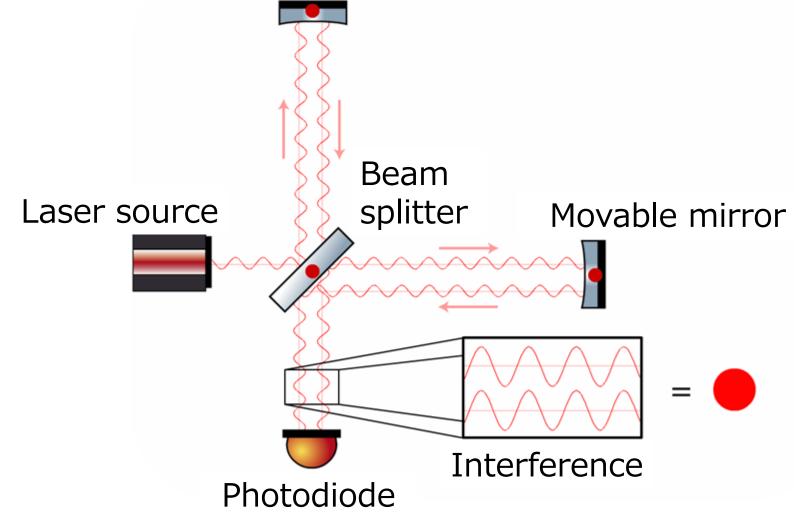
Ultralight dark matter searches

- Axion like particles (pseudoscalar)
- Scalar fields
- $U(1)_{B}$ and $U(1)_{B\text{-L}}$ gauge bosons (vector)
- Dark matter search with KAGRA
 - Current status of KAGRA
 - Prospected sensitivity for KAGRA
- Summary

