

Microwave Searches for Dark Matter

Axions, Shuket and all that ...

Tokyo-Paris seminars
September, 17th, 2021

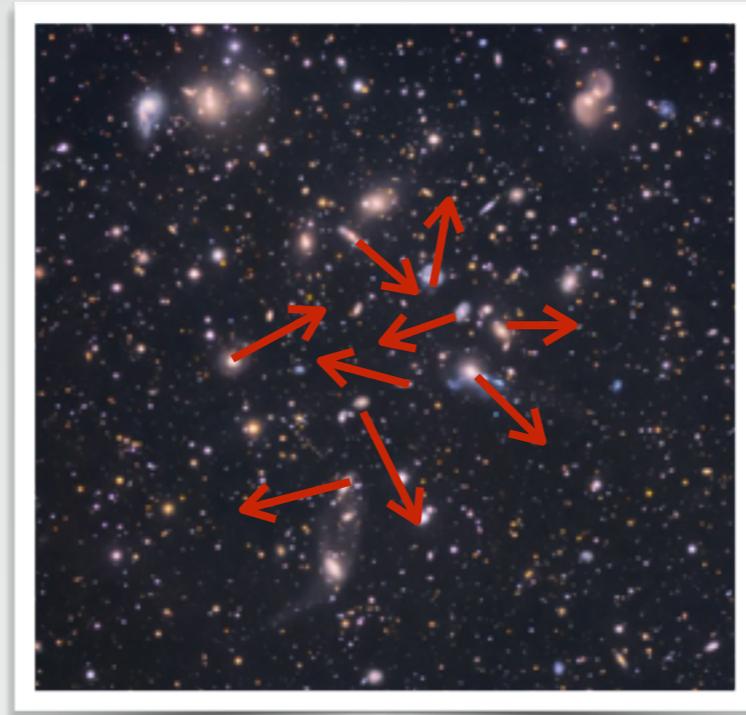
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CEA Paris-Saclay

Dark matter at various scales

Galaxies

Galaxy clusters



Rotation curves

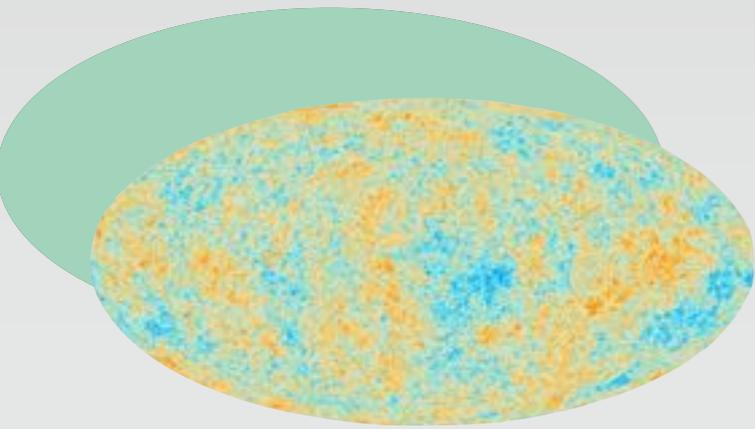
Velocity dispersion

Collision

New, dark component required

Dark matter at various scales

CMB emitted when first atoms formed

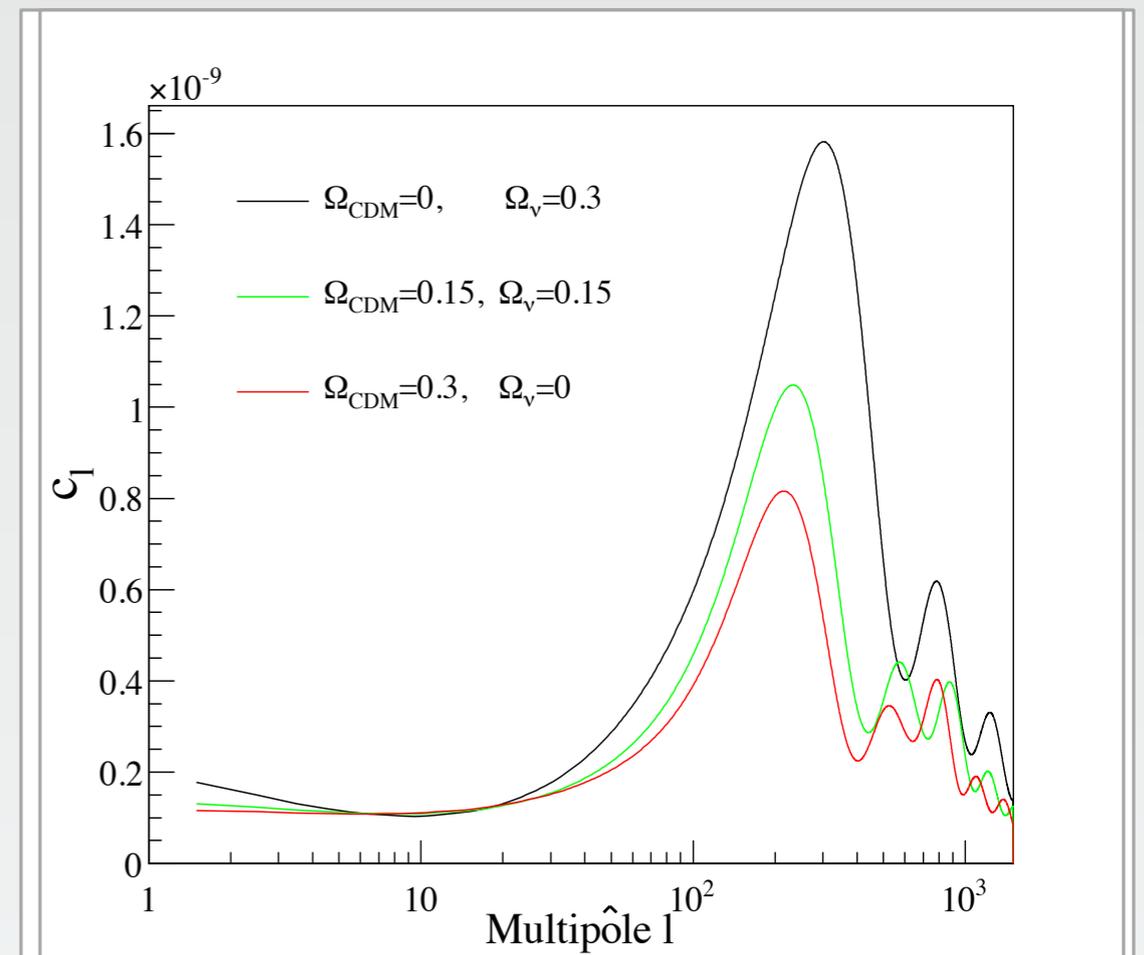
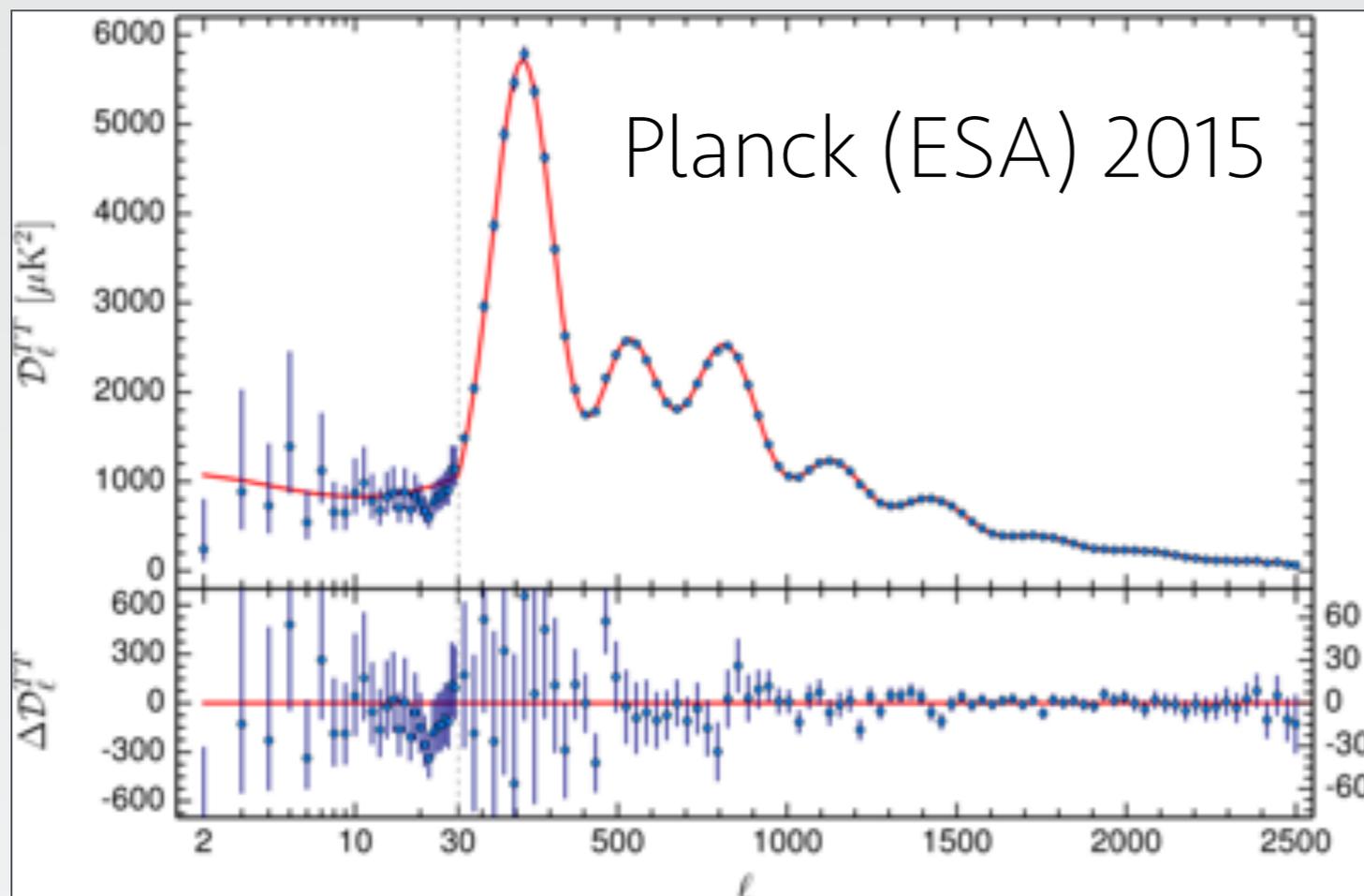


Very homogeneous and isotropic

Density oscillations before recombination



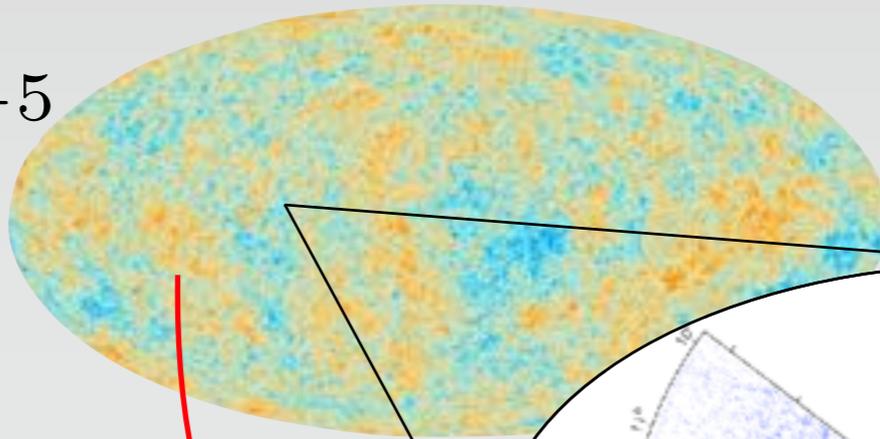
Temperature fluctuations



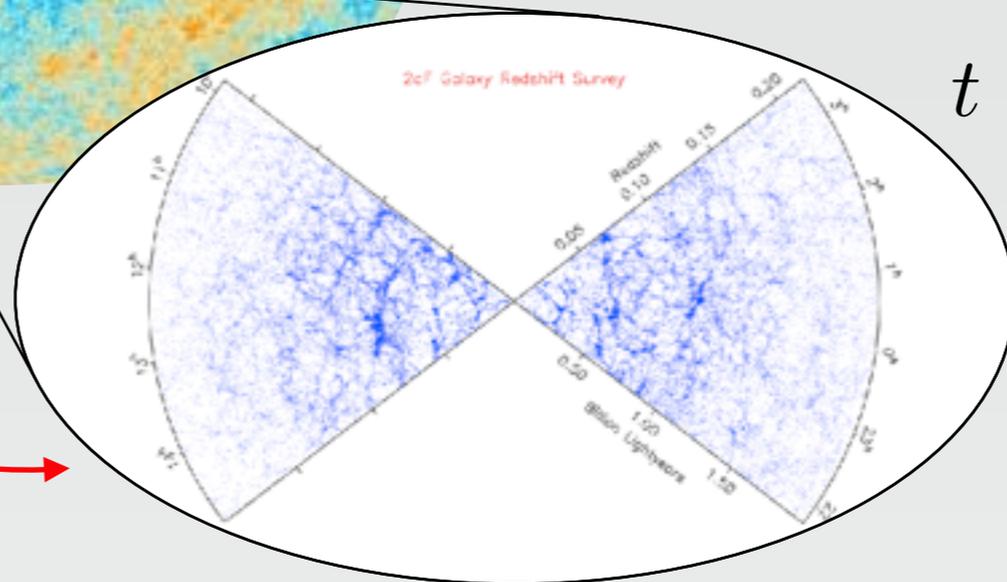
Dark matter at various scales

$$t = 4 \times 10^5 \text{ y}, a(t) \simeq 1/1000$$

$$\delta\rho/\rho \sim 10^{-5}$$



$$\delta\rho/\rho \propto a(t)$$



$$t = 10^{10} \text{ y}, a(t) \simeq 1$$

$$\delta\rho/\rho \gg 1$$

New type of matter required to explain growth of structures

Massive

Weakly interacting

Stable

Cold

84 % of total mass

CP symmetry

- ★ CP = Charge x Parity
- ★ CP violation required for matter domination
- ★ Weak interactions known to violate CP
- ★ Strong interactions, electromagnetic interactions, gravitation known to conserve CP

The strong CP problem

- ★ In general, QCD contains a term

$$\mathcal{L}_\theta = \frac{\theta}{32\pi^2} \text{Tr}(G_{\mu\nu} \tilde{G}^{\mu\nu})$$

$\theta \neq 0$ implies $\left\{ \begin{array}{l} \text{CP violation in strong sector} \\ \text{Non-vanishing neutron electric dipole} \end{array} \right.$

- ★ Strong CP violation not observed

- ★ Neutron electric dipole moment :

$$d_n \leq 2.9 \times 10^{-26} \text{ e.cm} \Rightarrow \theta \leq 10^{-10}$$

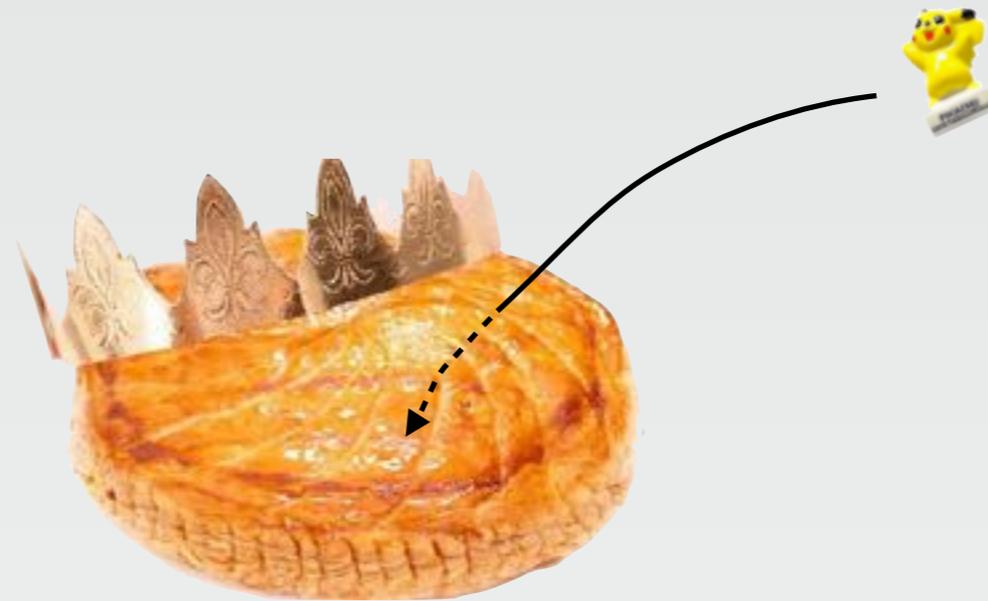
↑
Fine tuning problem

The strong θ parameter

- ★ θ is an angle
- ★ Why is $\theta=0$ a problem?

The strong θ parameter

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- ★ Why is $\theta=0$ a problem?



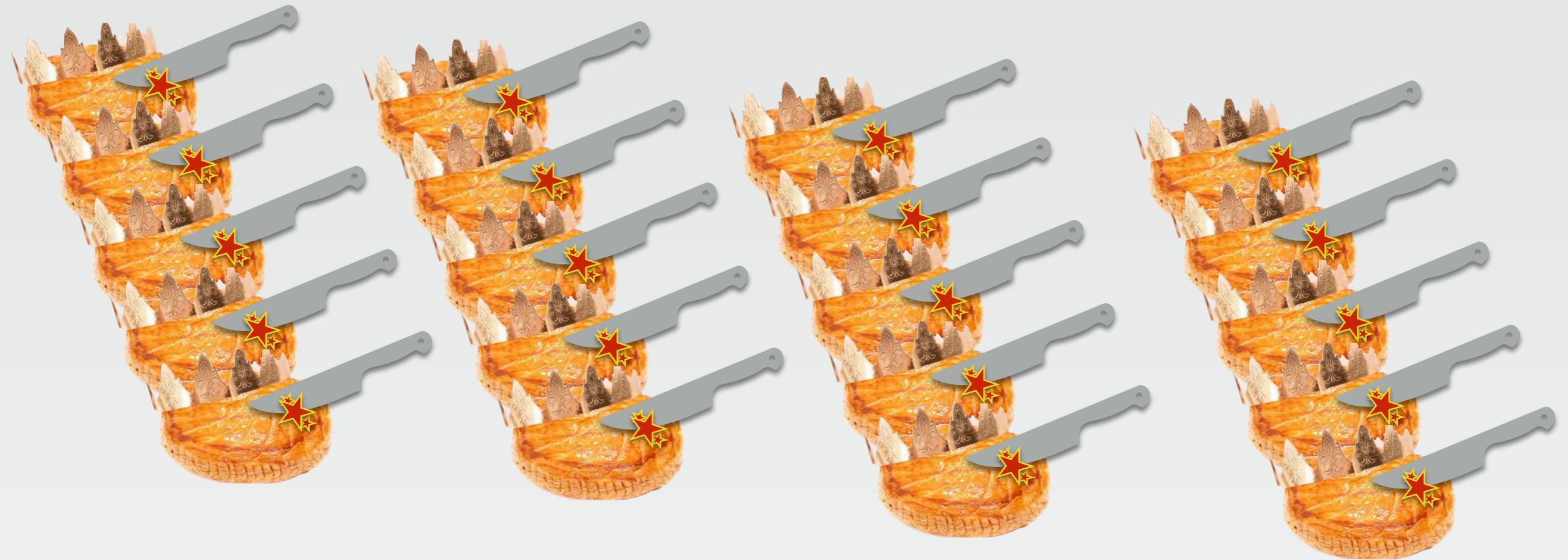
The strong θ parameter

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The strong θ parameter

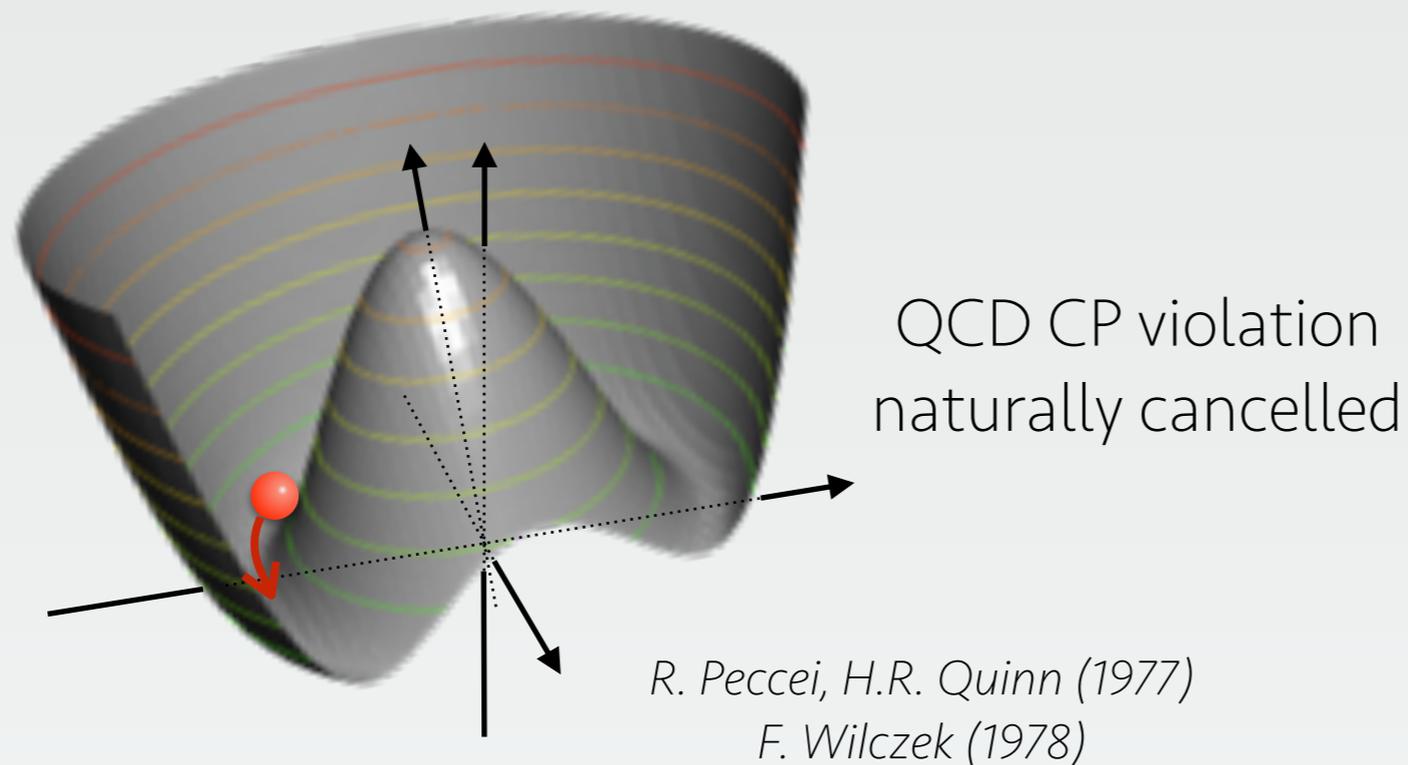
- ★ θ is an angle
- ★ Why is $\theta=0$ a problem?



An explanation is required: new force? new symmetry?

A theoretical prediction: the axion

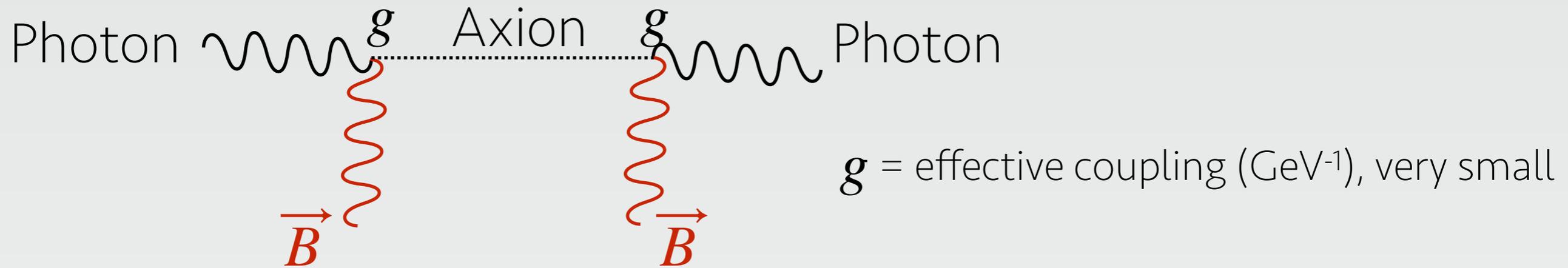
- ★ New symmetry is proposed
- ★ Strong CP is naturally suppressed
- ★ Cost: new particle



The Axion

Properties of the axion

- ★ Very light ($< eV/c^2$, electron $5 \times 10^5 eV/c^2$, proton $10^9 eV/c^2$)
- ★ Axions couple to photons in a magnetic field



- ★ Could modify stellar lifetime
- ★ Can be the dark matter!

Axions and axion-like particles

★ Axion models : θ is a dynamical field

- Minimum of the potential for $\langle \theta \rangle = 0$

$$\mathcal{L} = \frac{C_\gamma}{32\pi^2} \frac{\theta}{f} \text{Tr}(F_{\mu\nu} \tilde{F}^{\mu\nu}) \propto \theta \vec{E} \cdot \vec{B} \quad m_\theta \propto 1/f$$

Axions

★ Axions are also natural dark matter candidate :

- Extra dimensions 2 birds w/ 1 stone :

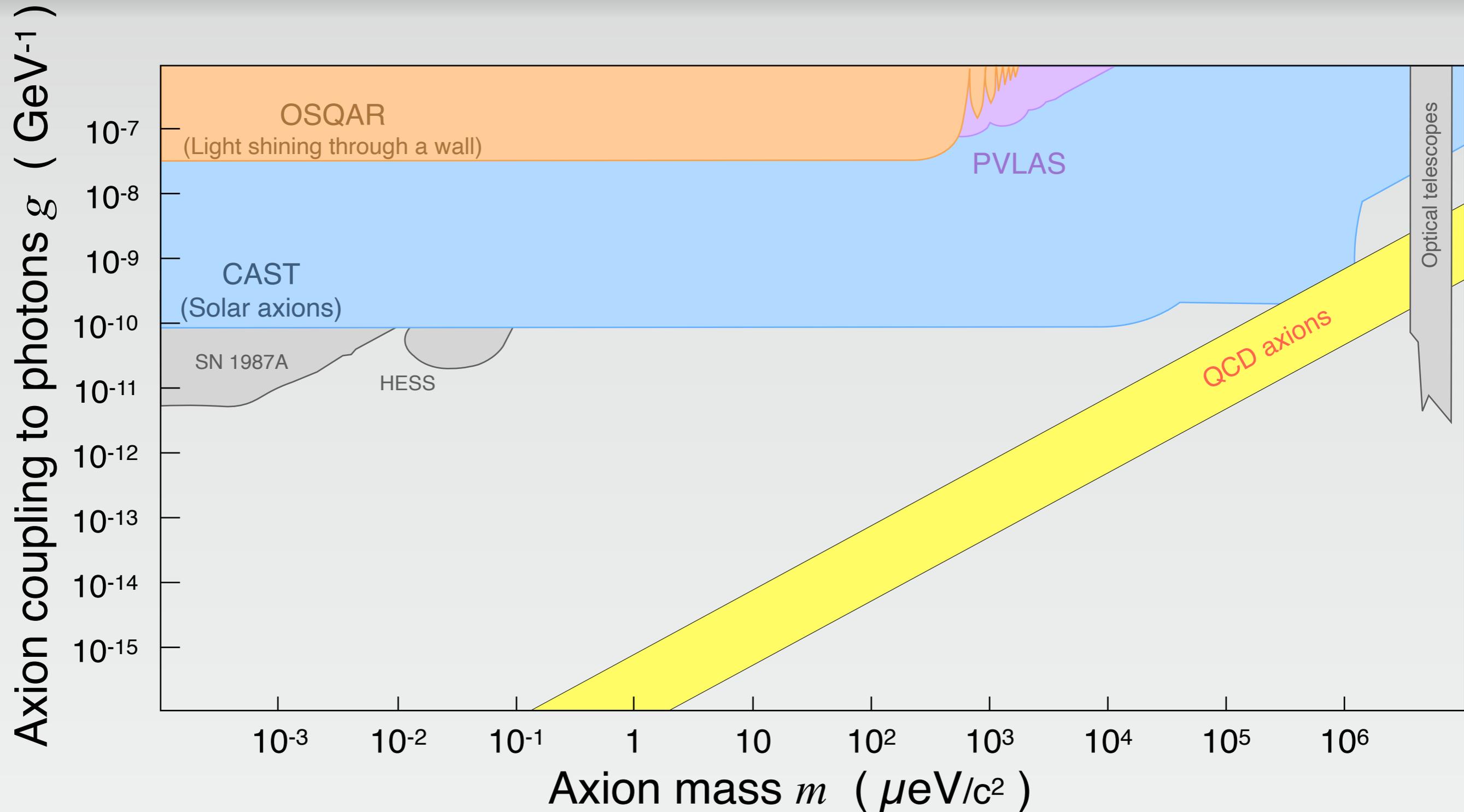
Strong CP and Dark Matter

- String theory, dilatons, etc.

$$m_\theta \propto 1/f$$

Axion-like
particles

Axion searches and constraints



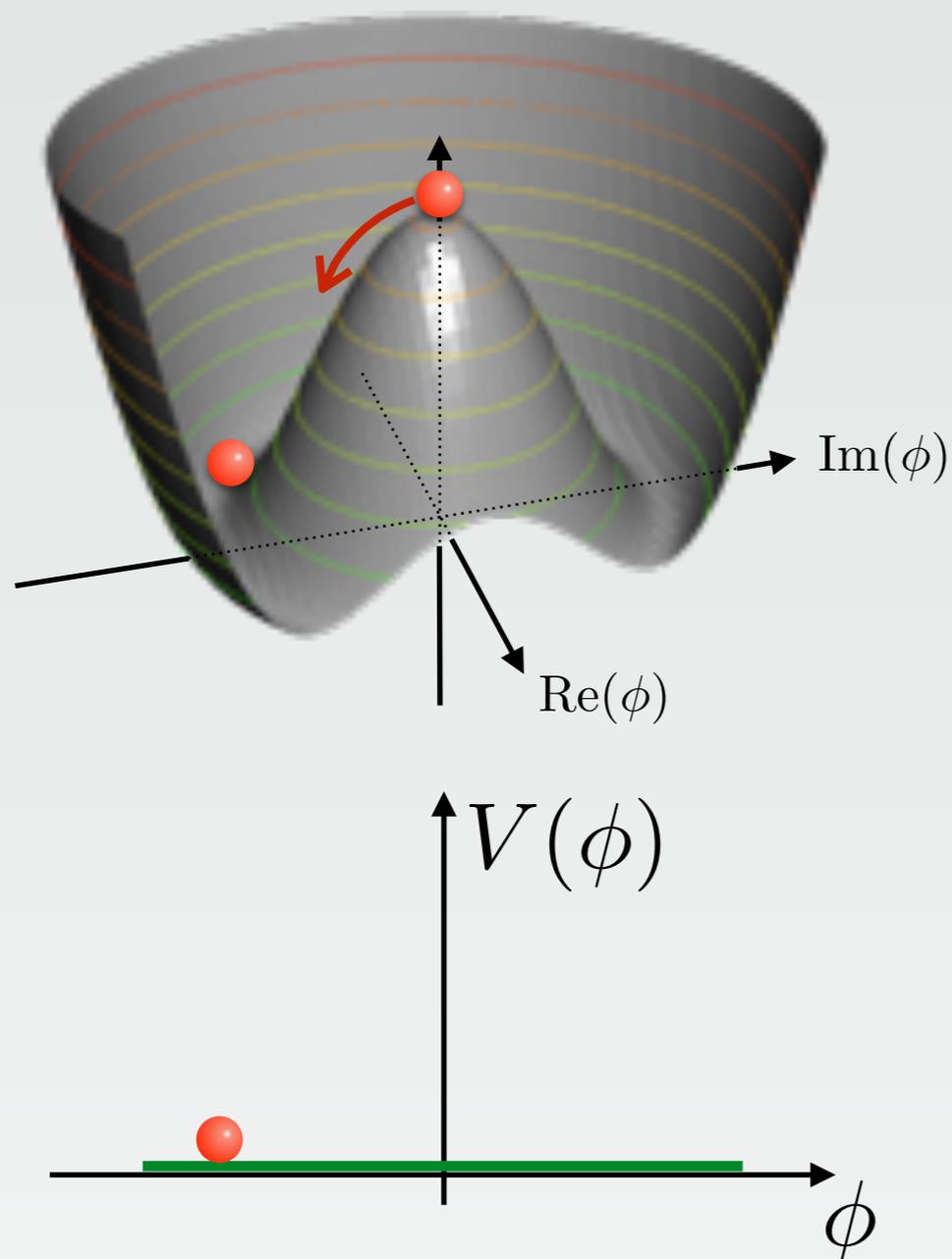
Axion Relic Density

- ★ Thermal population irrelevant to CDM
- ★ Different production mechanisms
 - Misalignment
 - Decay of topological defects if any
- ★ 2 classes of scenarios :
 - PQ symmetry breaking before/after inflation

Cosmological Axions

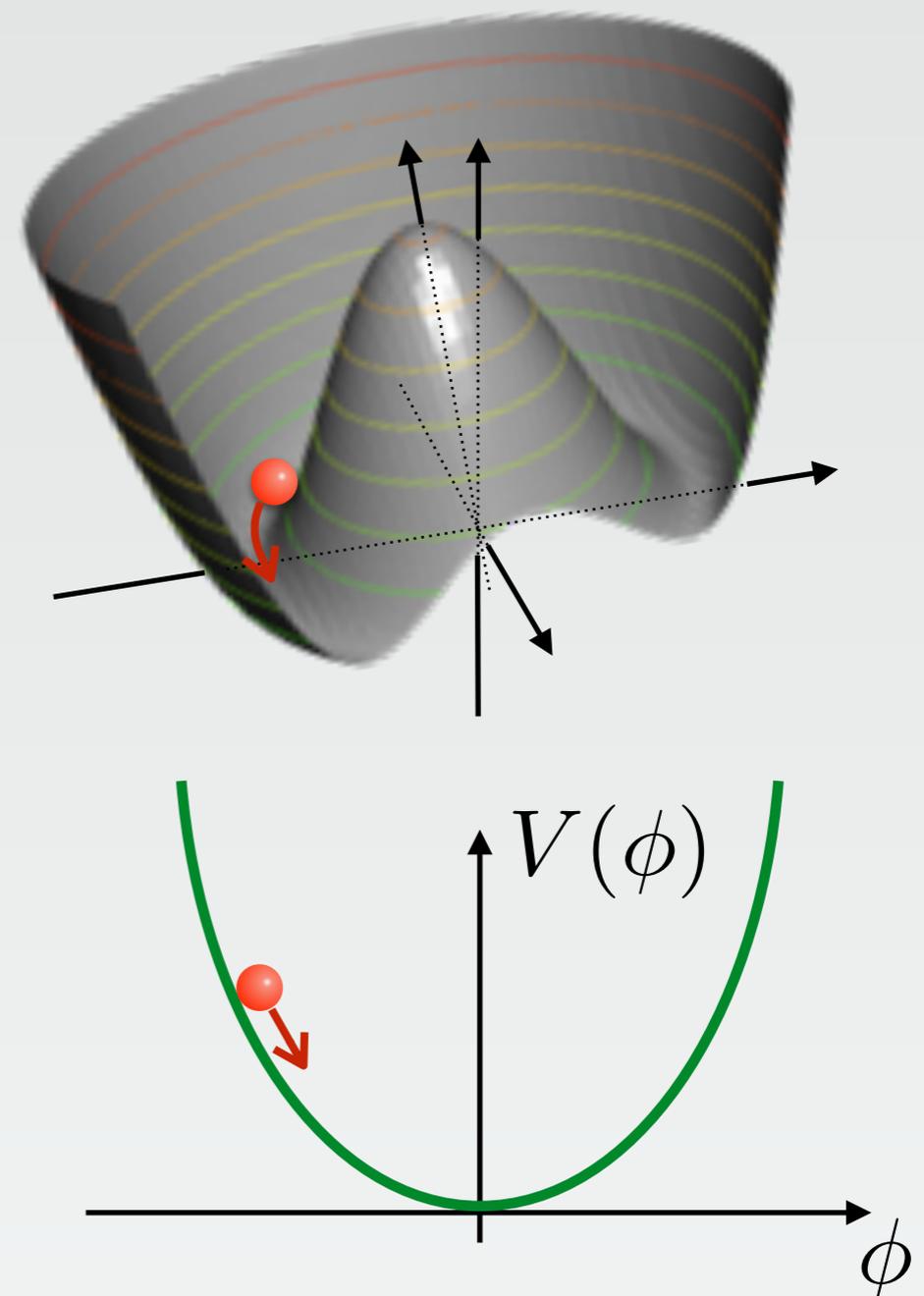
$$T = f_a \gg \Lambda_{QCD}$$

Breaking of Peccei-Quinn $U(1)_{PQ}$



$$T \sim \Lambda_{QCD}$$

QCD phase transition



Axion relic density

★ Occupation number is huge

★ Model as $\mathbf{k} = \mathbf{0}$ mode of a classical field:

$$\ddot{\phi} + 3H(t)\dot{\phi} + m^2\phi = (\text{terms with spatial derivatives}) = 0$$

★ Energy density: $\rho = \frac{1}{2}\dot{\phi}^2 + \frac{m^2}{2}\phi^2$

★ Pressure: $p = \frac{1}{2}\dot{\phi}^2 - \frac{m^2}{2}\phi^2$

★ Equipartition $\Rightarrow \frac{1}{2}\langle\dot{\phi}^2\rangle = \frac{m^2}{2}\langle\phi^2\rangle$

Axion relic density

★ Two regimes for $\ddot{\phi} + 3H(t)\dot{\phi} + m^2\phi = 0$

★ $m \ll 3H$ early times $t < t_1$ where $3H(t_1) = m$

$$\phi = cte$$

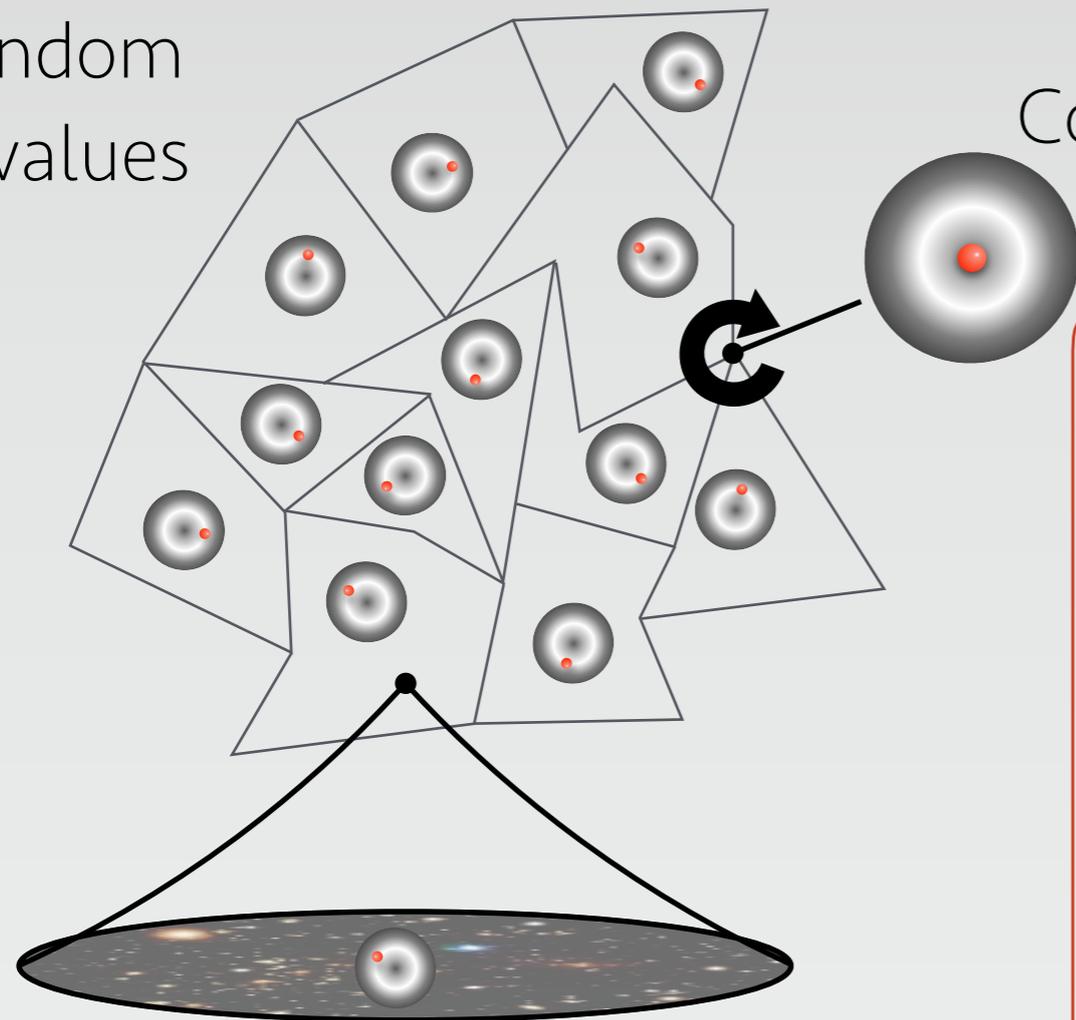
★ $m \gg 3H$ late times $t > t_1$

$$\phi(t) = \phi_1 \sqrt{\frac{a_1^3}{a(t)^3}} \cos mt$$

$$\rho = \rho_i \left(\frac{a_{\text{osc}}}{a} \right)^3 \text{ as expected for dark matter}$$

Two Classes of Scenarios

Random
 ϕ_i values



Cosmic axion string

Post-inflation PQ breaking
topological defects decay later

our Hubble volume
contains average over many ϕ_i

$$\Omega_\phi h^2 \sim 8 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{1.19} \times \left(\frac{g_{*,1}}{70} \right)^{-0.41} \left(\frac{\Lambda_{QCD}}{400 \text{ MeV}} \right)$$