Basics of laser interferometry

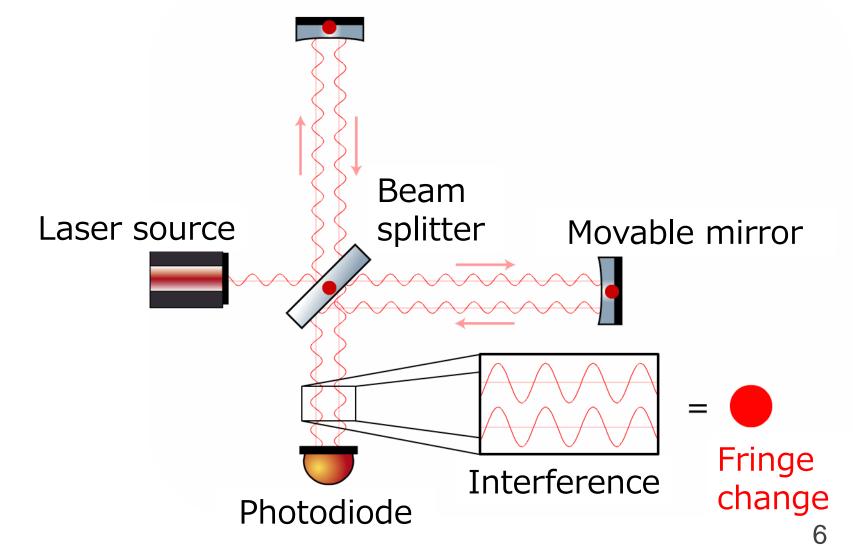
Michelson Interferometer

• measures differential arm length change



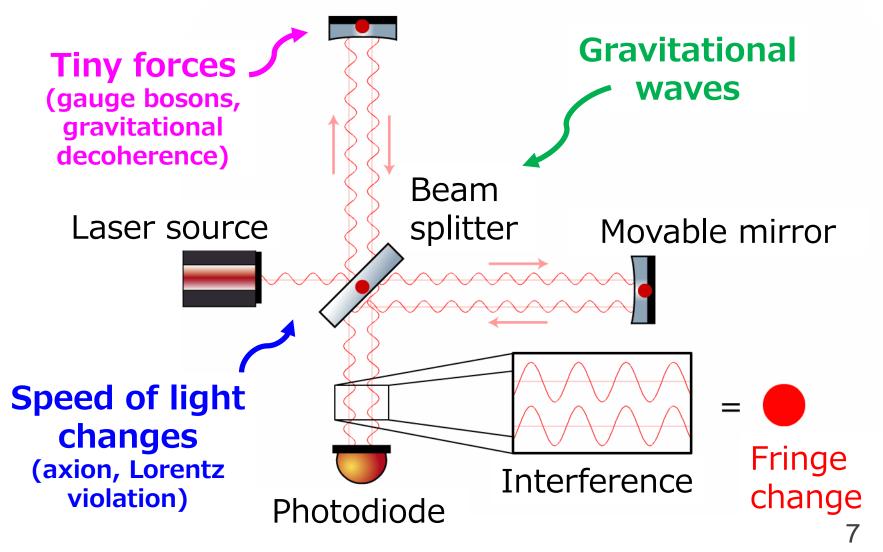
Michelson Interferometer

• measures differential arm length change



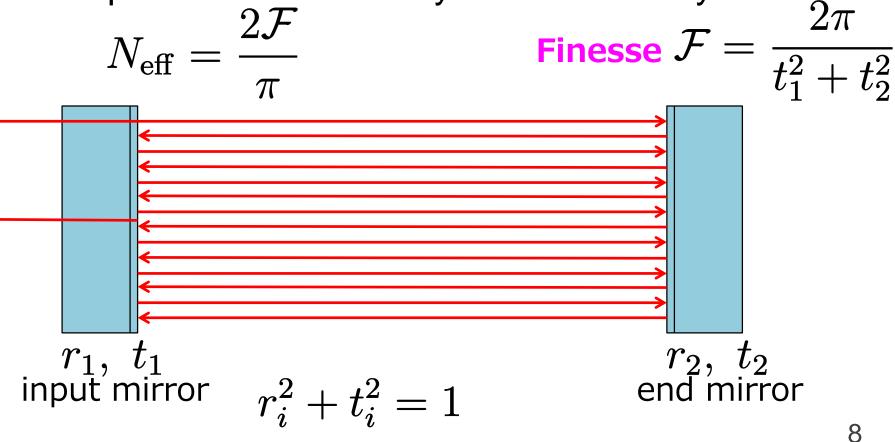
Michelson Interferometer

• measures differential arm length change

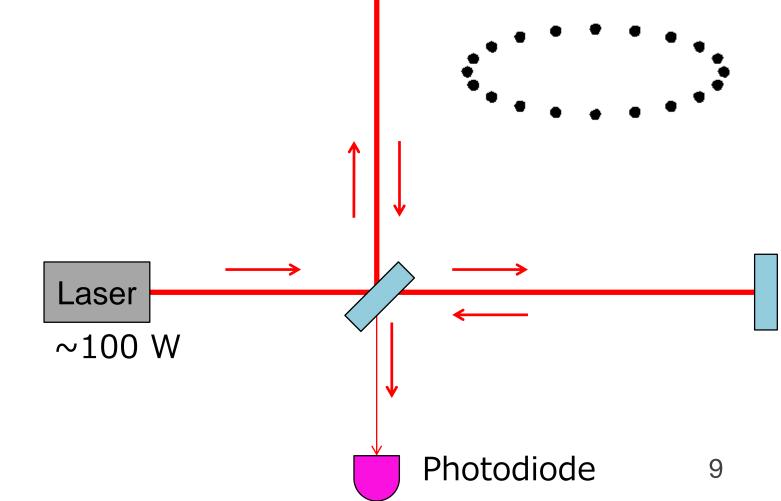


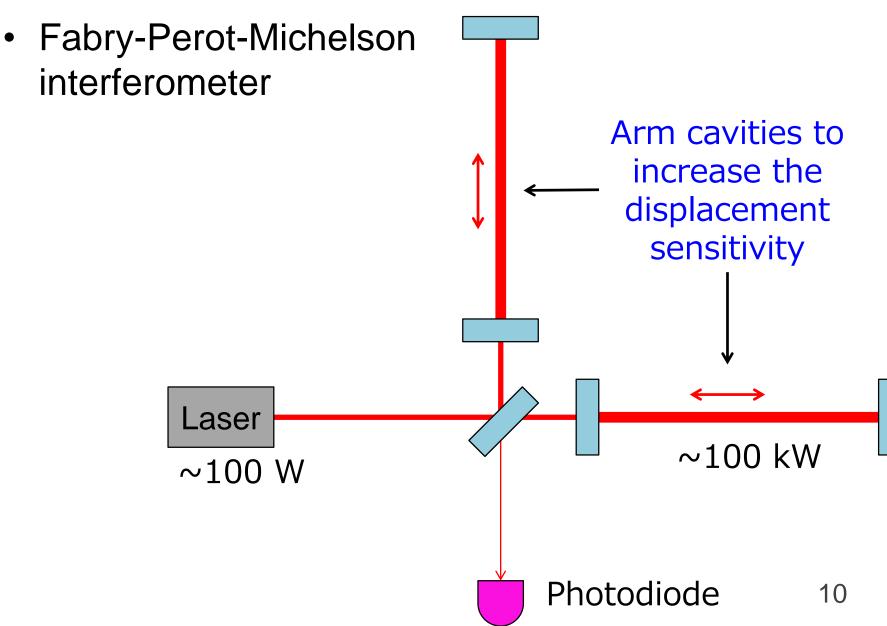
Fabry-Perot Cavity

- Two highly reflective mirrors
- Sense mirror displacement multiple times
- Displacement sensitivity is enhanced by

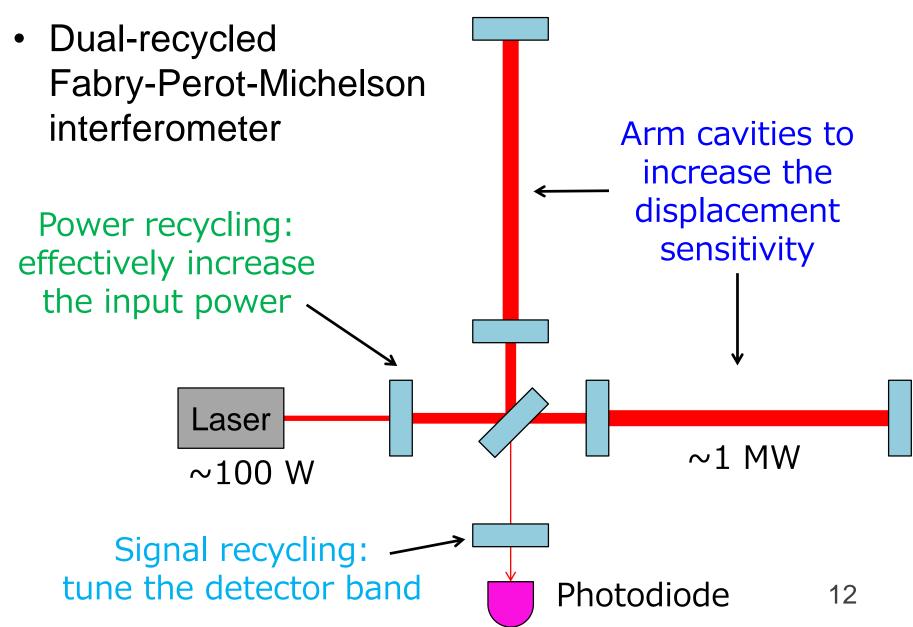


Michelson interferometer





 Power-recycled **Fabry-Perot-Michelson** interferometer Arm cavities to increase the displacement Power recycling: sensitivity effectively increase the input power Laser $\sim 1 \text{ MW}$ $\sim 100 \text{ W}$ Photodiode 11

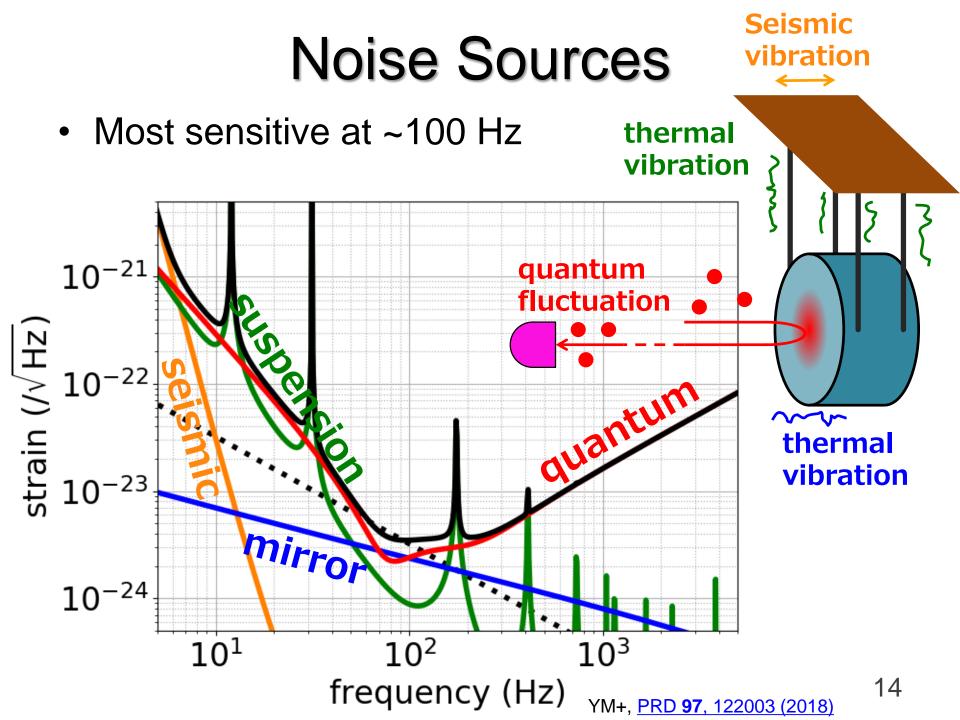


Global Network of GW Detectors

All are laser interferometric GW detectors

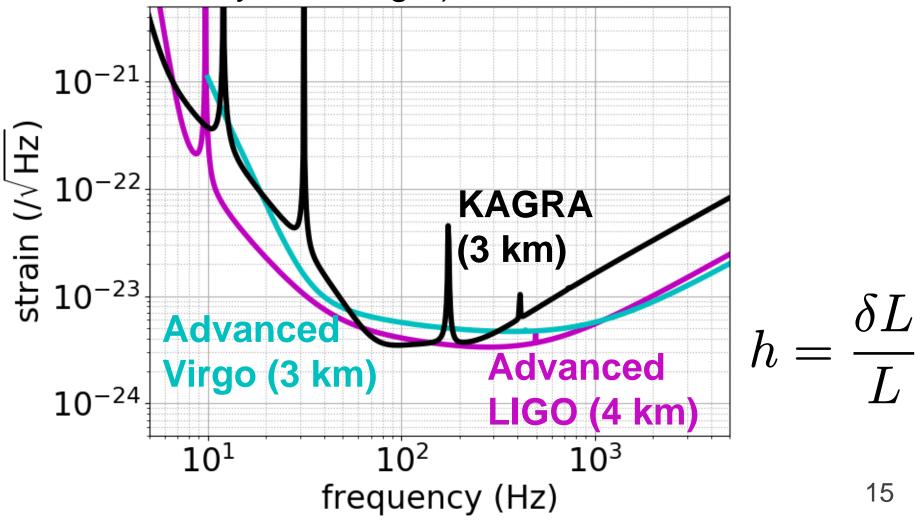
Advanced LIGO





Sensitivity of LIGO/Virgo/KAGRA

• Similar strain sensitivity (displacement sensitivity divided by arm length)



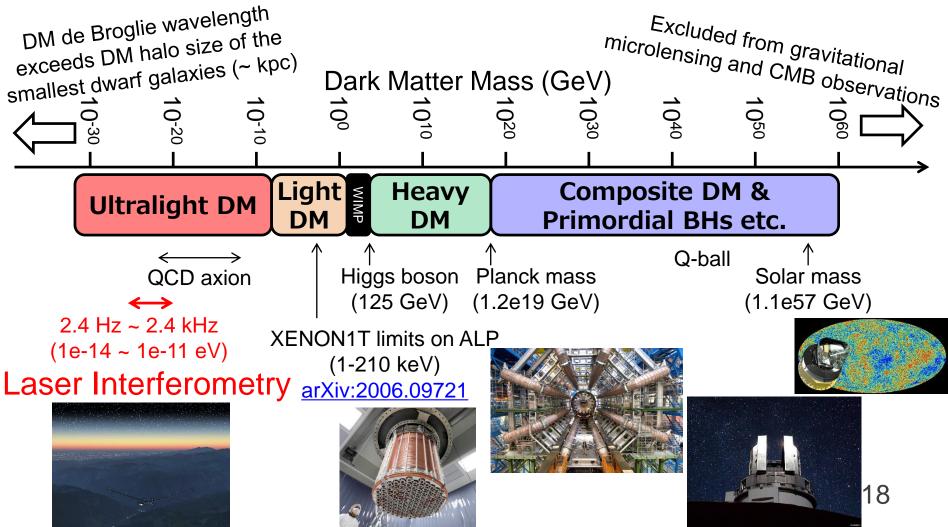
Key Aspects to Remember

- Michelson interferometer measure the differential length between two arms
 - insensitive to common length changes
- Optical cavities measure the distance (optical path length) between mirrors
 - insensitive to common displacements
- They are also sensitive to the changes in the speed of light $\frac{\delta L}{L} = \frac{\delta c}{c}$
- They are not sensitive to translational motion of mirrors (to the first order)

Ultralight dark matter searches

Dark Matter Models

- ~90 orders of magnitude
- Ultralight DMs behave as classical wave fields



Various Proposals

Axion-like particles

- W. DeRocco & A. Hook, PRD 98, 035021 (2018)
- I. Obata, T. Fujita, YM, PRL 121, 161301 (2018)
- H. Liu+, PRD 100, 023548 (2019)
- K. Nagano, T. Fujita, YM, I. Obata, PRL 123, 111301 (2019)
- D. Martynov & H. Miao, PRD 101, 095034 (2020)

Scalar fields

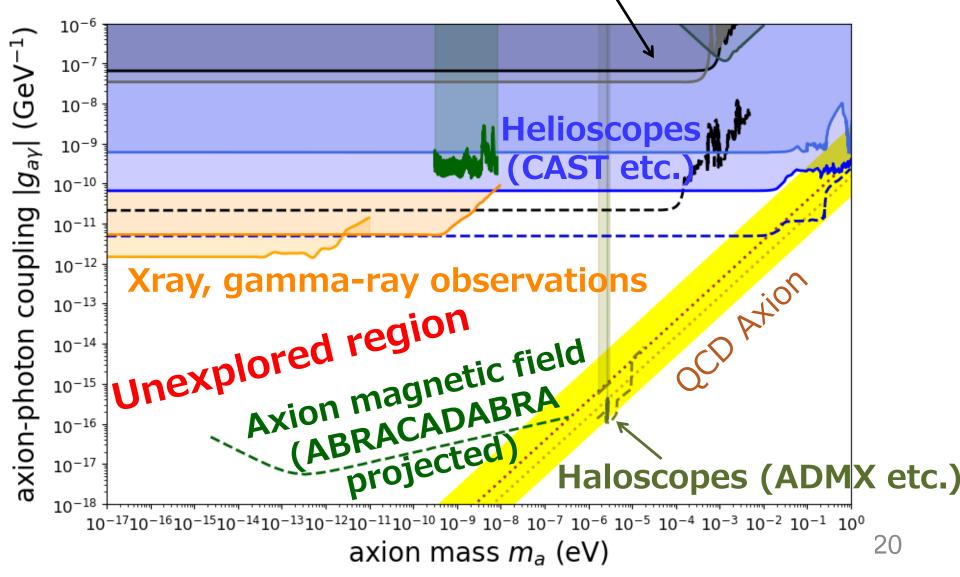
- Y. V. Stadnik & V. V. Flambaum, PRL 114, 161301 (2015)
- Y. V. Stadnik & V. V. Flambaum, PRA 93, 063630 (2016)
- A. A. Geraci+, PRL 123, 031304 (2019)
- H. Grote & Y. V. Stadnik, PRR 1, 033187 (2019)
- S. Morisaki & T. Suyama, PRD 100, 123512 (2019)

U(1)_B or U(1)_{B-L} gauge bosons

- P. W. Graham+, PRD 93, 075029 (2016)
- A. Pierce+, PRL 121, 061102 (2018)
- D. Carney+, <u>arXiv:1908.04797</u>

Not exhaustive. There are also proposals for heavier DM (I think they are not promising). The ones which require magnetic fields are not listed.

Search for Axion-Photon Coupling Light Shining through Wall (ALPS etc.)



Velocity of Circular Polarizations

• Axion-photon coupling $(\frac{g_{a\gamma}}{4}aF_{\mu\nu}\tilde{F}^{\mu\nu})$ gives different phase velocity between left-handed and right-handed circular polarizations

$$c_{\mathrm{L/R}} = \sqrt{1 \pm \frac{g_{a\gamma}a_0m_a}{k}} \sin(m_a t + \delta_{\tau})$$
coupling constant axion field axion mass

 Measure the difference as resonant frequency difference in an optical cavity

Laser

 $\nu_{\rm R}$

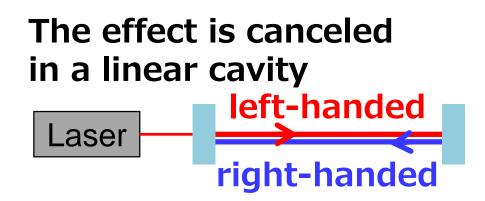
$$\frac{\delta c}{c} = \frac{\nu_{\rm L} - \nu_{\rm R}}{\nu}$$

• Search can be done without magnetic field

Our Ideas

Use of bow-tie cavity





Not canceled in a bow-tie cavity

left-handed

Laser

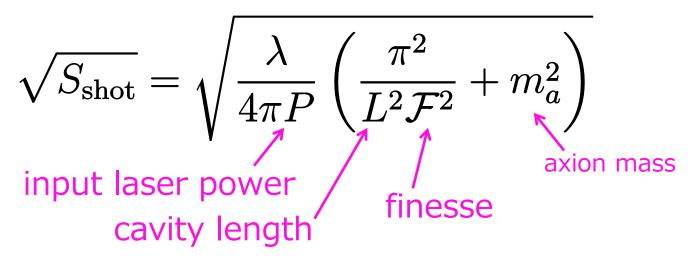
Use of double-pass configuration
 Transmitted beam is reflected back into the same cavity as different polarization to realize a null measurement of the resonant frequency difference
 Y.M+, PRL 110, 200401 (2013)