Pre-inflation PQ breaking

one ϕ_i singled out randomly
same initial conditions everywhere

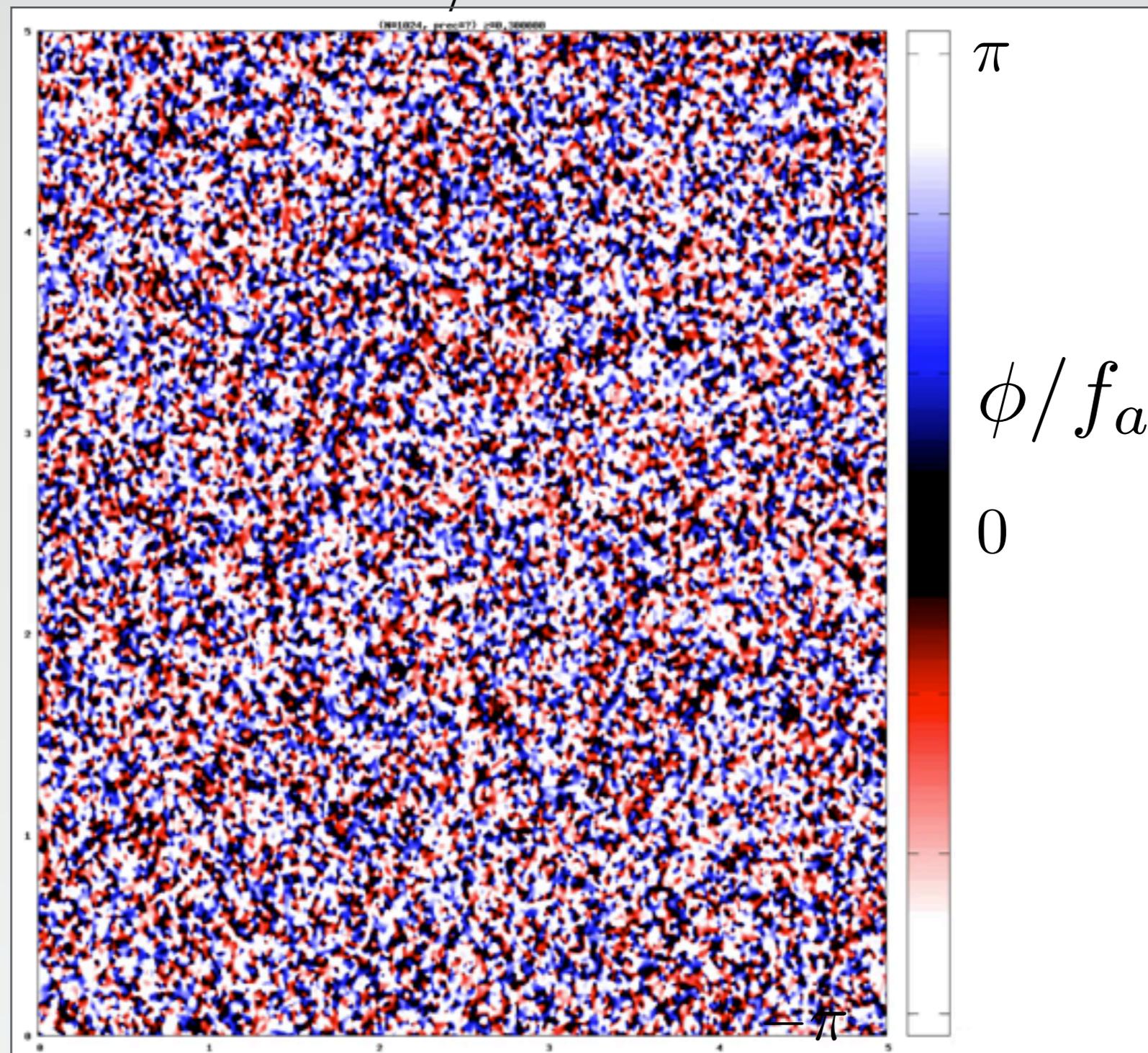
$$\Omega_\phi h^2 = 0.11 \phi_i^2 \left(\frac{10 \mu\text{eV}}{m} \right)^{1.184}$$

Axion Production from Cosmic Strings

~ 10 comoving
solar systems

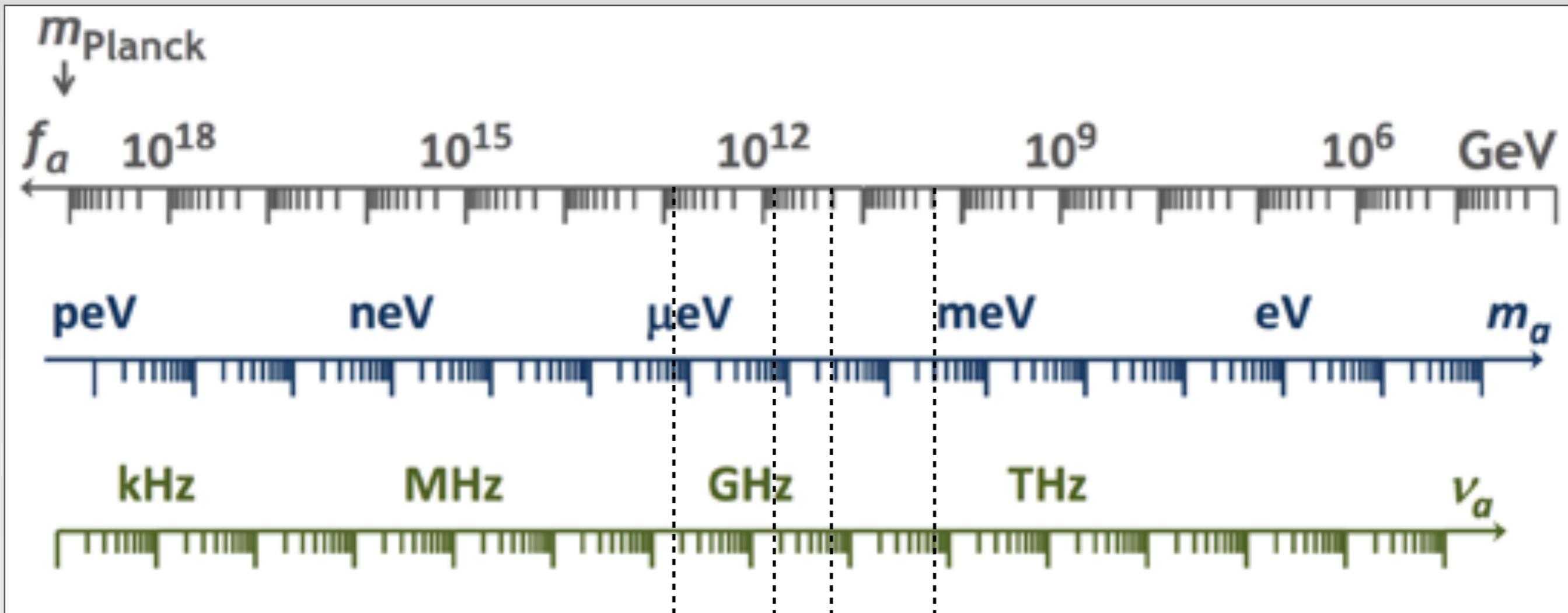
$T_0 \sim 4 \text{ GeV}$

fraction of a second



courtesy J. Redondo

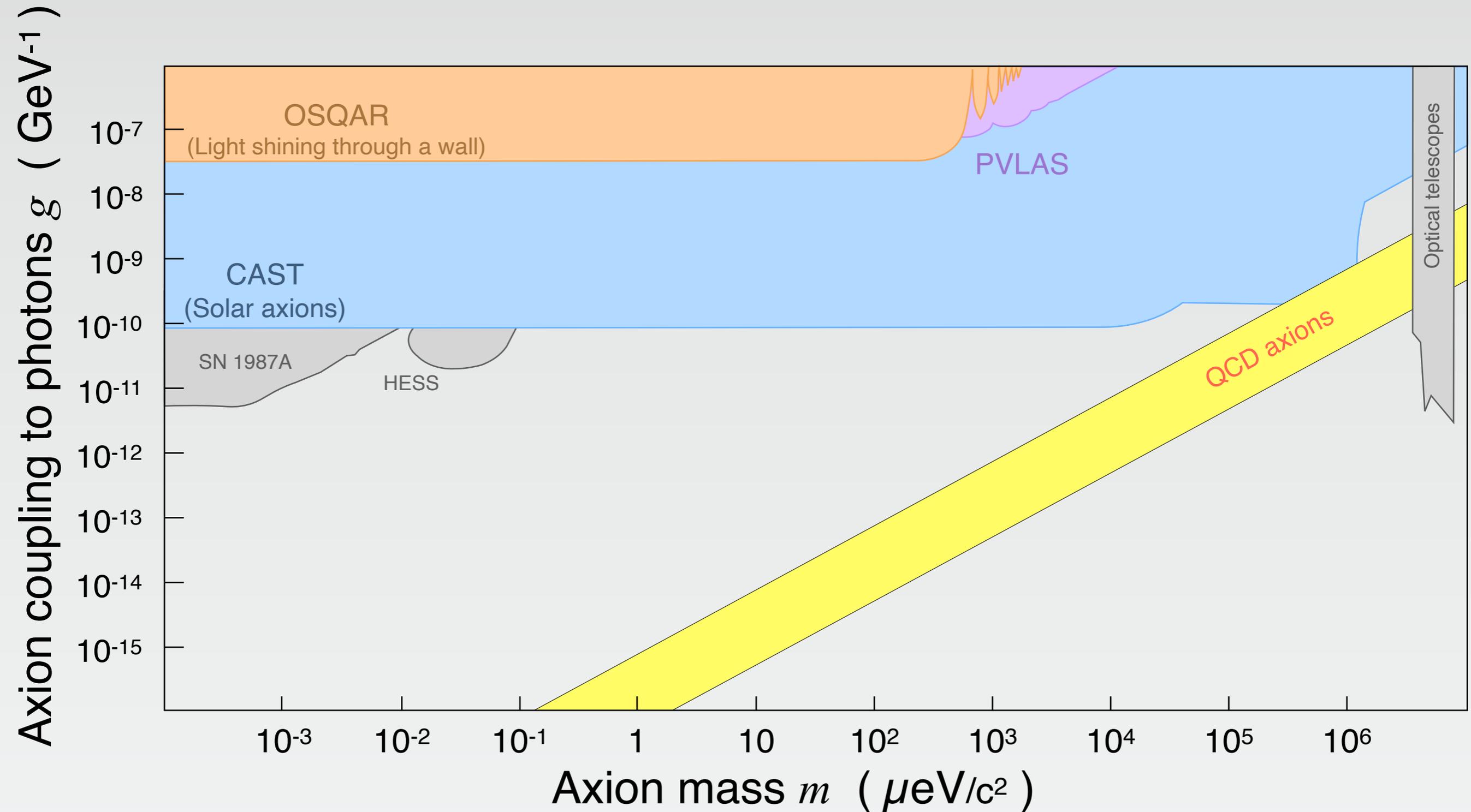
Relevant masses & frequencies



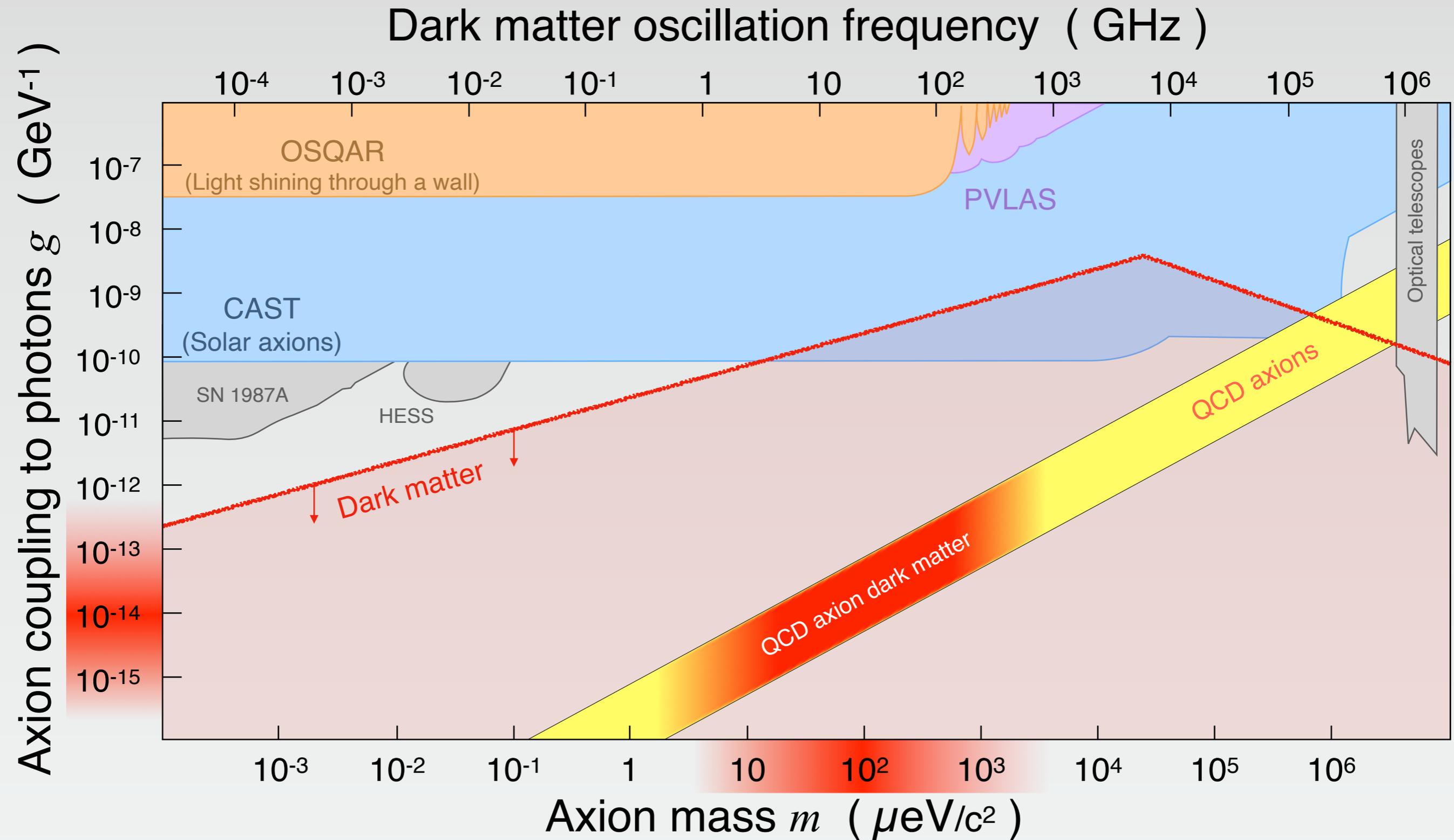
Pre-inflation
PQ breaking

Post-inflation
PQ breaking

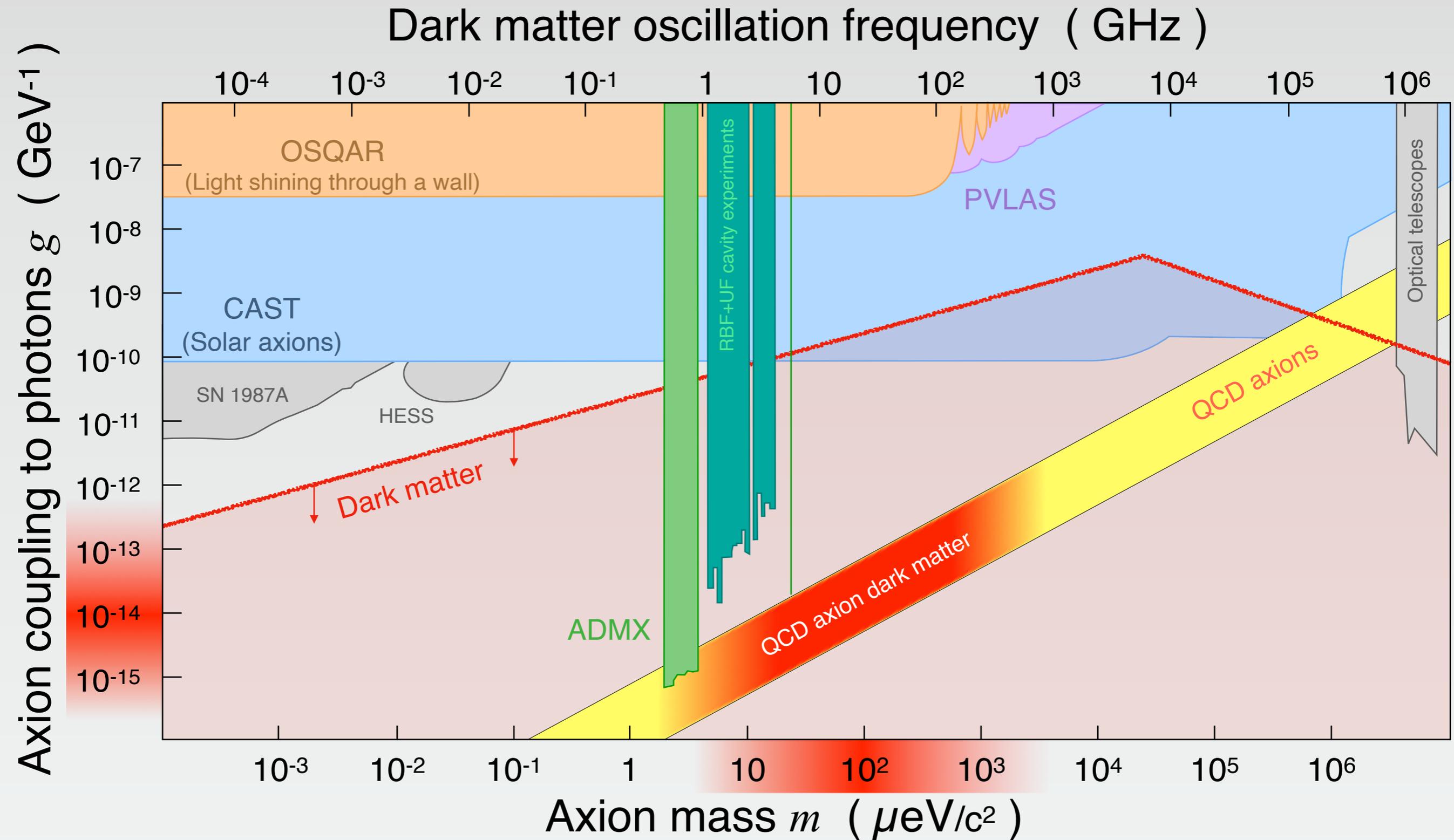
Axion searches and constraints



Axion searches and constraints



Axion searches and constraints



Axions as dark matter

Phase space density is huge...

$$N_{\text{particles}} \simeq \frac{10^{12} M_{\odot}}{m} \simeq 10^{83} \frac{10 \mu\text{eV}/c^2}{m}$$

$$N_{\text{cells}} \simeq \frac{\frac{4\pi}{3} p_{\text{max}}^3 \times \frac{4\pi}{3} R^3}{(2\pi\hbar)^3} \simeq 2 \times 10^{59} \times \left(\frac{m}{10 \mu\text{eV}/c^2} \right)^3$$

$p_{\text{max}} = mv_{\text{escape}} = m \times 550 \text{ km/s}$
 $R \sim 50 \text{ kpc}$

$$\frac{N_{\text{particles}}}{N_{\text{cells}}} \simeq 5 \times 10^{23} \times \left(\frac{10 \mu\text{eV}/c^2}{m} \right)^4$$

Local Relic Axions

★ Axion field oscillates on scale $>$ solar system

★ Velocity of the lab is $< 300 \text{ km/s} = 10^{-3} c$

$$v = \frac{kc^2}{\omega}, \quad \omega = \frac{mc^2}{\hbar} \left(1 + \frac{v^2}{2c^2} + o\left(\frac{v^2}{c^2}\right) \right)$$

★ Oscillation frequency set by axion mass

$$k = 0 \Rightarrow \omega = \frac{mc^2}{\hbar} \quad f = \frac{mc^2}{2\pi\hbar} \quad f = 2.4 \text{ GHz} \times \frac{m}{10 \mu\text{eV}}$$

★ Known signal dispersion $\frac{\delta f}{f} = \frac{v^2}{2c^2} < 10^{-6}$

Classical description

$$\mathcal{L}_{\text{a}\gamma} = g_{\text{a}\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} \propto g_{\text{a}\gamma} \vec{E} \cdot \vec{B} \quad \Rightarrow \text{modified Maxwell equations}$$

$$\vec{\nabla} \cdot \vec{B} = \vec{0} \quad \vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \left[-g_{\text{a}\gamma} c \vec{\nabla} \phi \cdot \vec{B} \right]$$