Double-Pass Configuration

 Axion signal is extracted from the cavity reflection (null measurement) **Frequency servo Photodiode** -dser CW laser $\nu_{\rm T}$ left-handed High common mode rejection due to the common path $\nu_{\rm R}$ right-**Double-pass Axion signal** handed configuration $-\nu_{\rm F}$ 23

Sensitivity Calculation

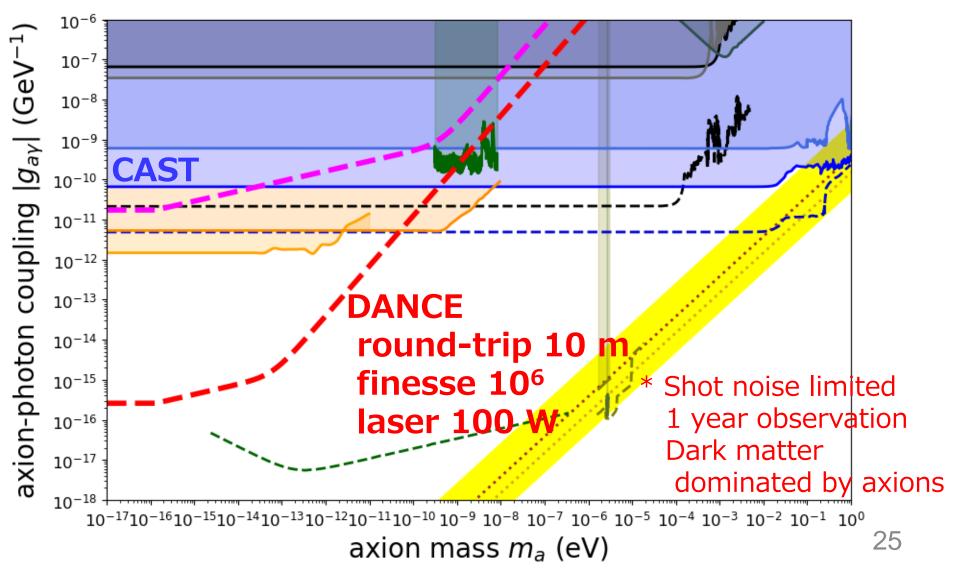
- Cavity length changes (displacement noises) will not be a fundamental noise due to common mode rejection
- Ultimately limited by quantum shot noise



 Sensitivity to axion-photon coupling can be calculated by assuming axion density = dark matter density

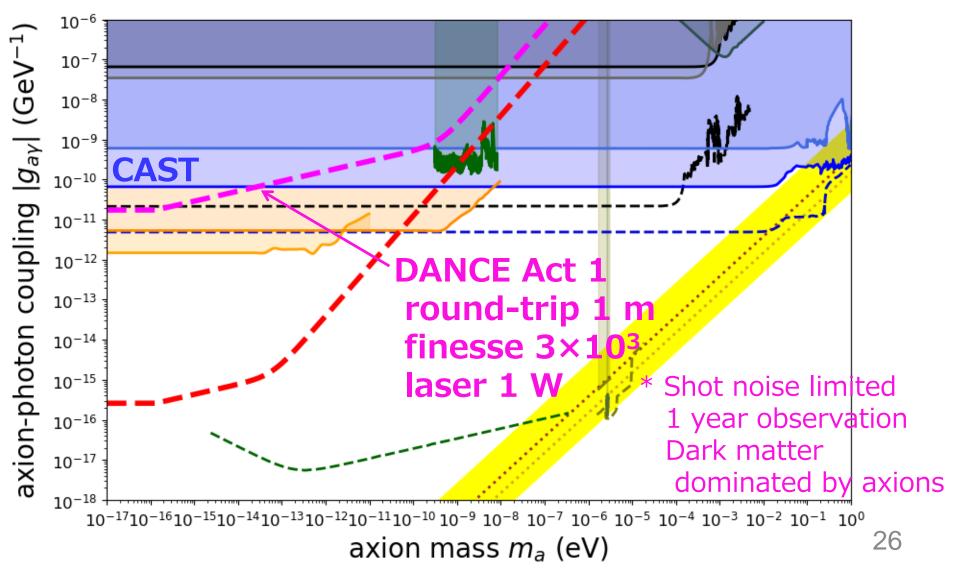
Search for Unexplored Region

Dark matter Axion search with riNg Cavity Experiment



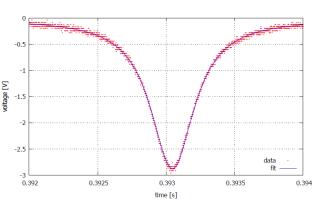
Prototype Experiment

Dark matter Axion search with riNg Cavity Experiment



DANCE Act 1

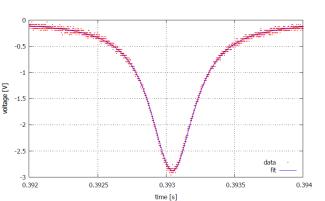
- Completed the assembly of optics
- Finesse measured to be 515 +/- 6 (design: 3×10^3)
- Having trouble with stable lock
- Aiming for first run in 2020

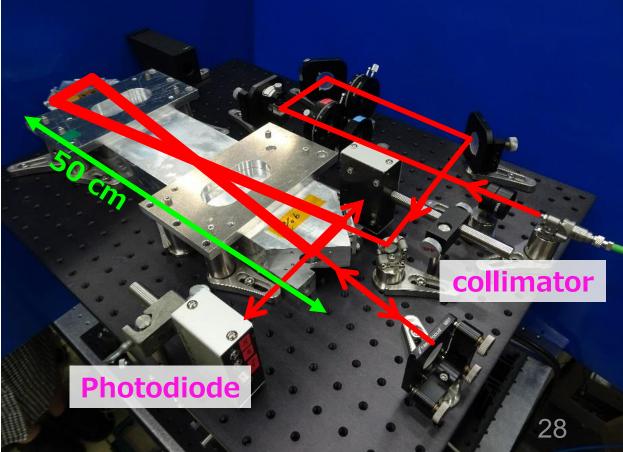




DANCE Act 1

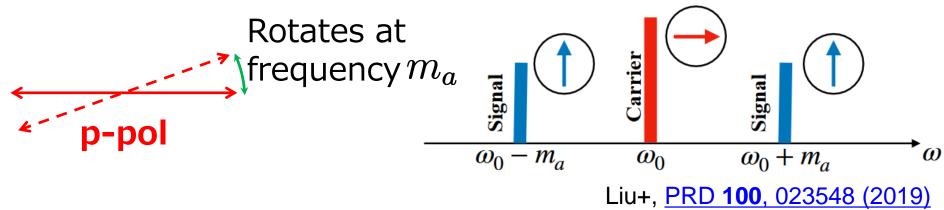
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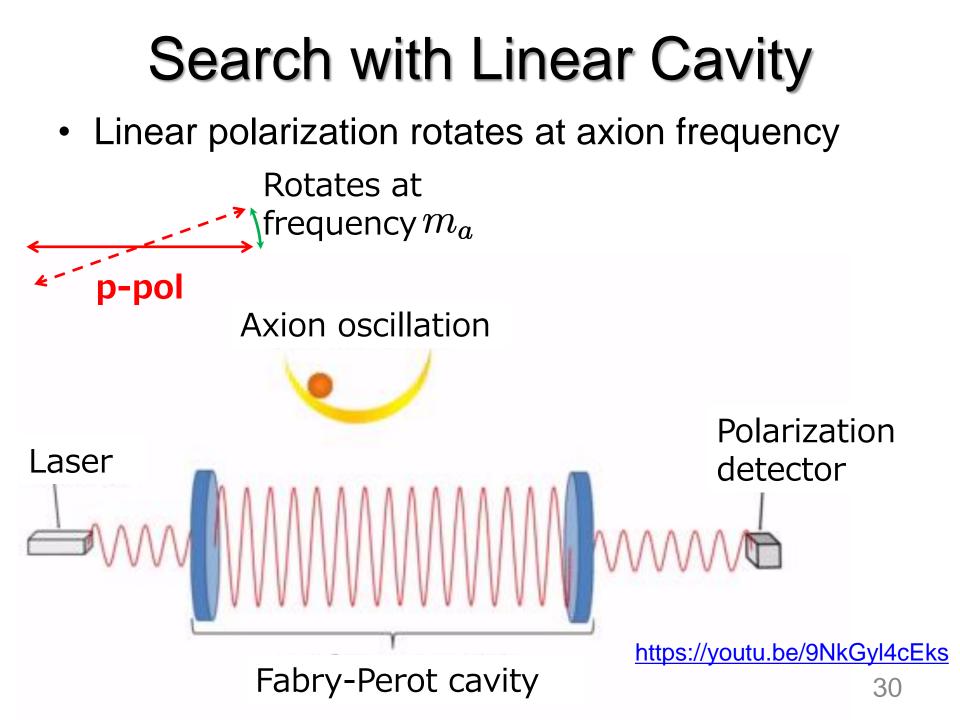
Search with Linear Cavity

• Linear polarization rotates at axion frequency



 Sensitive when axion oscillation period and roundtrip time of optical cavity is the same





Linear Cavity in GW Detectors

- Suitable because of 10⁻¹⁰
 long arms and high 10⁻¹¹
 power
 long 10⁻¹¹
- Can be done simultaneously with GW observation
- Considering of applying to KAGRA

Helioscope KAGRA SN1987A DECIGC **Cosmic Explorer** buildno 10⁻¹⁴ 10⁻¹⁶ 10⁻¹⁷ 10⁻¹⁶ 10⁻¹⁵ 10⁻¹¹ 10⁻¹⁴ 10⁻¹³ 10⁻¹² 10⁻¹⁰ Axion mass [eV]

 10^{-9}

S-pol S-pol S-pol S-pol (GW signal) (Axion signal)

Other Recent Proposals

 There are also different proposals for axion dark matter search with laser interferometers DeRocco & Hook, PRD 98, 035021 (2018); Liu+, PRD 100, 023548 (2019); Martynov & Miao, PRD 101, 095034 (2020) frequency (Hz) Laser 10^1 10^2 10^3 10^4 $10^{-2} \ 10^{-1} \ 10^{0}$ 10⁵ 10^{6} 10^{7} 10^{-9} FIG. 3. A diagram of our proposed axion interferometer \BRA_10cm Laser 10^{-10} CAST ALPS-1 10^{-11} īΣ₹Č 10^{-12} $\ell = 2 \text{ m}$ Axion Interferometr ADBO **ADBC Experiment** 10-13 D PD FIG. 2: Schematic of the ADBC experiment. The red optical 10^{-14} DANCE Act-1 Advanced LIGC = 1 m. $F = 3 \times$ cavity 10^{-15} Main cavity 10^{-16} ABRACADABBA Squeezed light S-polarised -P-polarised Auxiliary beam 10^{-17} Half waveplate Polarising beam splitter Jantum Enhanced 10^{-18} $10^{-17}10^{-16}10^{-15}10^{-14}10^{-13}10^{-12}10^{-11}10^{-10}10^{-9}10^{-8}10^{-7}$

axion mass *m* (eV)

32

YM+, JPCS 1468, 012032 (2020)

(GeV

axion-photon coupling $g_{a\gamma}$

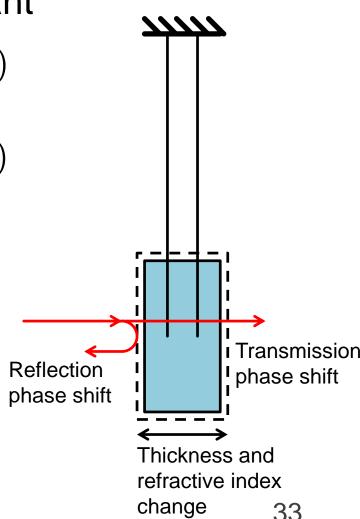
Search for Scalar Dark Matter

 Dilaton-like scalar DM drives oscillations in electron mass and fine structure constant

$$\frac{\delta m_{\rm e}}{m_{\rm e}} = \frac{1}{\Lambda_f} \phi_0 \cos\left(m_{\phi} t - \vec{k} \cdot \vec{r}\right)$$
$$\frac{\delta \alpha}{\alpha} = \frac{1}{\Lambda_{\gamma}} \phi_0 \cos\left(m_{\phi} t - \vec{k} \cdot \vec{r}\right)$$

- This drives oscillations in the Bohr radius $a_0 = \frac{\hbar}{\alpha m_e c}$
- The size and refractive index of mirrors changes

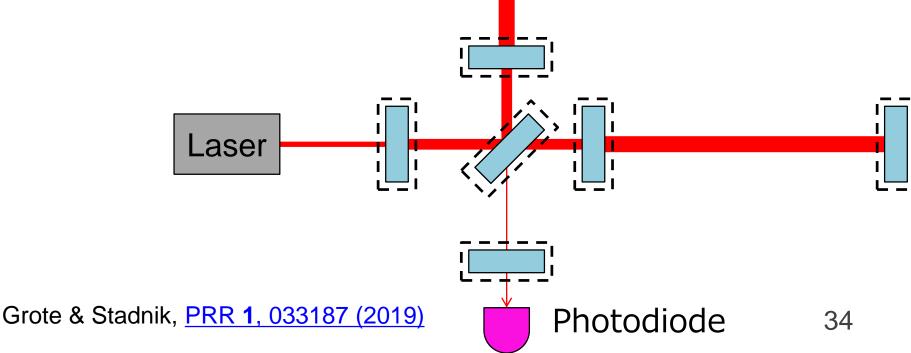
$$\frac{\delta l}{l} = \left(-\frac{\delta \alpha}{\alpha} - \frac{\delta m_{\rm e}}{m_{\rm e}}\right)$$



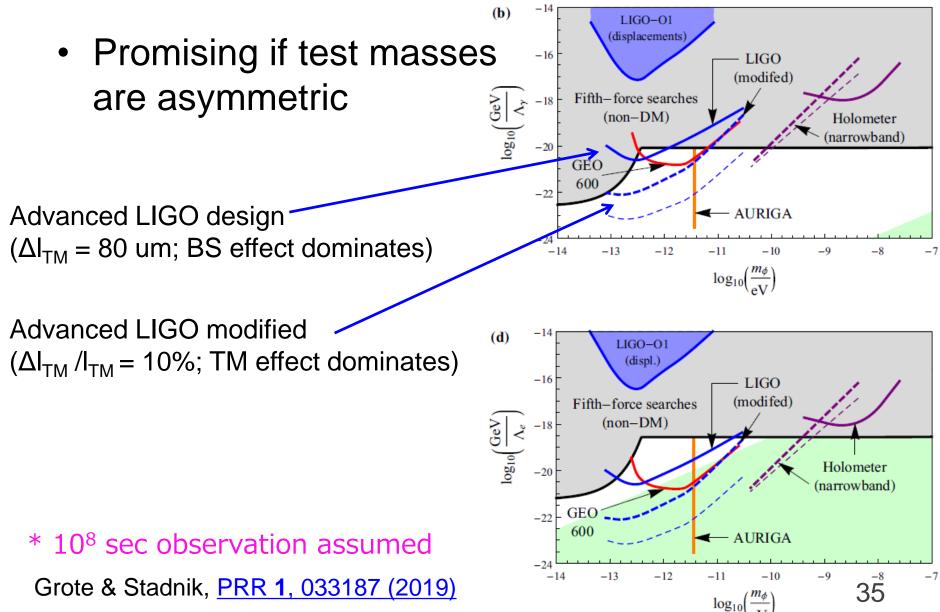
Search with GW Detectors

- Thickness changes in beam splitter is attenuated by $N_{\rm eff}=\frac{2\mathcal{F}}{-1}$
- Changes in arm length are common
- Sensitive only if the thickness of test masses are asymmetric