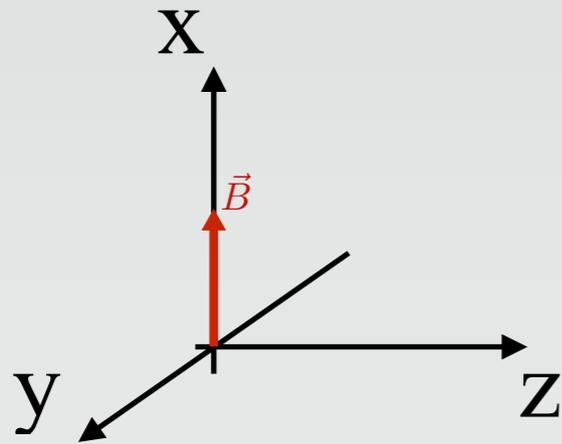
$$\vec{\nabla} \wedge \vec{E} = -\dot{\vec{B}} \quad \vec{\nabla} \wedge \vec{B} = \mu_0 \vec{j} + \frac{1}{c^2} \dot{\vec{E}} \left[+ \frac{g_{\text{a}\gamma}}{c} \left(\vec{\nabla} \phi \wedge \vec{E} + \dot{\phi} \vec{B} \right) \right]$$

In an external magnetic field \vec{B}_{ext} with $\phi(t)$ as for dark matter:

$$\vec{E}_{\text{DM}} = -\phi_0 g_{\text{a}\gamma} c \vec{B}_{\text{ext}} \cos \left(\frac{m_a c^2}{\hbar} t \right)$$

Equations of Motion

- ★ Consider a plane wave



$$\Psi = \begin{pmatrix} \vec{A}_{||} \\ \phi \end{pmatrix} \exp(-i(\omega t - kz))$$

- ★ Leads to the equation of motion :

$$\left[(\omega^2 - k^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} 0 & -g|\vec{B}|\omega \\ -g|\vec{B}|\omega & m \end{pmatrix} \right] \begin{pmatrix} \vec{A}_{||} \\ \phi \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Photon and axion mix, angle prop. to B field

Solutions of EoM

★ First line of EoM :

$$(\omega^2 - k^2)A_i + gB_i\omega\phi = 0$$

$$\Rightarrow \vec{A}_{DM} = -\phi_0 \frac{gB_x\omega}{\omega^2 - k^2} e^{-i(\omega t - kz)} \vec{e}_x$$

★ To ϕ is associated a small electric field :

$$\vec{E}_{DM} = -\partial_0 \vec{A}_{DM} = i\phi_0 g B_x e^{-i(\omega t - kz)} \vec{e}_x$$

Dark Matter Solution of EoM

★ DM velocity is small $v = \frac{kc^2}{\omega} \ll c$

★ For $v = 10^{-3}c$, $\lambda_{DM} \simeq 30$ m

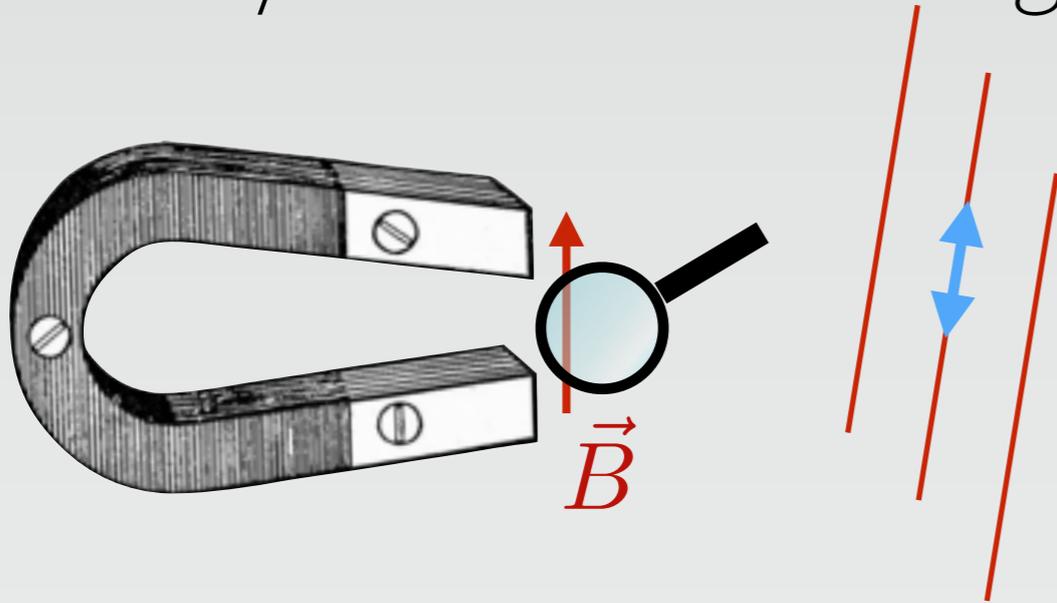
★ Practically a coherent oscillation :

$$\vec{E}_{DM} = i\phi_0 g B_x e^{-i\omega t} \vec{e}_x$$

up to $\mathcal{O}(10^{-3})$ terms

More Concretely...

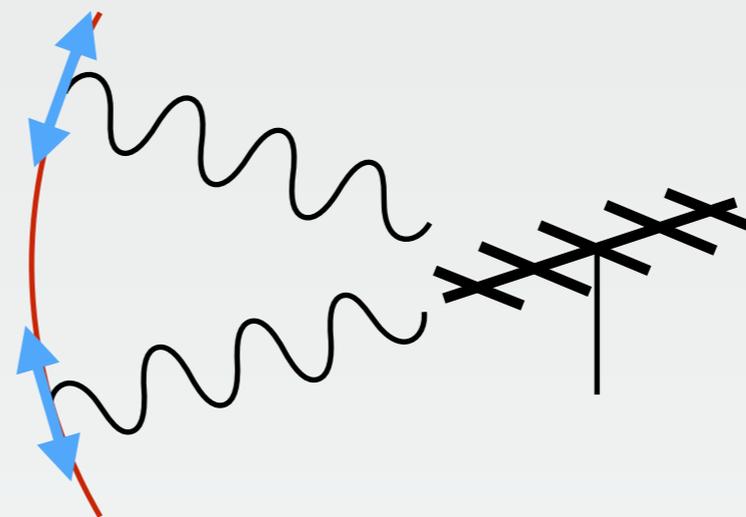
Everywhere in the background : $\phi = \phi_0 \exp(-i\omega t)$



$$g F_{\mu\nu} \tilde{F}^{\mu\nu} \phi \propto g \vec{E} \cdot \vec{B} \phi$$

small oscillating electric field along each B-field line

Possible lab experiment :
bend field lines,
detect EM radiation from this

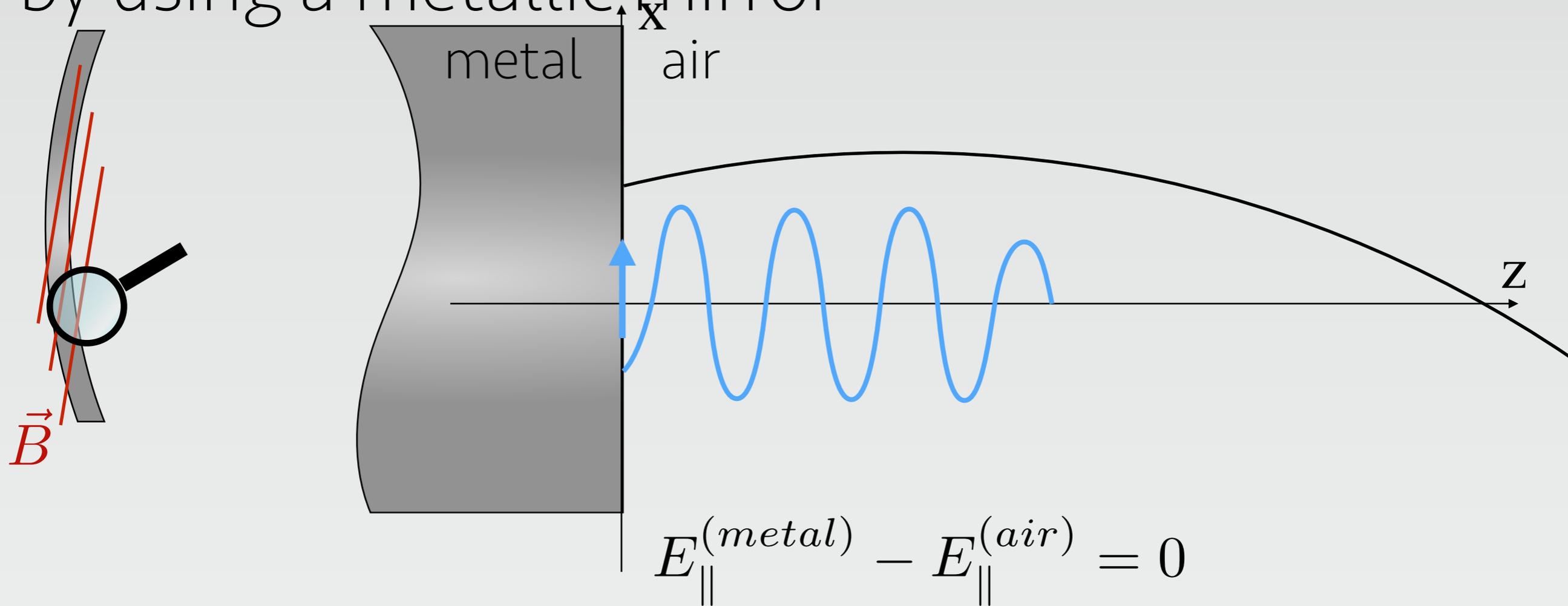


~10 GHz
signal

Polarized

Imposing Boundary Conditions

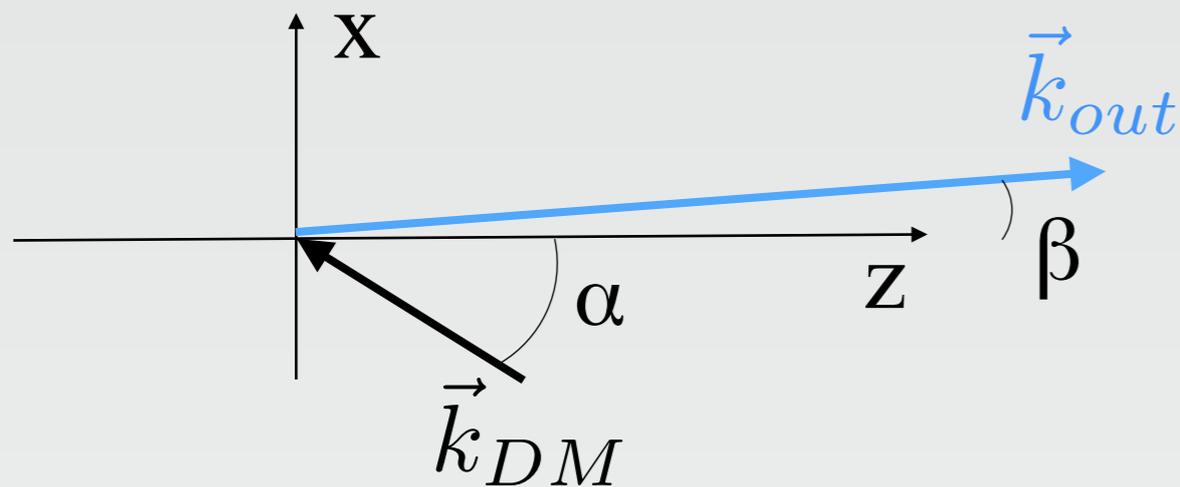
.. by using a metallic mirror



$$\vec{E}_{tot} = \underbrace{i\phi_0 g B_x e^{-i(\omega t)}}_{\vec{E}_{DM}} \vec{e}_x - \underbrace{i\phi_0 g B_x e^{-i(\omega t - kz)}}_{\vec{E}_{out}} \vec{e}_x$$

Experiment Principle

- ★ Radio emission is normal to the surface
- ★ Another derivation show



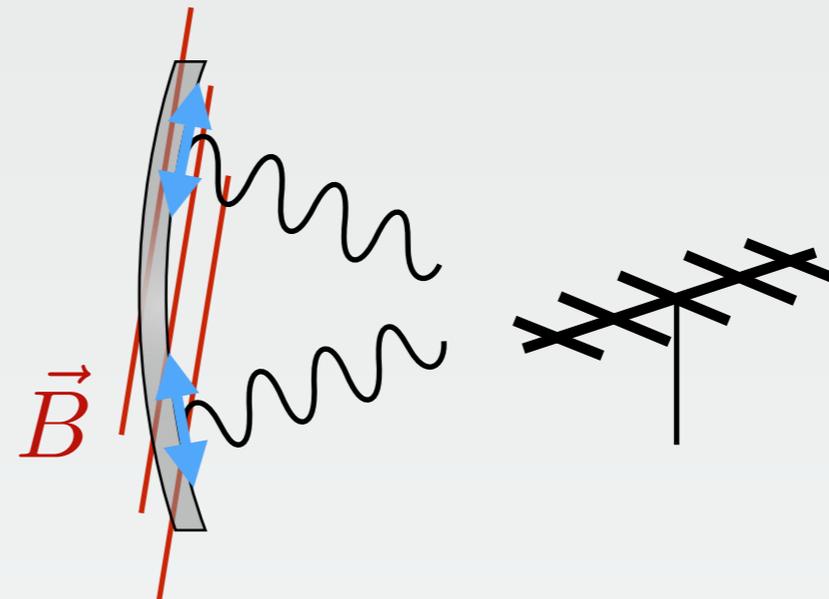
$$\sin(\beta) = \sin(\alpha) \frac{v}{c\sqrt{1 - \left(\frac{v}{c}\right)^2}}$$

- ★ Proposed experiment

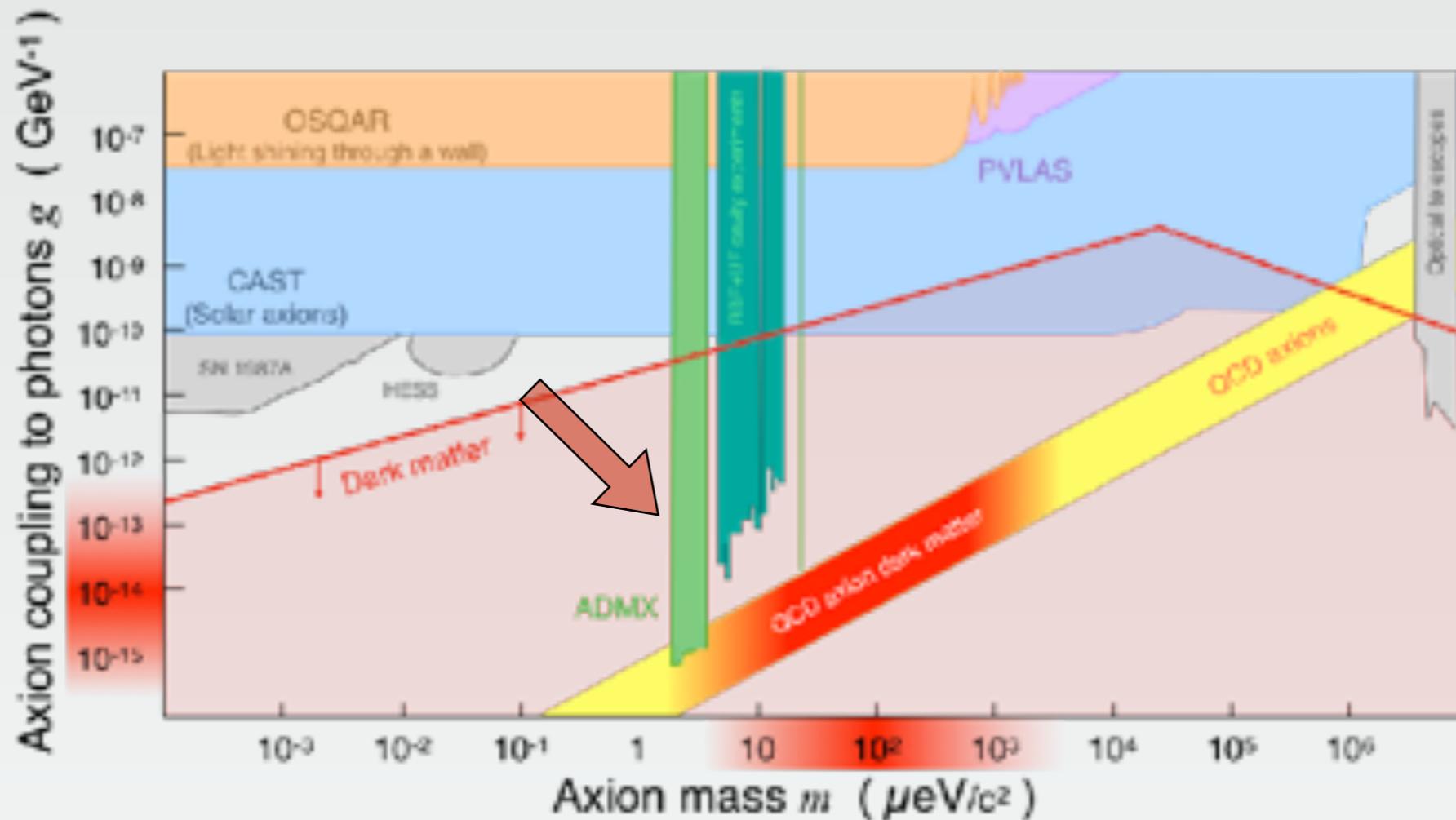
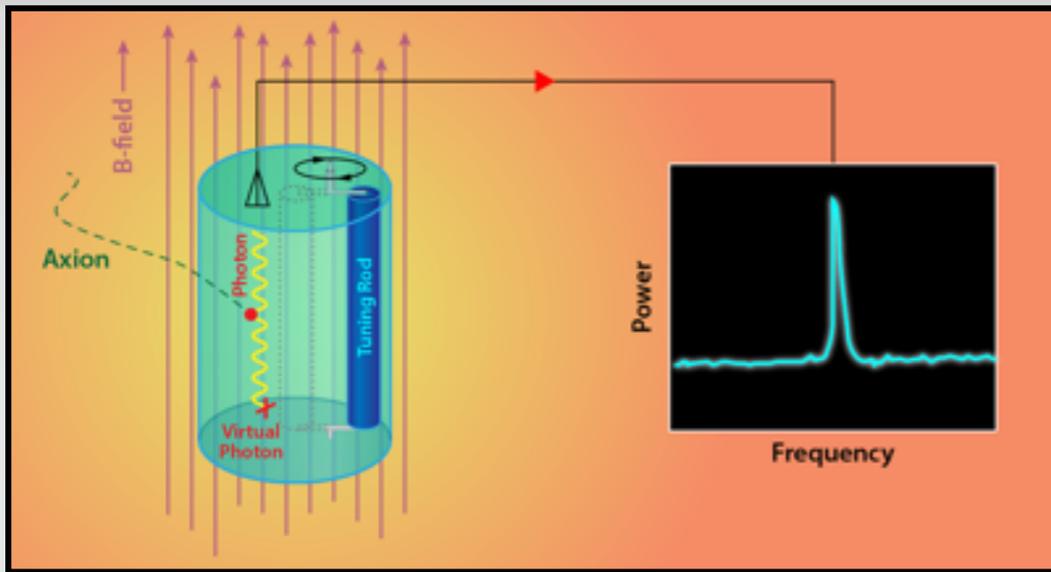
D. Horns et al. JCAP 2012