 π



Sensitivity to Scalar DM



Search for Vector Dark Matter

- Dark gauge bosons (e.g. dark photon)
- New gauge symmetry: B-L (baryon number minus lepton number)
 - B-L is conserved in standard model
 - B-L could be the charge for dark gauge boson
- Gauge bosons give acceleration to mirrors

Dimensionless coupling strength Dark charge

$$\vec{a}(t, \vec{r_i}) = \epsilon_{\mathrm{B-L}} e \frac{q_{\mathrm{B-L},i}}{M_i} m_{\mathrm{A}} \vec{A_0} \cos\left(m_{\mathrm{A}} t - \vec{k}_{\mathrm{A}} \cdot \vec{r_i}\right)$$

• Arm cavity differential length change If q/M is the same for all the mirrors

$$\sqrt{\langle \delta L^2 \rangle} = \frac{\sqrt{2}}{3} \frac{|a|kL}{m_A^2}$$
 ~10⁻⁶ for r and

Averaged over all directions of a and k

A. Pierce+, <u>PRL 121, 061102 (2018)</u> 36

for mA = 100 Hz = 4e-13 eV

and L=4 km (Advanced LIGO)

Sensitivity to $U(1)_{B-L}$ Gauge Boson

- The effect to $\vec{a}(t, \vec{r}_{\rm EY})$ differential arm length change depend on the direction of \vec{k} and \vec{a}
- Length change rely on phase $\vec{a}(t, \vec{r}_{\text{IY}}) \rightarrow \vec{a}$ differences

Laser

No change in Y-arm length

 $\vec{a}(t, \vec{r}_{\mathrm{IX}})$

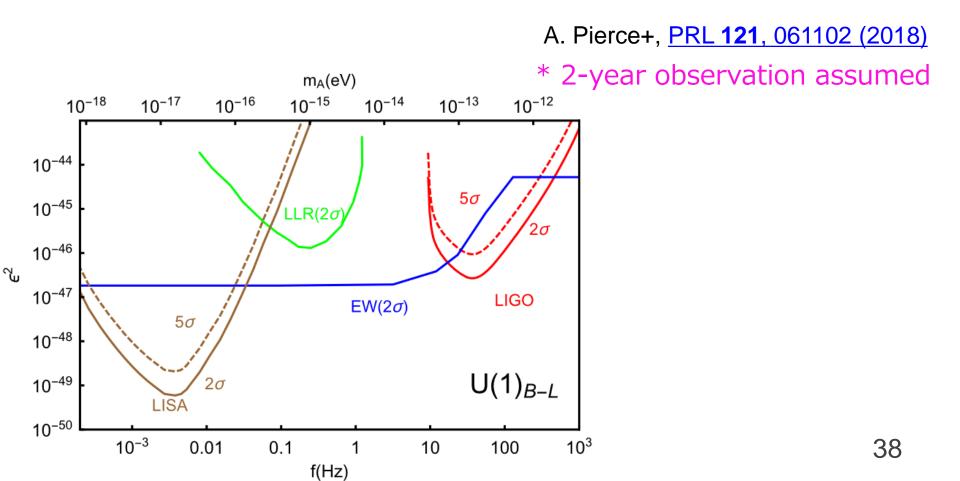
 \vec{k}



 $\vec{a}(t, \vec{r}_{\mathrm{EX}})$

Sensitivity to U(1)_{B-L} Gauge Boson

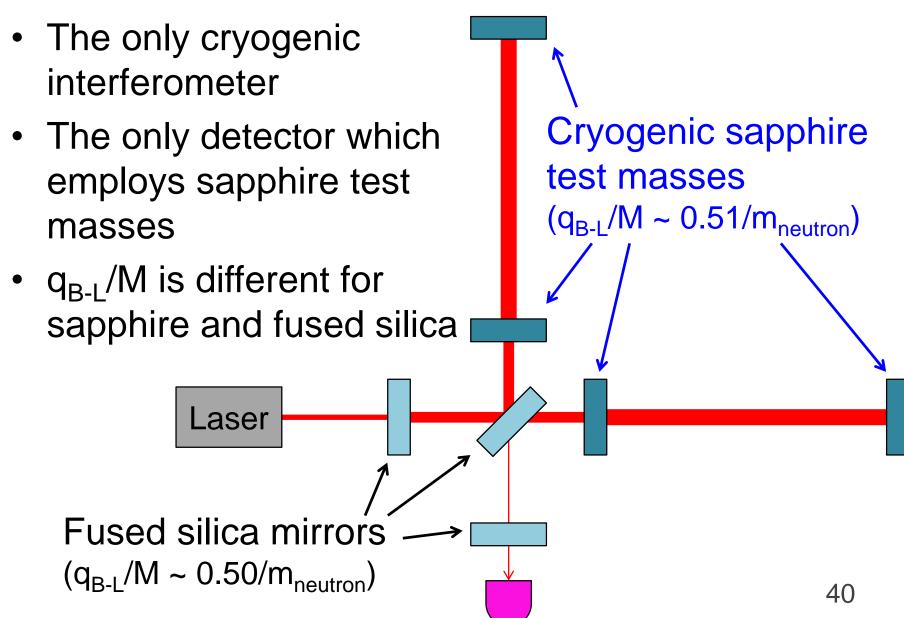
Sensitivity can be better than equivalence principle tests with torsion pendulum



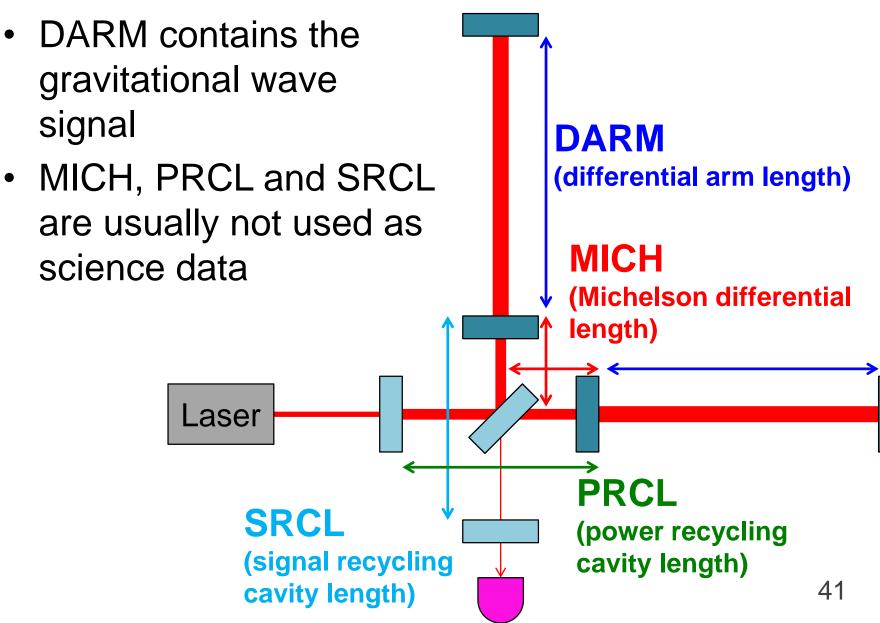
Our New Idea

- Issues with previous proposals for scalar and vector dark matter search
 - the effect is common to test masses
 - the effect is common to both arm cavities
- The sensitivity rely on slight asymmetry in the arm cavities or small phase differences in distant test masses
- How about using auxiliary length signals to enhance the sensitivity?
 - power/signal recycling cavity
 - Michelson interferometer

KAGRA Interferometer

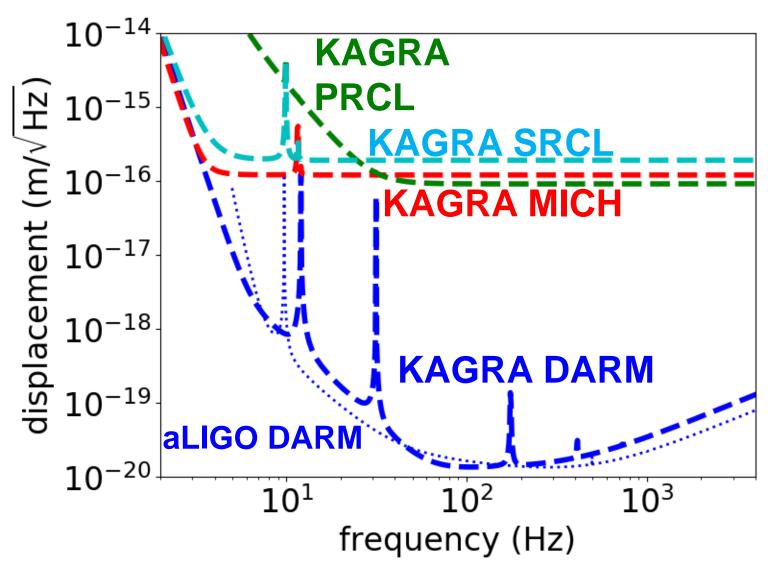


Auxiliary Length Signals



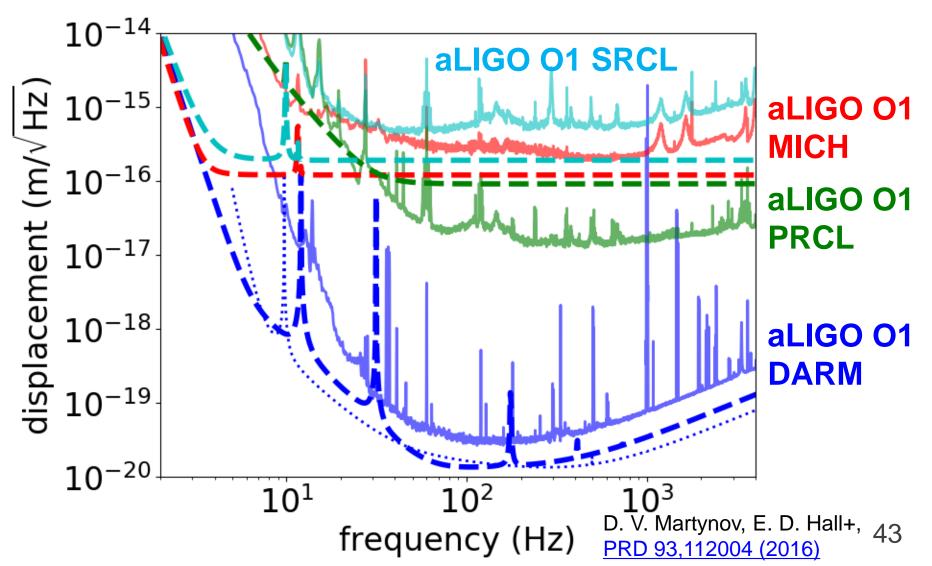
Displacement Sensitivity

• Auxiliary signals are not so sensitive at high freq.



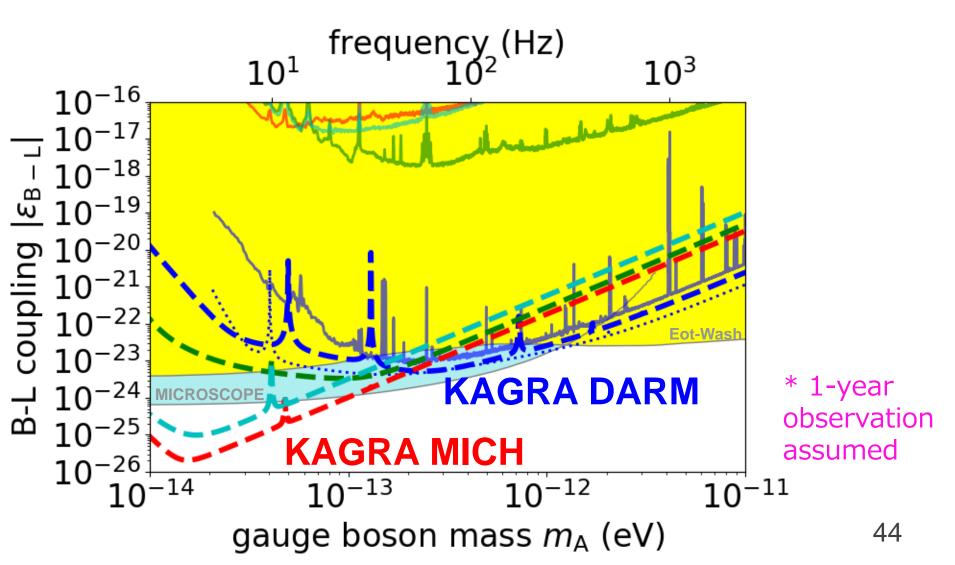
Displacement Sensitivity

• Auxiliary signals are not so sensitive at high freq.