J. Jaeckel & J. Redondo, JCAP 2013

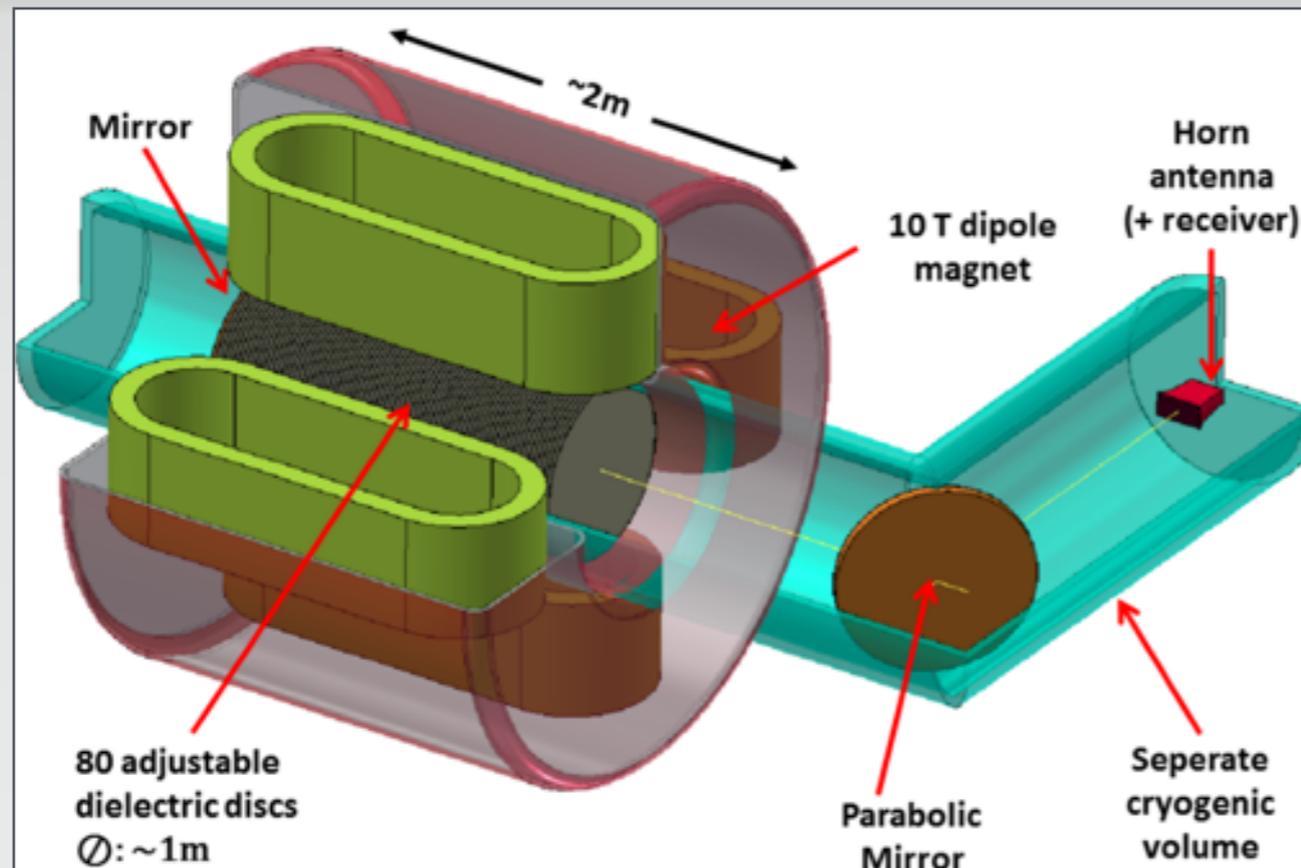
J. Jaeckel & S. Knirck, Patras 2016



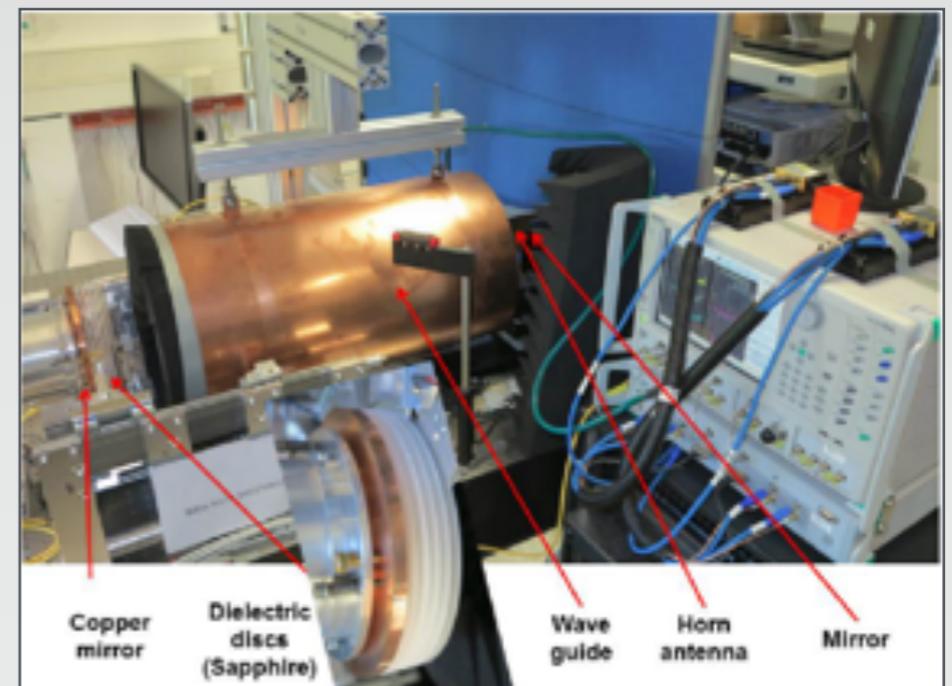
Narrowband experiments: ADMX



Narrowband experiments: MADMAX



Prototype at MPI Munich:



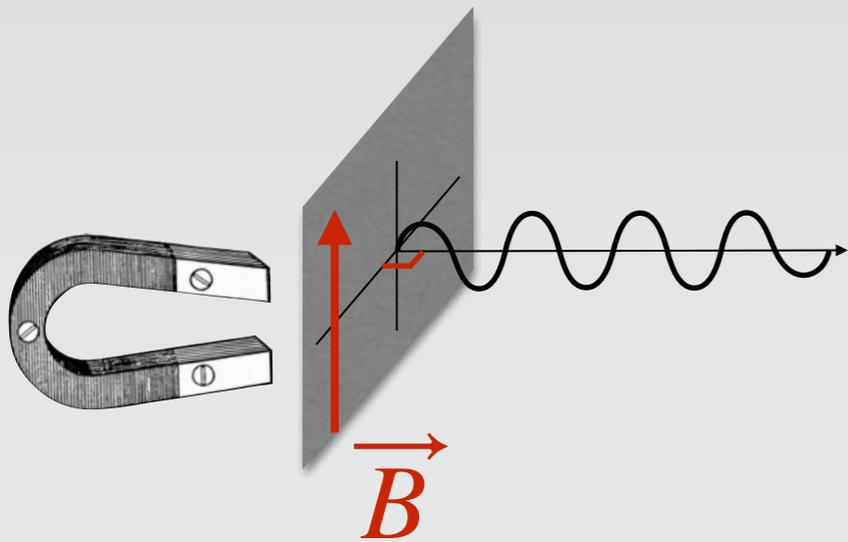
Very expensive magnet, long to produce
Challenging mechanics, requires lots of development

Do not work for $m_a > 200 \mu\text{eV}$ (i.e. $\nu > 50 \text{ GHz}$)

Possible synergies in the detection strategies

Proposed experiment: principle

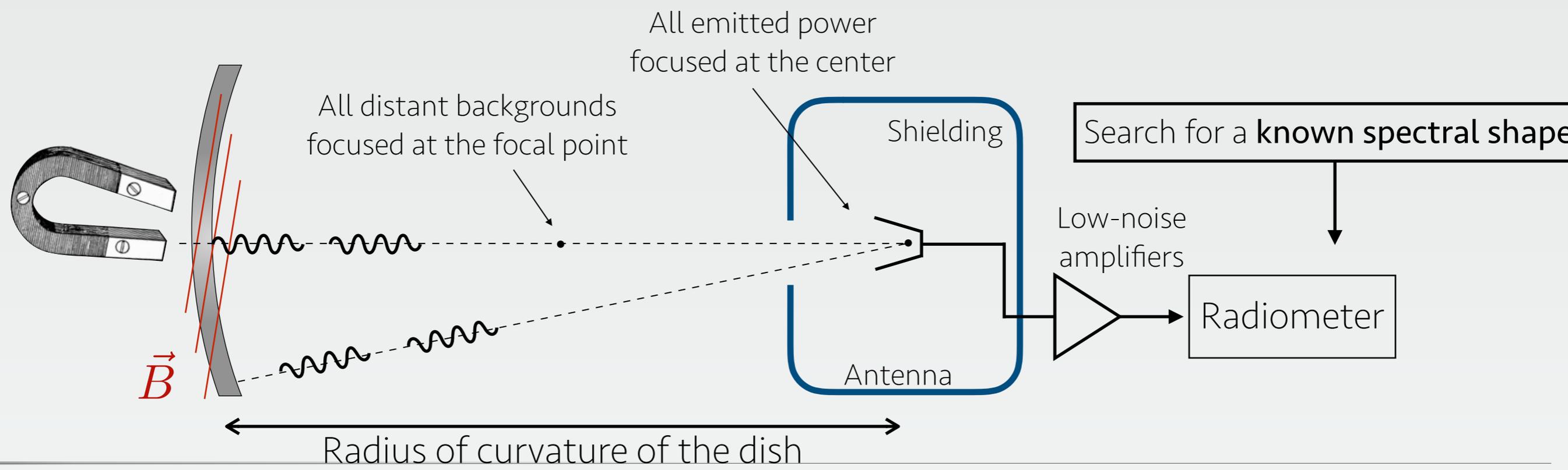
Axion dark matter \Rightarrow magnetized conductor emits continuous microwaves



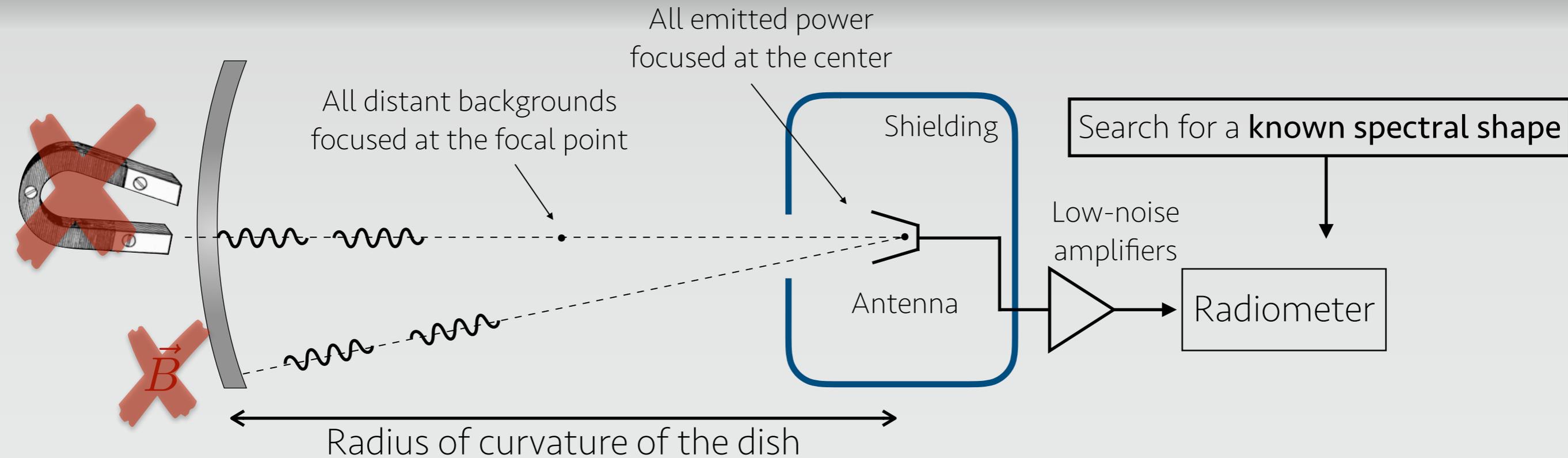
$$f = 2.4 \text{ GHz} \times \frac{m c^2}{10 \mu\text{eV}}$$

$$\frac{\delta f}{f} = 10^{-6}$$

Simplest possible experimental setup (Horns et al., JCAP 2014):



Pathfinder: SHUKET



Poorman's experiment: no magnetic field

Sensitive to other type of dark matter: « dark photons »

SHUKET

=



Search for U(1) dark matter with an Electromagnetic Telescope



Hidden photons

New vector field ϕ :

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}\phi_{\mu\nu}\phi^{\mu\nu} - \frac{m^2}{2}\phi_\mu\phi^\mu - \frac{\chi}{2}F_{\mu\nu}\phi^{\mu\nu}$$

Relatable to axion \downarrow coupling

$$-\frac{1}{4}gF_{\mu\nu}\tilde{F}^{\mu\nu}\phi \leftrightarrow -\frac{\chi}{2}F^{\mu\nu}X_{\mu\nu}$$

Same phenomenology (misalignment, oscillations, etc.) with

$$\vec{E}_{\text{DM}} = \chi\omega\vec{\phi}$$

χ could be 10^{-12} / 10^{-3}

K.R. Dienes *et al.*, Nucl Phys B 1997

M. Goodsell *et al.*, JHEP 2009

M. Goodsell *et al.*, JHEP 2012

Sensitivity to coupling parameters

★ Axion-like particles

$$\mathcal{L} = -\frac{1}{2}\partial_\mu\partial^\mu\phi + \frac{1}{2}m^2\phi^2 - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}gF_{\mu\nu}\tilde{F}^{\mu\nu}\phi$$

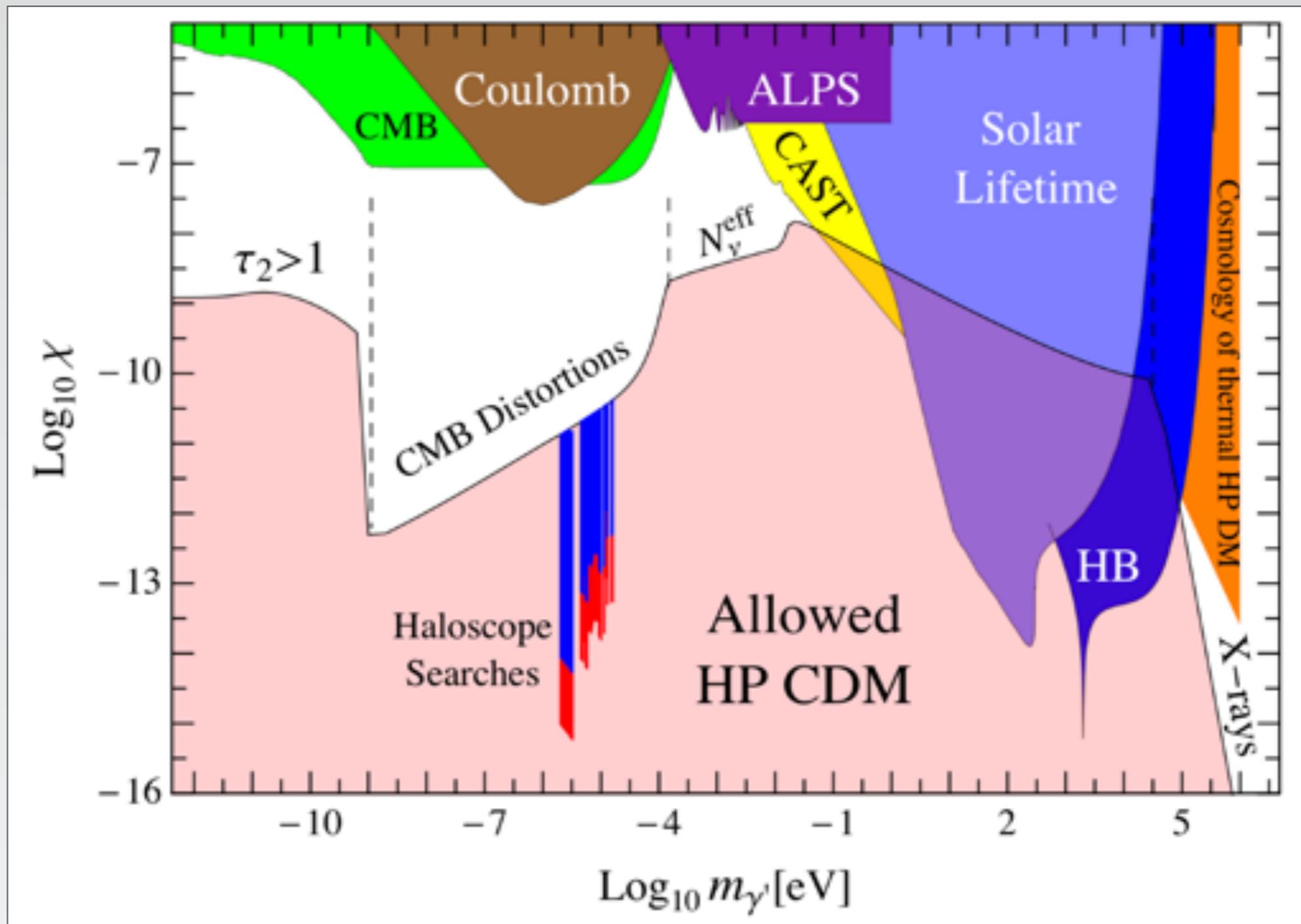
$$g = \frac{3.6 \times 10^{-8}}{\text{GeV}} \left(\frac{5 \text{ T}}{B}\right) \left(\frac{P_{det}}{10^{-23} \text{ W}}\right)^{1/2} \left(\frac{m}{\text{eV}}\right) \left(\frac{0.3 \text{ GeV/cm}^3}{\rho_\odot}\right)^{1/2} \left(\frac{1 \text{ m}^2}{A_{dish}}\right)^{1/2}$$

★ Hidden photons

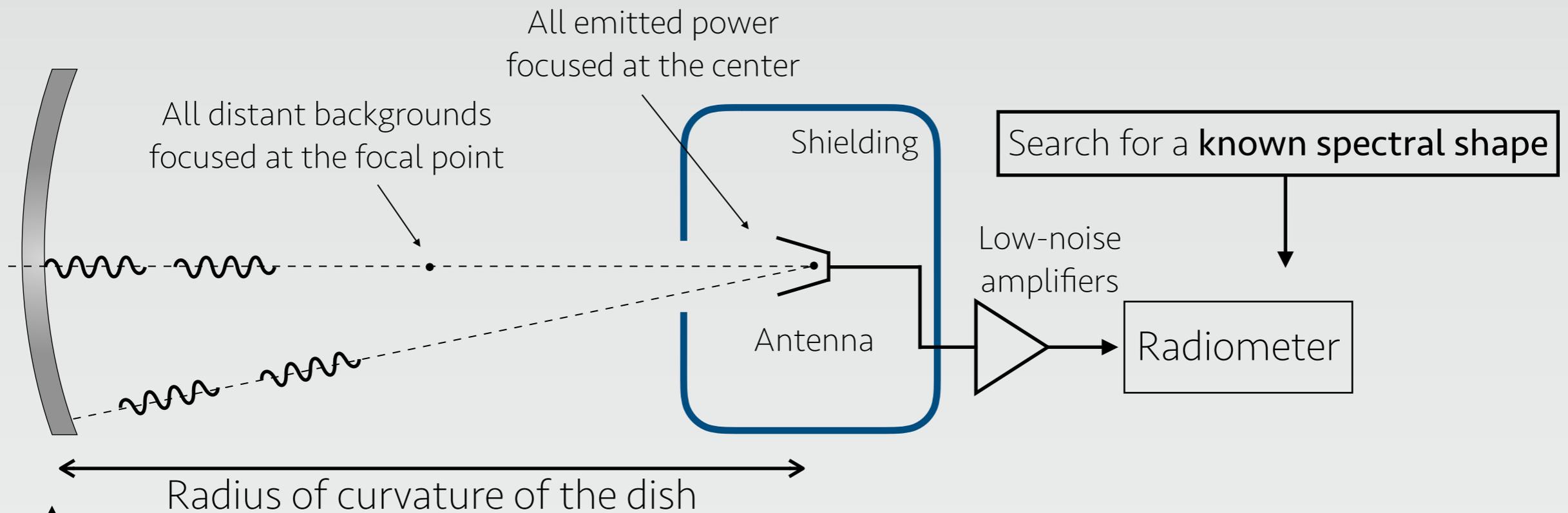
$$\mathcal{L} = -\frac{1}{4}\tilde{X}_{\mu\nu}\tilde{X}^{\mu\nu} + \frac{m^2}{2}\tilde{X}_\mu\tilde{X}^\mu - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{\chi}{2}F_{\mu\nu}\tilde{X}^{\mu\nu}$$


$$\chi = 4.5 \times 10^{-14} \left(\frac{P_{det}}{10^{-23} \text{ W}}\right)^{1/2} \left(\frac{0.3 \text{ GeV/cm}^3}{\rho_\odot}\right)^{1/2} \left(\frac{1 \text{ m}^2}{A_{dish}}\right)^{1/2} \left(\frac{\sqrt{2/3}}{\alpha}\right)$$