Sensitivity to U(1)_{B-L} Gauge Boson

Auxiliary signals can have better sensitivity



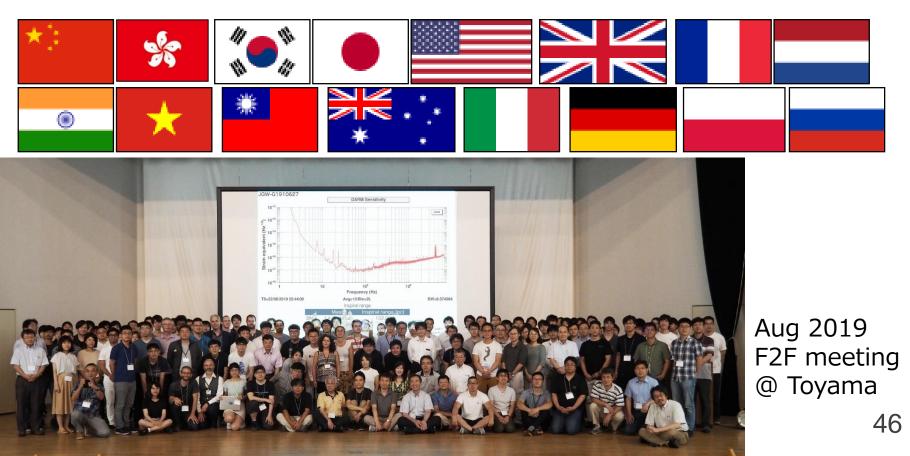
Dark matter search with KAGRA

KAGRA Project KAGRA

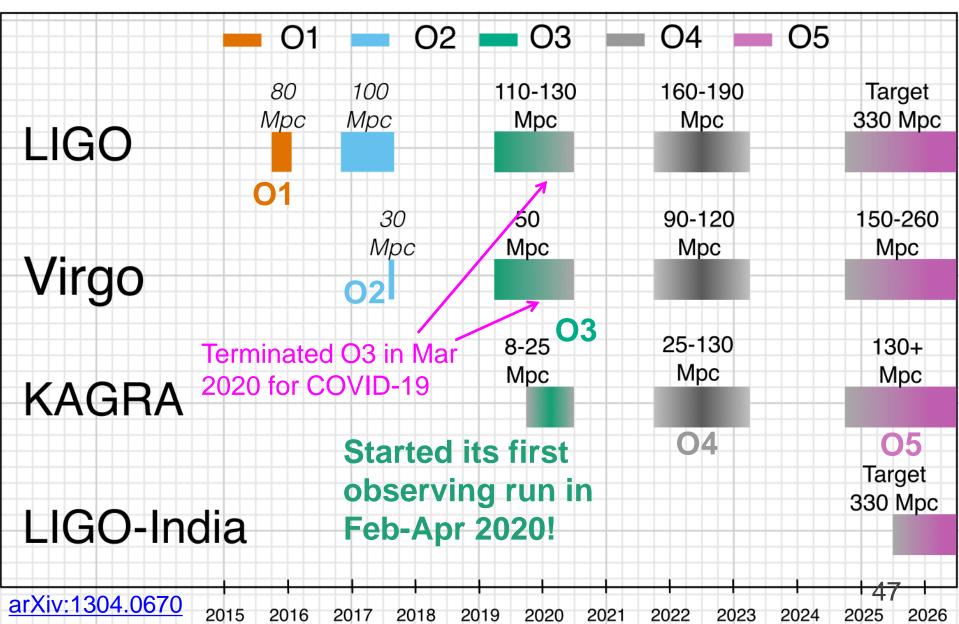


- Budget approved in 2010
- 110 institutes, 450+ collaborators (200 authors)
- Cryogenic and underground

Join us!



Observing Plans



First KAGRA Observing Run

- Officially started on February 25, 2020
- But soon stopped to resume interferometer tuning
- Observing run restarted on April 7 to April 21
 - with ~0.6 Mpc in binary neutron star range
 - ~170 hours of science data

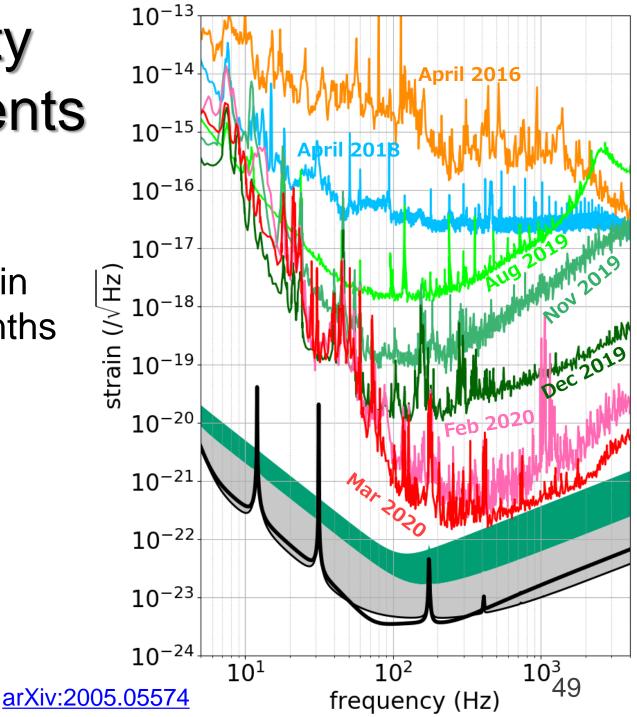
arXiv:2005.05574

 Joint observation with LIGO and Virgo was not possible



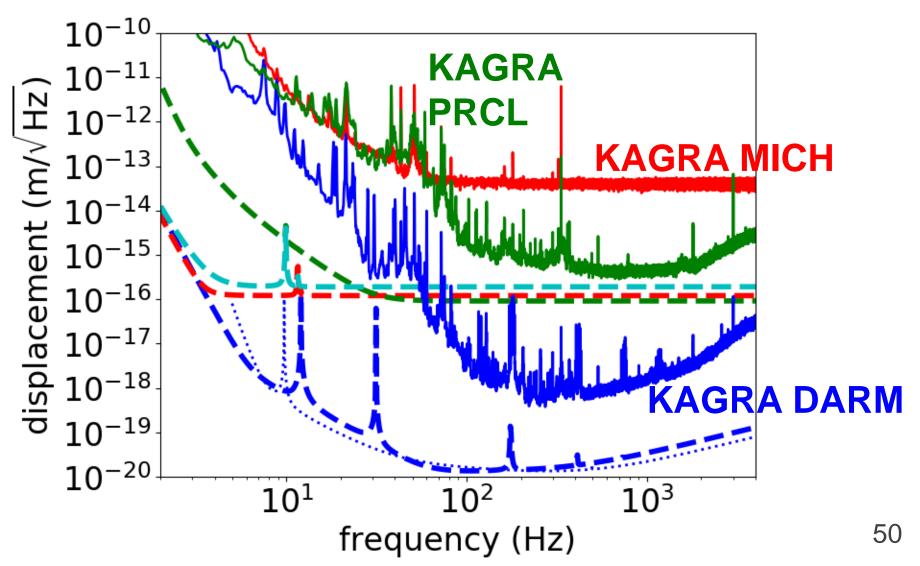
Sensitivity Improvements

 Dramatic sensitivity improvement in the last 7 months



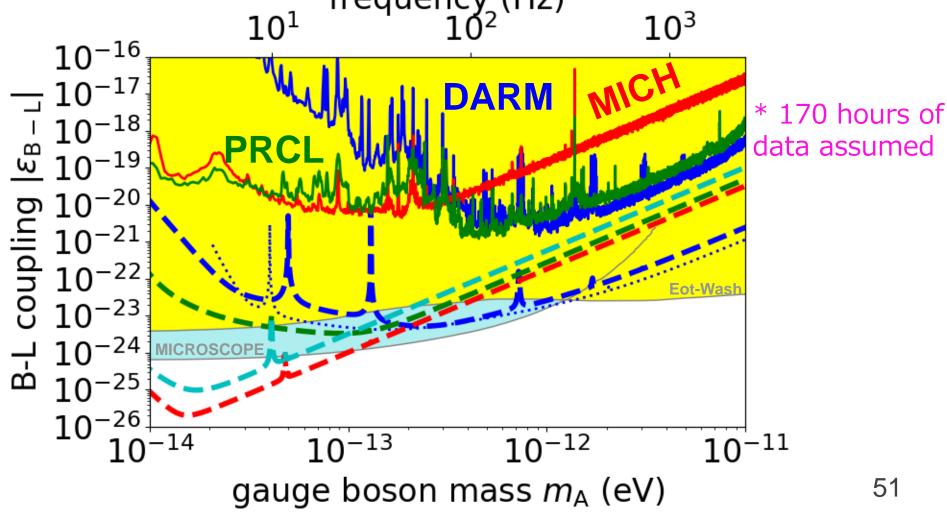
Best Sensitivity on March 20, 2020

~2 orders of magnitude worse than the designed



Prospected Sensitivity for U(1)_{B-L}

 Actual limit would be worse due to sensitivity fluctuations and intermittent science mode frequency (Hz)



Plans for Dark Matter Search

- Development of data analysis code underway
- First try with O3 data

- this will be the first search for U(1)_{B-L} gauge boson with laser interferometer

- Planning to install polarization optics to search for axion-like particles by O4 (~2021)
- Stay tuned for dark matter signals from gravitational wave detectors!

Summary

- Laser interferometers are attractive tools to search for ultralight dark matter
- Axion-like particles can be searched with polarization measurements
- Scalar fields can be searched through variations in the optical thickness of mirrors
- Gauge bosons can be searched through measuring forces acting on mirrors
- KAGRA can do unique searches because of the use of sapphire mirrors