Constraints on hidden photons

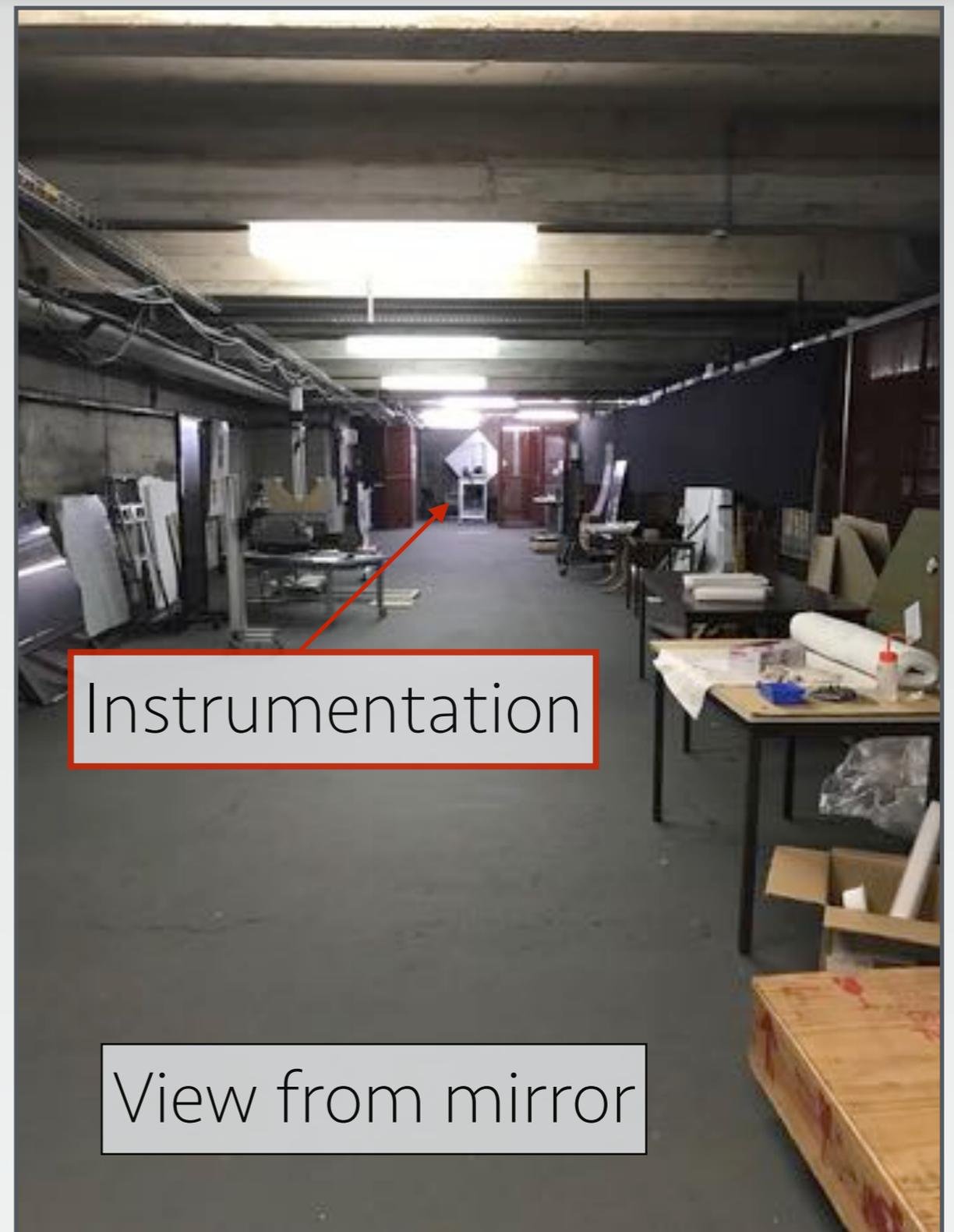
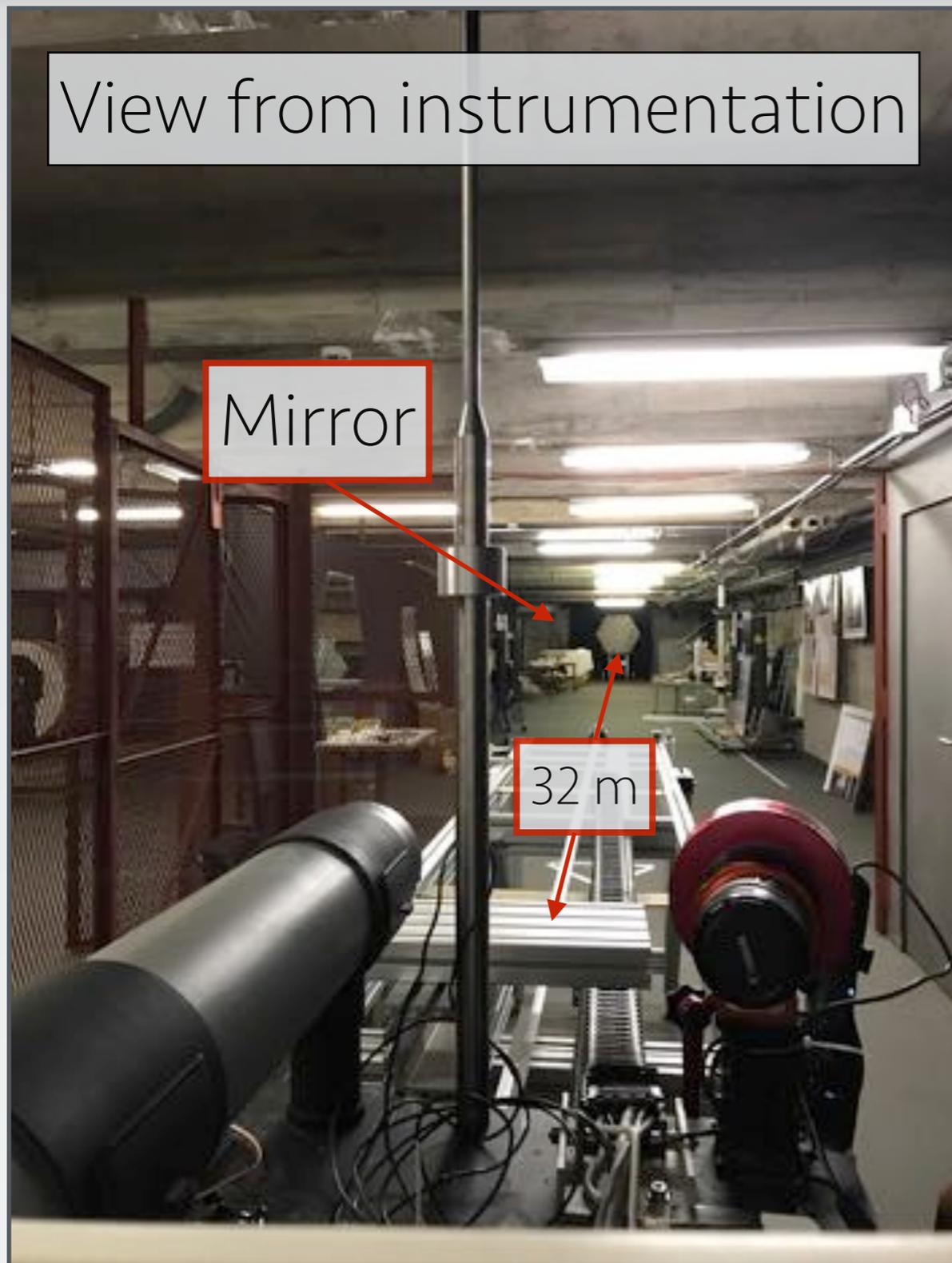


Interesting from model point of view, but no hint on the mass

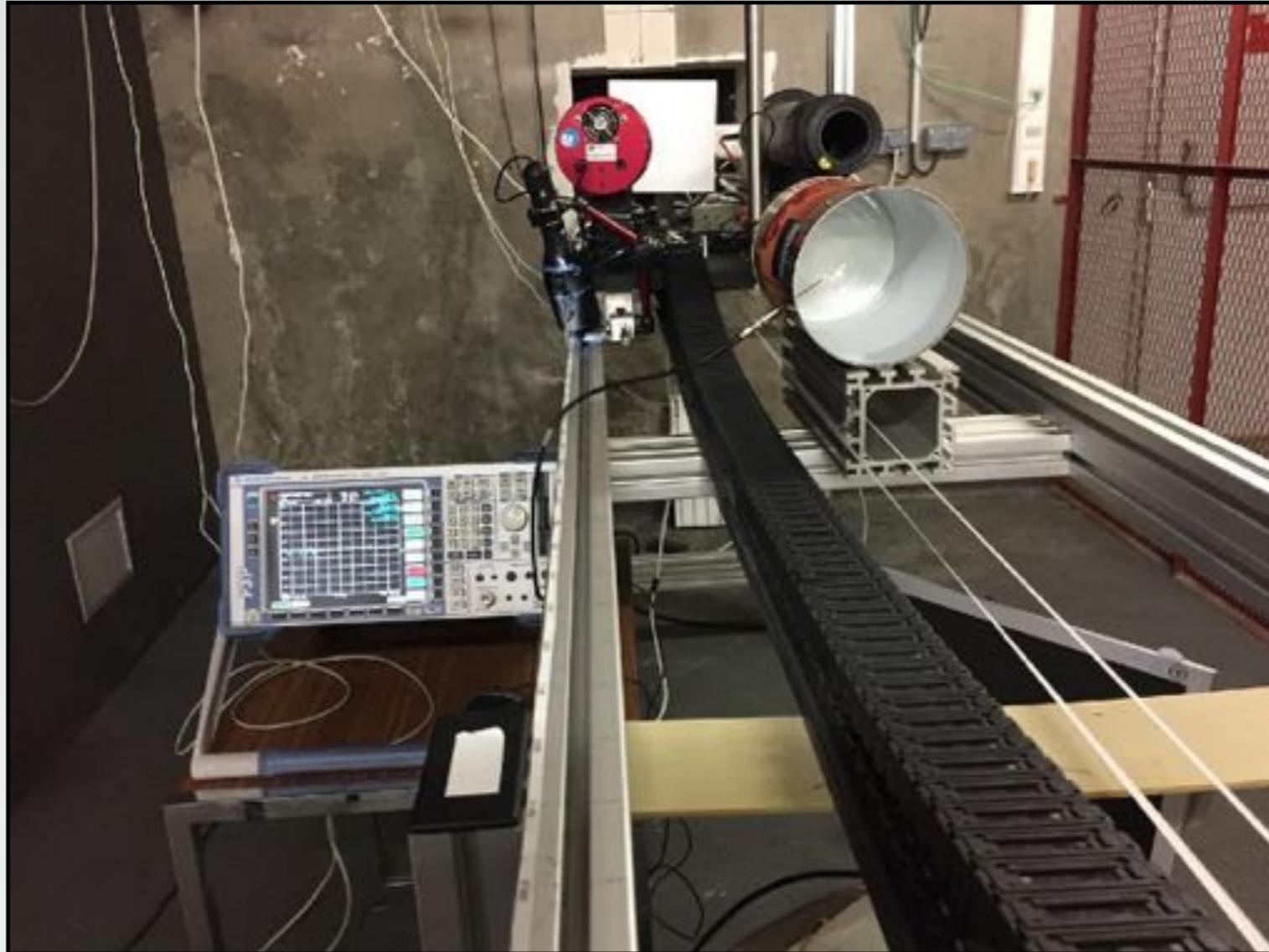


Mirror from a gamma-ray astronomy experiment (CTA)

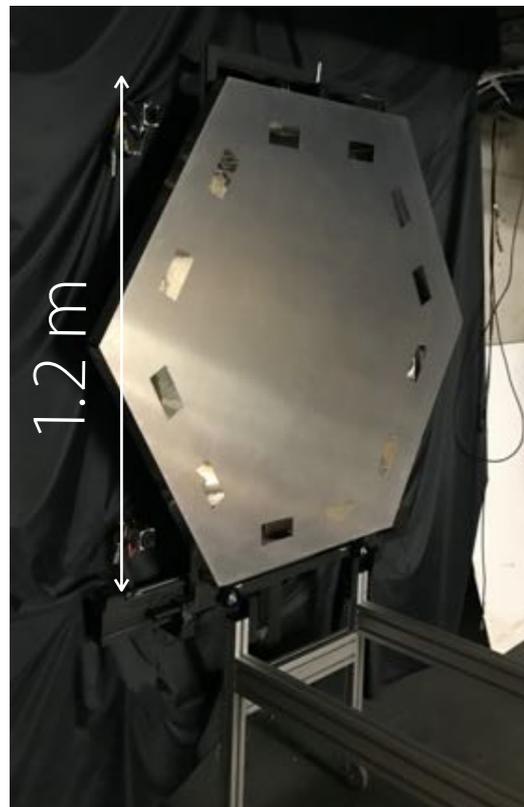
Optical Test Bench at Saclay



First setup



SHUKET improved setup



Mirror

32 m



Horn antenna + shielding

Amplifiers
+
Spectrum analyzer

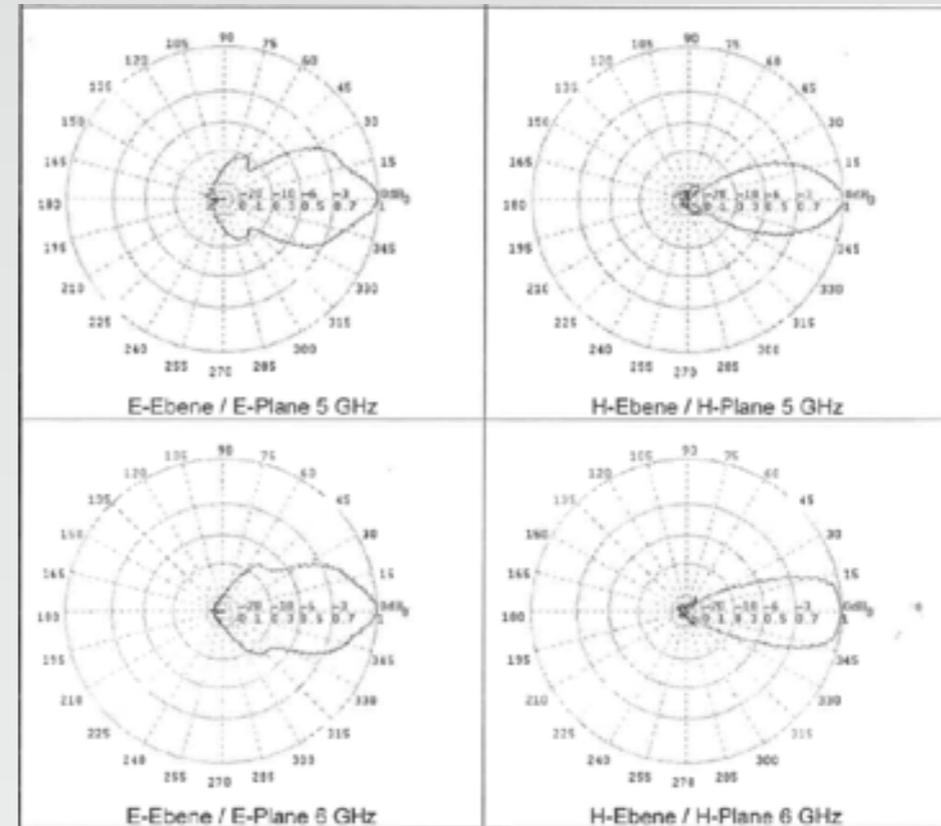


Achieved 10^{-22} W/Hz sensitivity
Constraints on hidden photons dark matter

Horn Antenna

Polarized horn antenna

1-18 GHz



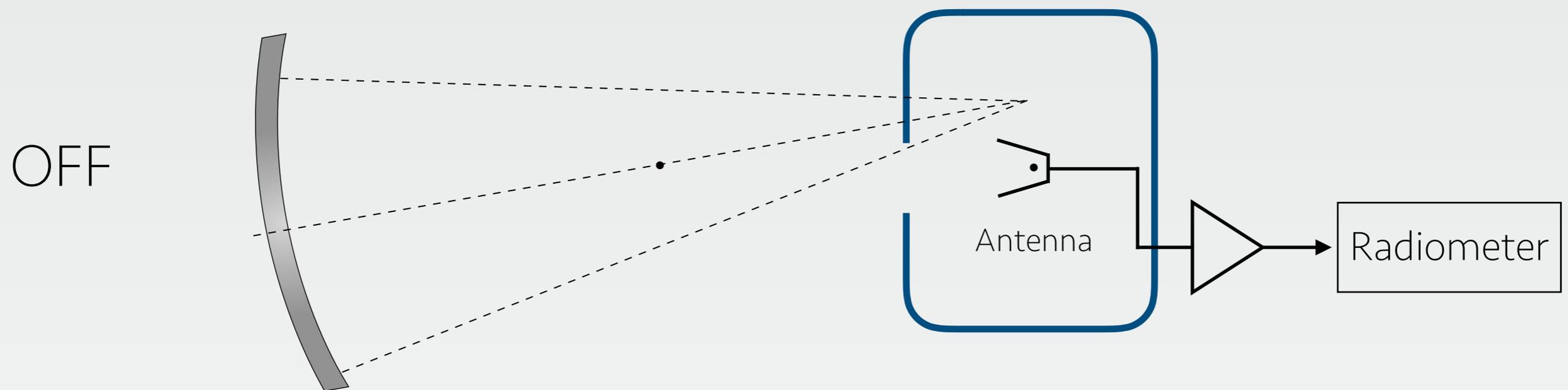
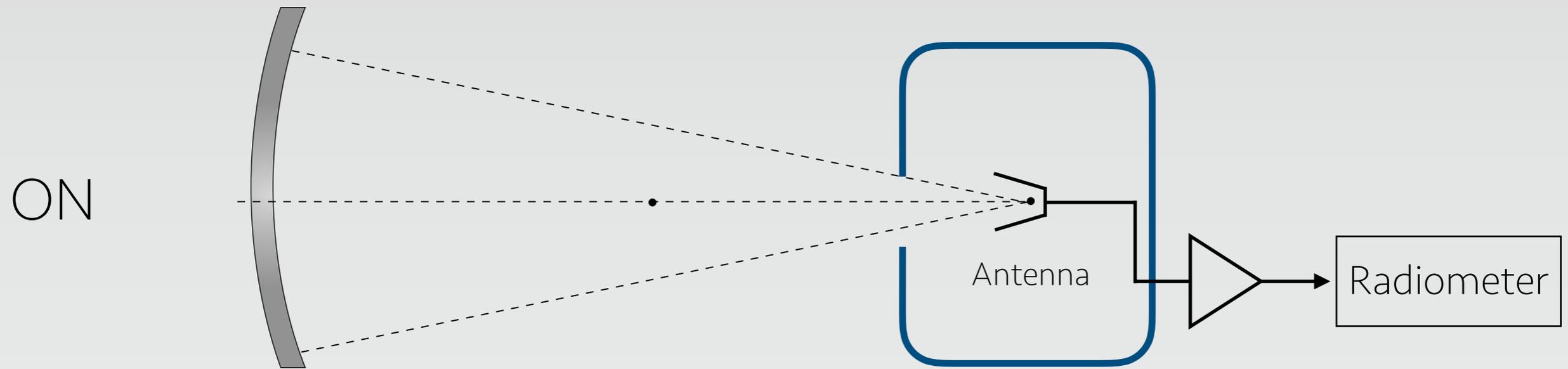
- 2017 data taking : 5 GHz - 7 GHz
- Incoming spherical wave, tiny solid angle

Spectrum analyzer

- ★ Room-temperature LN amplifier ~ 30 dB
- ★ Loan from Rhodes & Schwarz
 - 3 GHz - 20 GHz FSP
- ★ Performs FFT + power measurement in 1 Hz bins
- ★ 2 weeks : ~ 2 GHz bandwidth analyzed (5-7 GHz)

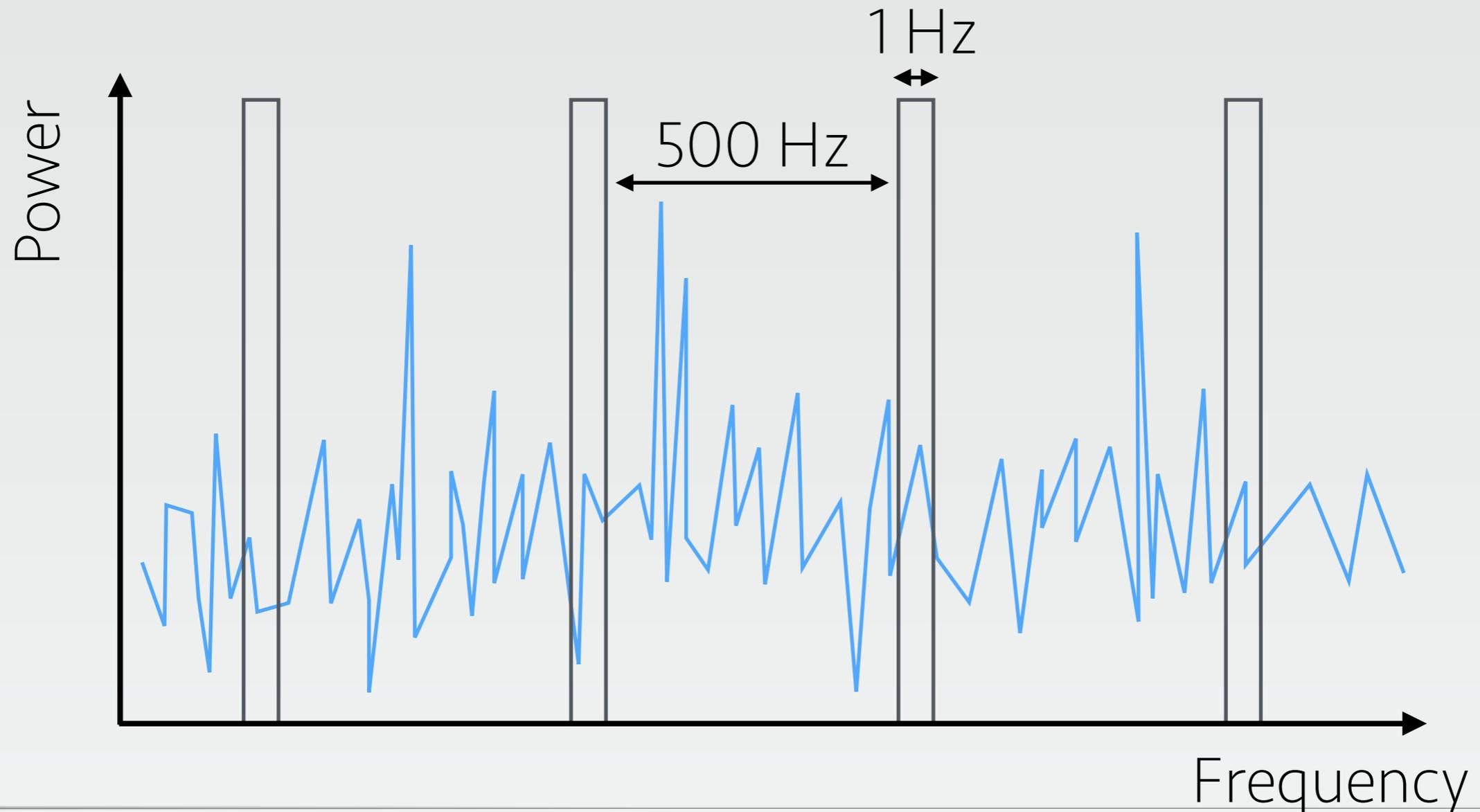


ON and OFF runs



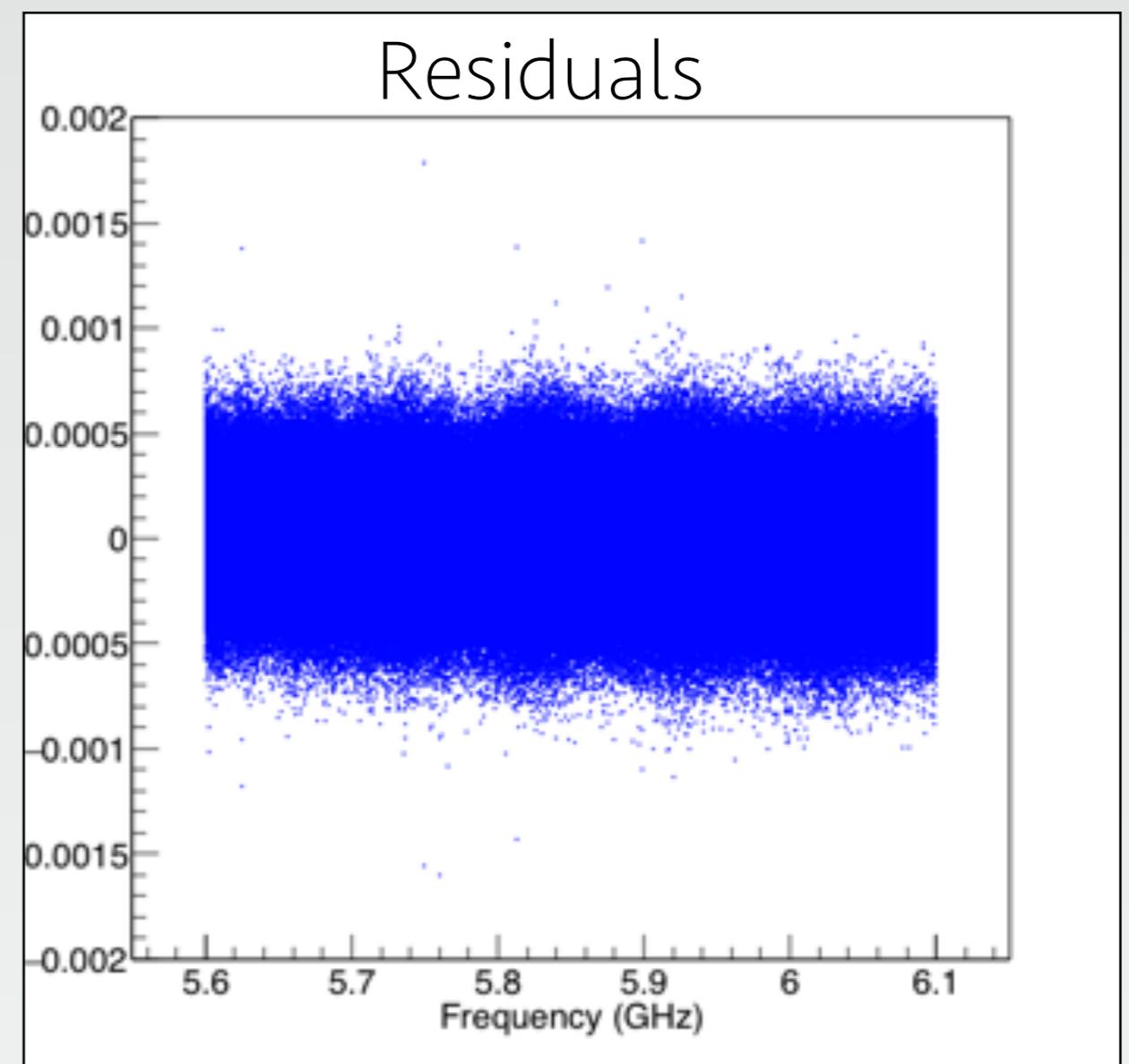
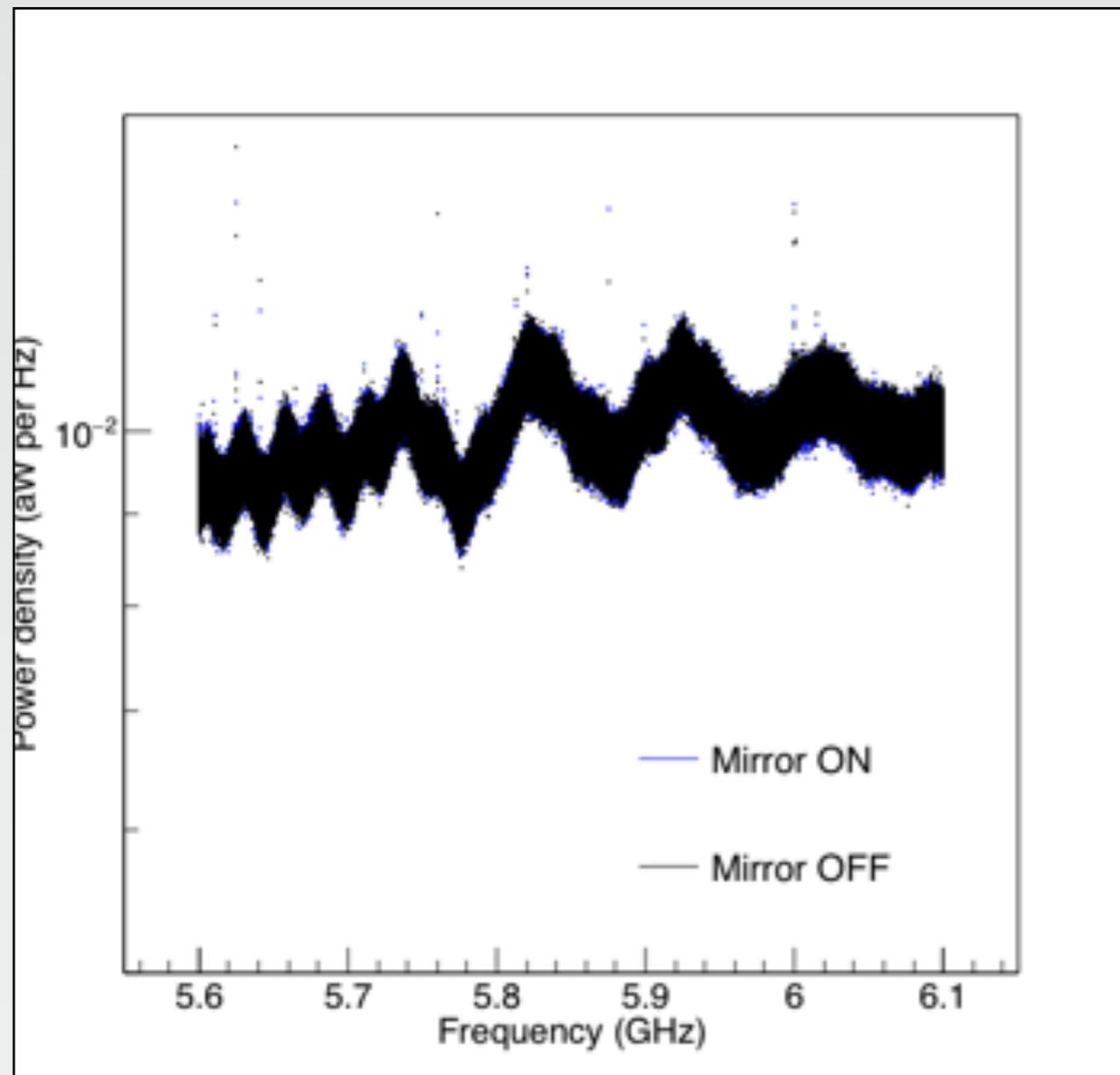
Search optimization

- ★ Signal is constant, broad band (~ 5 kHz)
- ★ Most backgrounds removed by sampling



Results

Power measurement w & w/o mirror



Expected sensitivity

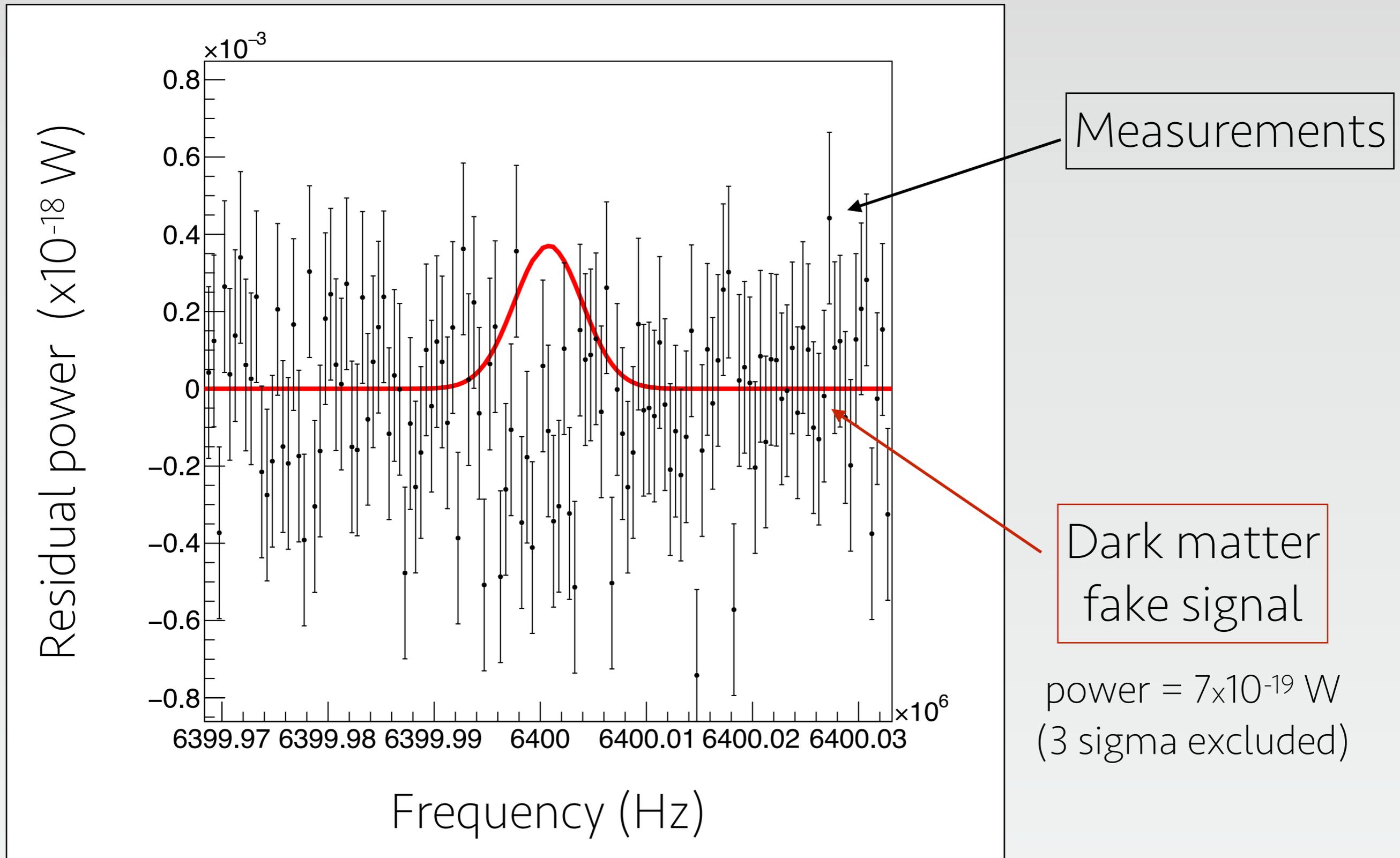
System noise temperature :

290 K environment

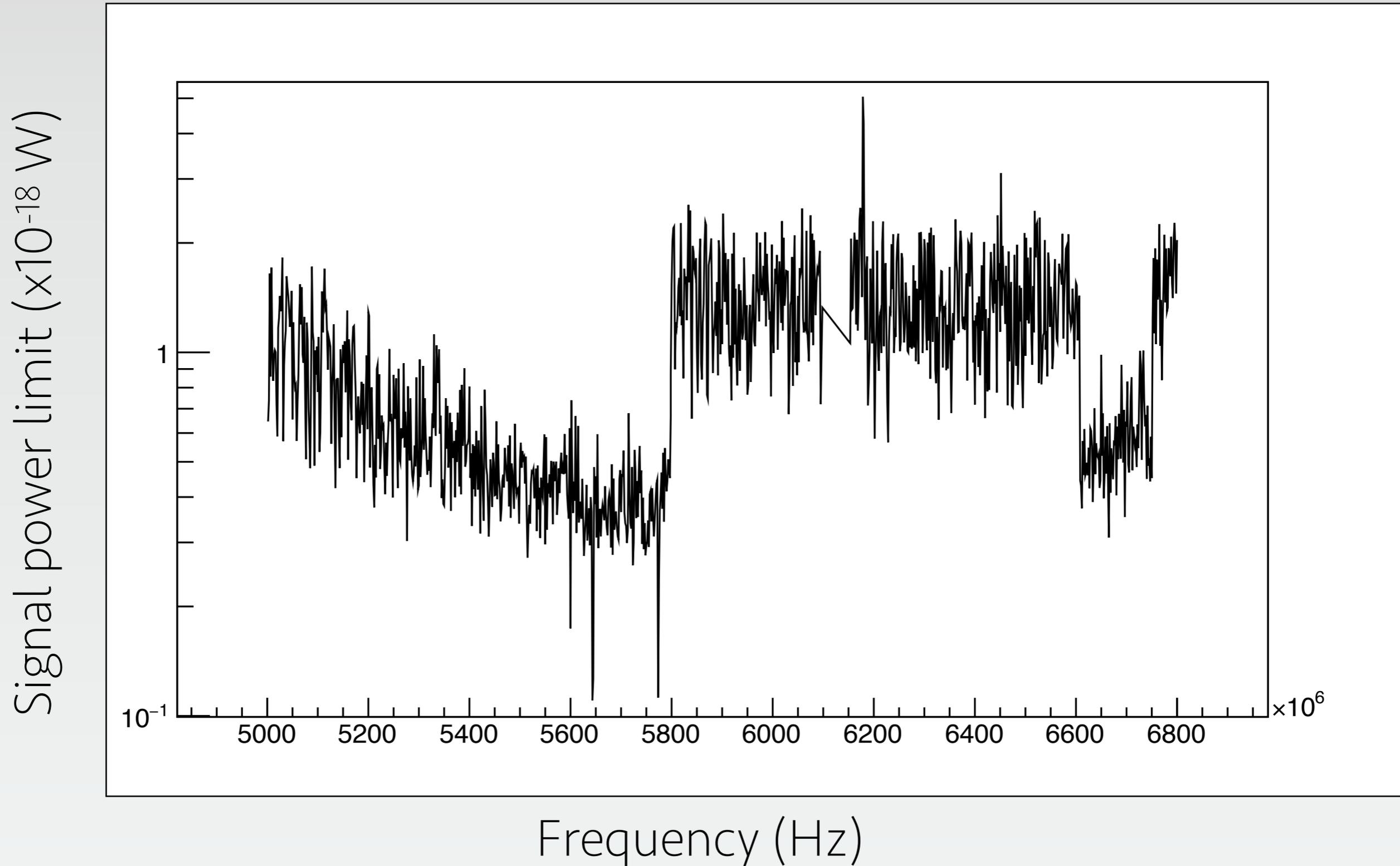
214 K LNA

50 K spectrum analyser

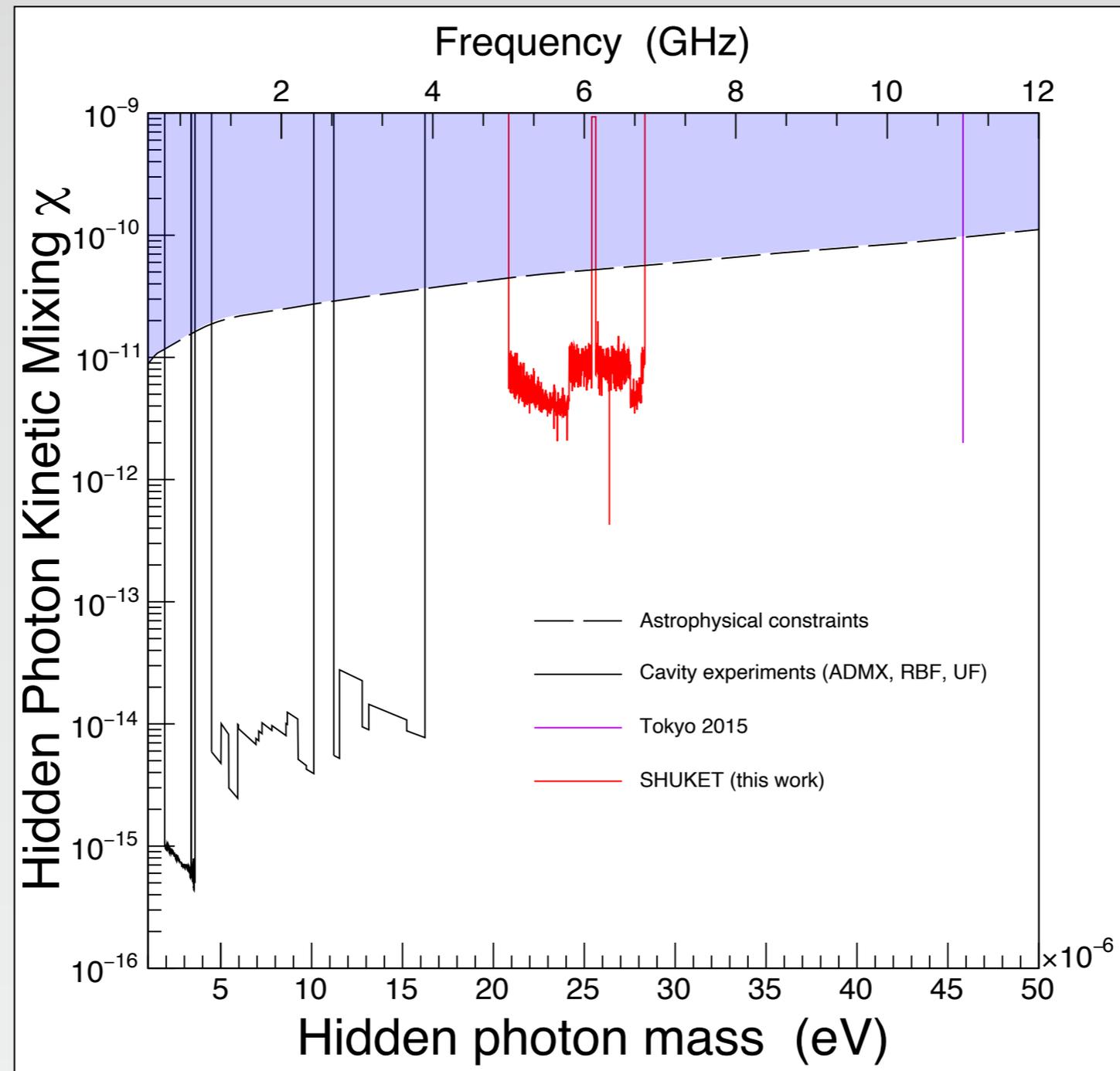
Limits on potential signals



Constraints on extra-power



Constraints on hidden photons mixing

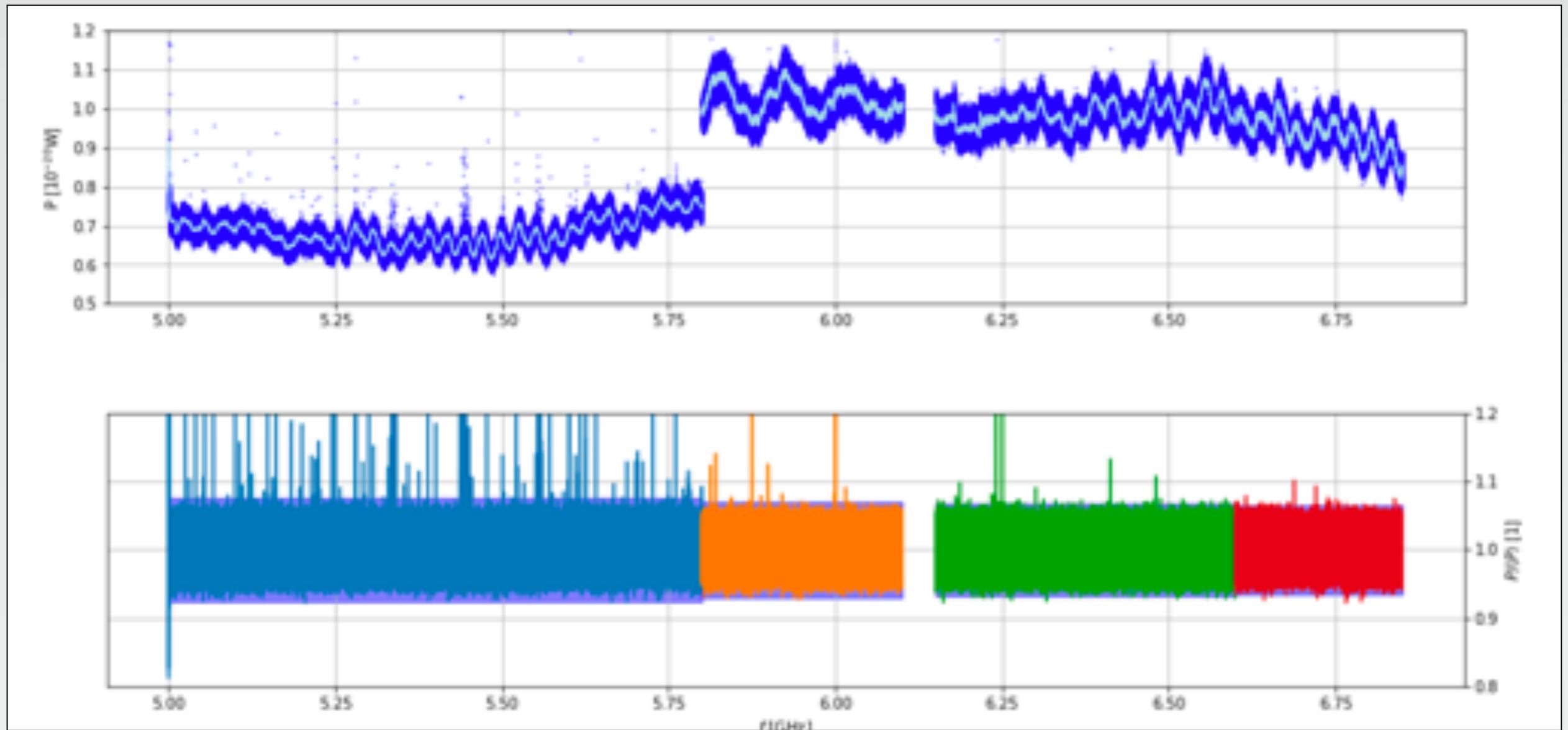


P. Brun, L. Chevalier, C. Flouzat, *Phys. Rev. Lett.* 122, 201801 (2019)

SYRTE re-analysis of Shuket

- ★ Use only ON data
- ★ Smooth-on data
- ★ Build a TS based on DM halo distribution as prior
- ★ χ is a built-in parameter

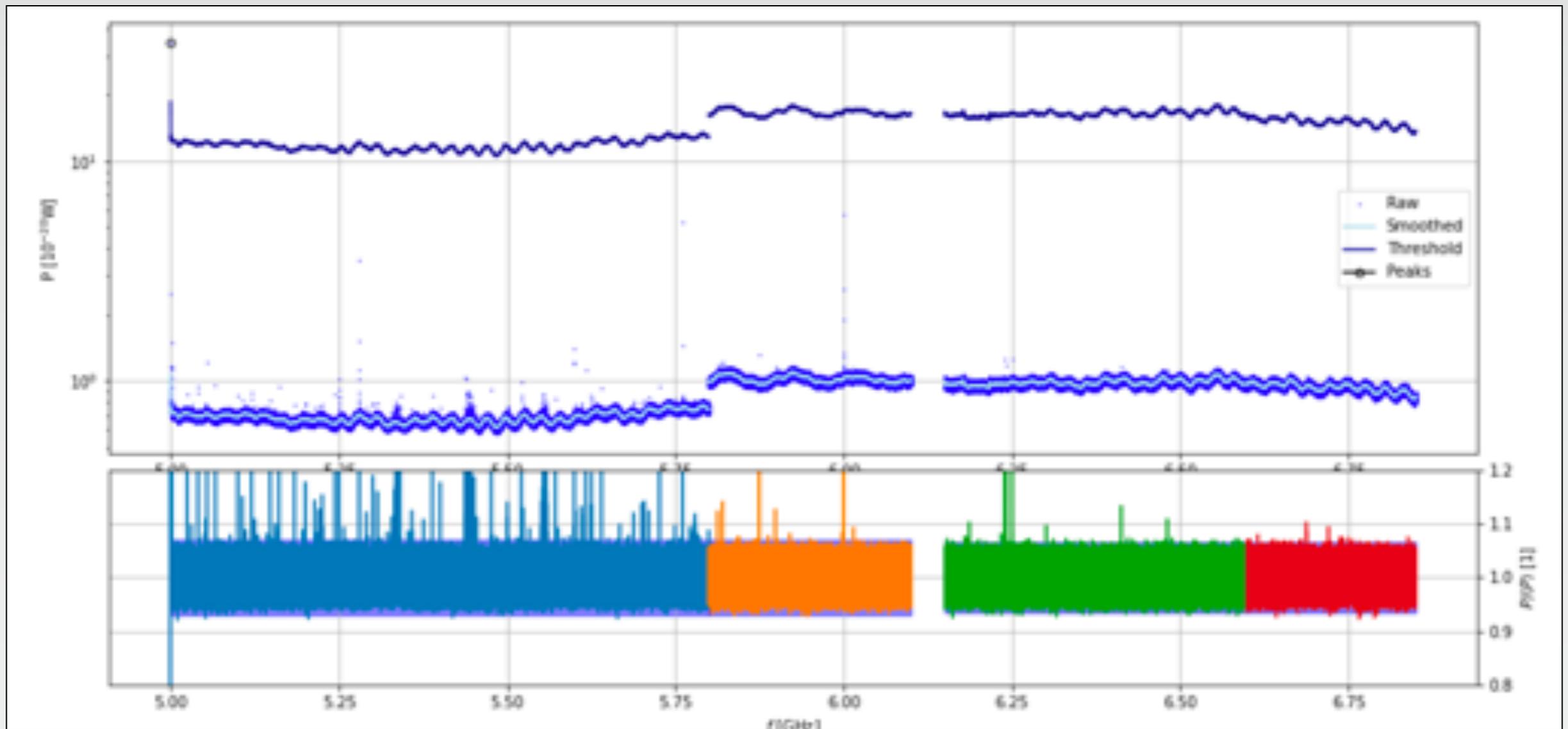
SYRTE re-analysis of Shuket



A. HEES, E. SAVALLE, P. WOLF

SYRTE re-analysis of Shuket

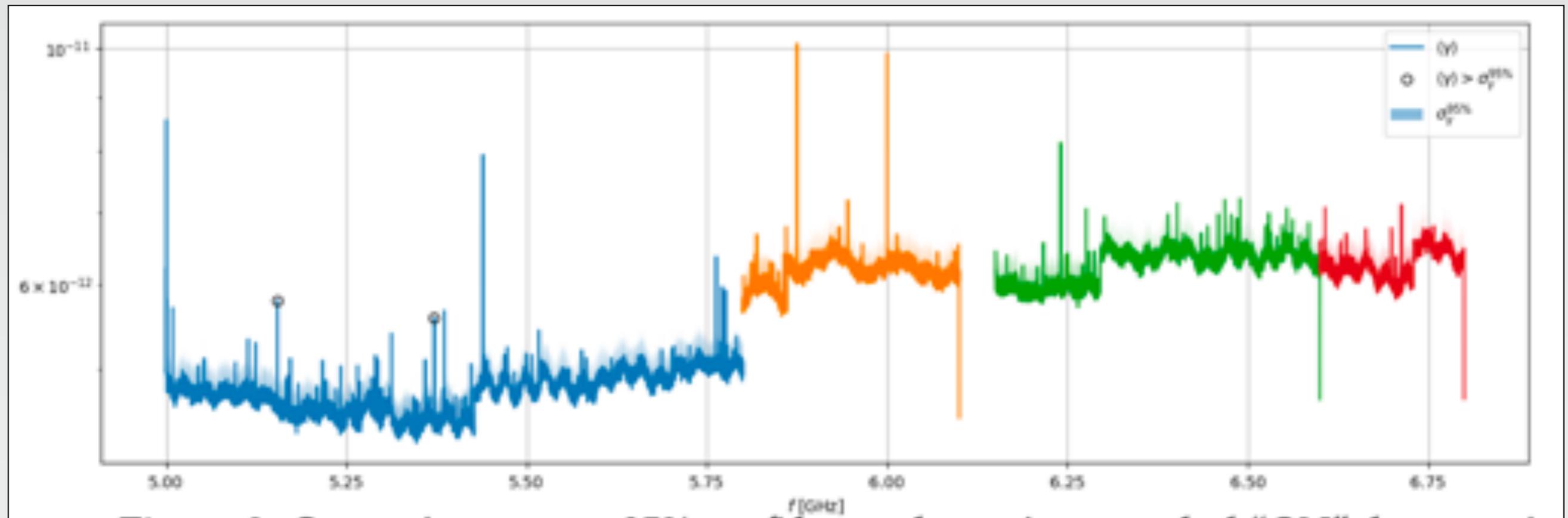
95% threshold



A. HEES, E. SAVALLE, P. WOLF

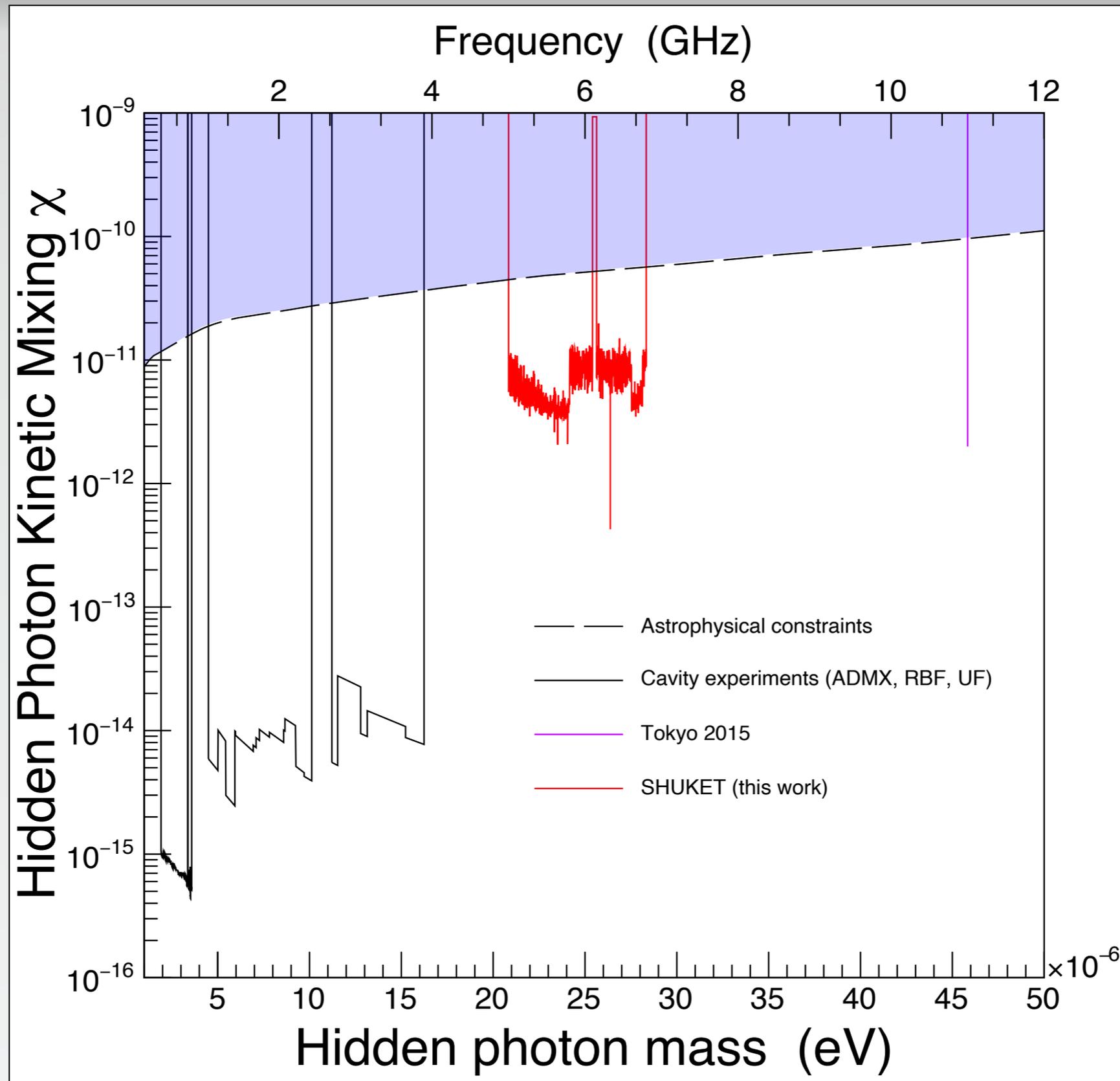
SYRTE re-analysis of Shuket

Exclusion lines on χ

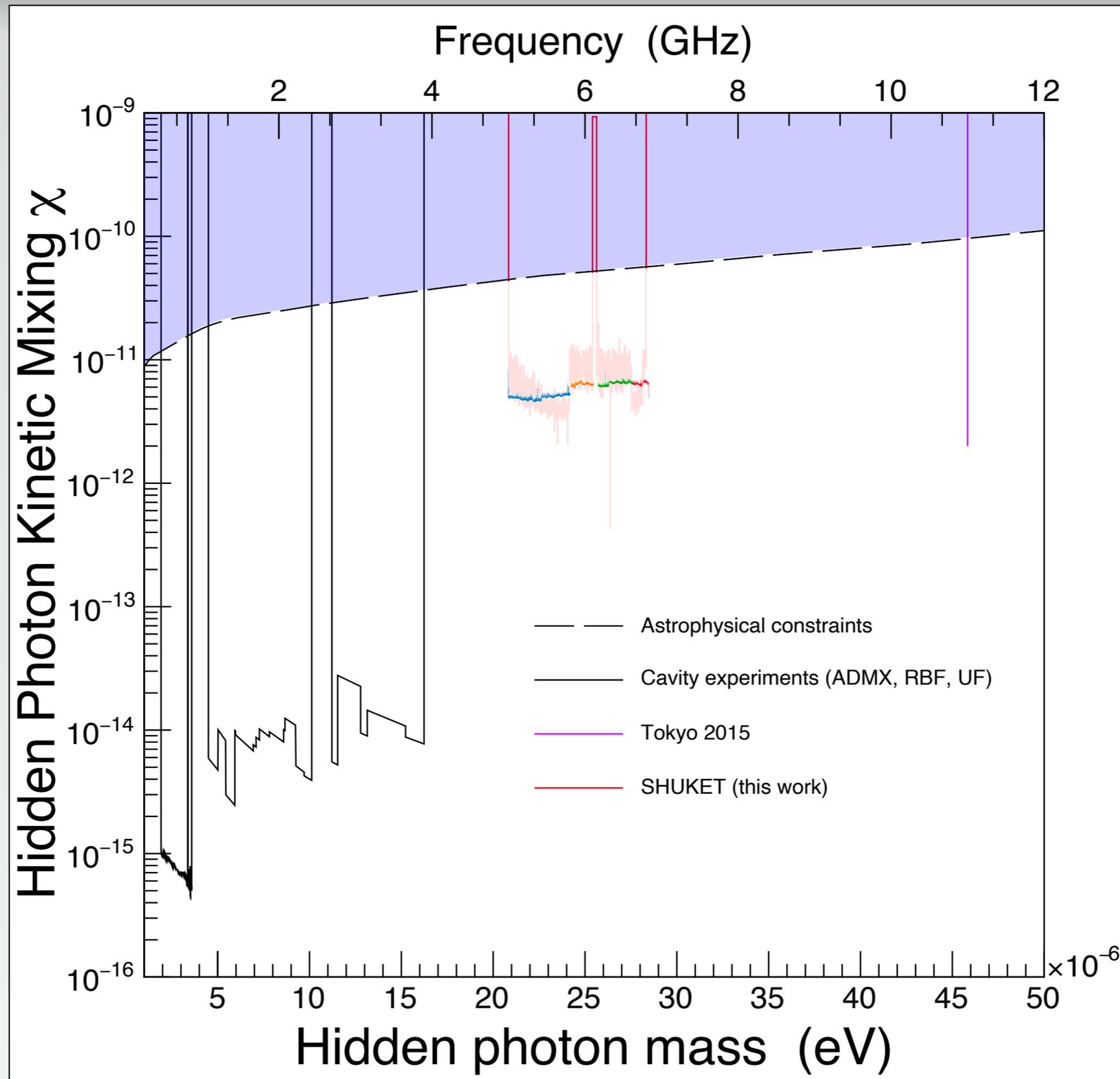


A. HEES, E. SAVALLE, P. WOLF

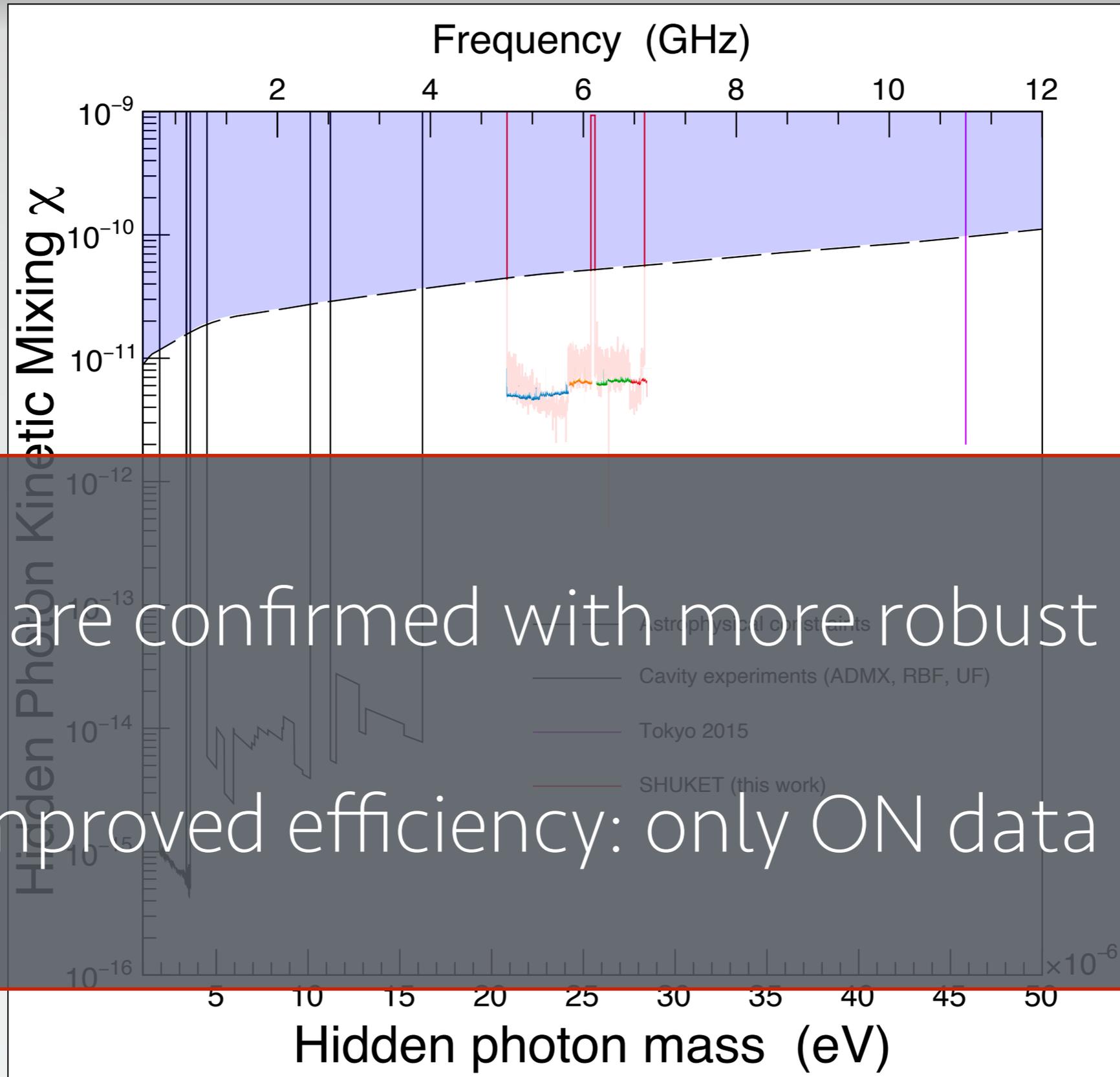
Constraints comparison



Constraints comparison



Constraints comparison



Results are confirmed with more robust method

x2 improved efficiency: only ON data used

Future plans : Hidden photons

- ★ CEA/SYRTE collaboration
- ★ New Shuket runs 10 GHz - 20 GHz
- ★ New, cryogenic low-noise amplifiers
- ★ Data taking optimized for SYRTE-style analysis
- ★ New reflector (could be cold)
- ★ New antenna
- ★ New spectrum analyzer



$$T_{\text{noise}} < 5 \text{ K}$$

Future plans : axions

- ★ Build a magnetic experiment
- ★ Use all that was learnt on the detection part
- ★ Solve the strong CP problem
- ★ Find out what dark matter is