Displacement Noises in Laser Interferometric Gravitational Wave Detectors

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Slides

http://granite.phys.s.u-tokyo.ac.jp/ michimura/lectures/ GWPhysics20171212.pdf Or https://tinyurl.com/GWPhysics20171212



Advanced LIGO Design Sensitivity



Advanced LIGO Design Sensitivity



Advanced Virgo and KAGRA



Fundamental Noises

- Quantum noise quantum fluctuation of photons
- Thermal noise thermal vibration of mirror surface and suspension system
- Seismic noise ground vibration
- Gravity gradient noise Newtonian force on a mirror



Suspension thermal noise

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Quantum radiation

pressure noise

Shot noise Mirror Gravity thermal noise gradient noise

Contents

- Thermal noise
 - Brownian motion, Johnson-Nyquist noise
 - Fluctuation-dissipation theorem
 - Suspension thermal noise
 - Mirror thermal noise
 - Coating thermal noise
- Seismic noise
 - Seismic vibration
 - Vibration isolation
 - Gravity gradient noise
 - Underground site

Thermal Noise

Brownian Motion

- 1827 Robert Brown Discovered random motion of pollen grains in water
- Also proved that this motion exists in non-organic matter → not related to life







Molecular-Kinetic Theory

- 1905 Albert Einstein Molecular-Kinetic Theory
 - → Quantitative calculations on the movement of small particles in a liquid (Ph. D. Thesis)
- 1908 Jean B. Perrin Experimental verification using colloid liquid
 - → Existence of atoms Statistical physics

Johnson-Nyquist Noise

- 1926 John B. Johnson observed voltage fluctuation across a resister
- 1928 Harry Nyquist

Fig. 4. Voltage-squared vs. resistance component for various kinds of conductors. J. B. Johnson, <u>Phys. Rev. 32, 97 (1928)</u>

Generalization of Brownian Motion

 1951-52 Herbert Callen et al. Fluctuation-Dissipation Theorem

Fluctuation-Dissipation Theorem

• Linear System

• System impedance

$$Z(\omega) = \frac{F(\omega)}{v(\omega)} = \frac{F(\omega)}{i\omega x(\omega)}$$

- Thermal fluctuating force for x is $F_{\rm th}^2(\omega) = 4k_{\rm B}T{\rm Re}(Z(\omega))$
- Analogy with Johnson-Nyquist noise Voltage V ↔ Force F Current I ↔ Velocity v Resistor R ↔ Real part of impedance Z

Fluctuation-Dissipation Theorem

- Thermal fluctuating force for x is $F_{
 m th}^2(\omega) = 4k_{
 m B}T{
 m Re}(Z(\omega))$
- Thermal fluctuation of x will be

$$x_{\rm th}^2(\omega) = \frac{F_{\rm th}^2(\omega)}{\omega^2 |Z(\omega)|^2}$$
$$= \frac{4k_{\rm B}T}{\omega^2} \operatorname{Re}(Y(\omega))$$
admittance
$$Y(\omega) \equiv Z^{-1}(\omega)$$

- Damped harmonic oscillator (viscous damping) $m\ddot{x} + \gamma\dot{x} + kx = F$ $_{
 m Spring} k$ • Admittance $\begin{array}{l} \text{Ittance} \\ Y(\omega) = \frac{i\omega x(\omega)}{F(\omega)} \end{array}$ m $i\omega$ Damper $-\omega^2 m + i\omega\gamma + k$ $= \frac{\gamma + i(k/\omega - \omega m)}{\gamma^2 + (k/\omega - \omega m)^2}$
- Thermal noise

$$x_{\rm th}^2(\omega) = \frac{4k_{\rm B}T}{\omega^2} \operatorname{Re}(Y(\omega)) = \frac{4k_{\rm B}T\gamma}{\gamma^2\omega^2 + (k - m\omega_{15}^2)^2}$$

Resonant frequency and Q value

$$\omega_0 = \sqrt{\frac{k}{m}} \qquad \qquad Q = \frac{m\omega_0}{\gamma}$$

• Thermal noise

$$x_{\rm th}^2(\omega) = \frac{4k_{\rm B}T\gamma}{\gamma^2\omega^2 + (k - m\omega^2)^2}$$
$$= \frac{4k_{\rm B}T\omega_0}{mQ} \frac{1}{(\omega^2 - \omega_0^2)^2 + \frac{\omega_0^2\omega^2}{Q^2}}$$
$$\omega \ll \omega_0 \qquad \omega = \omega_0 \qquad \omega \gg \omega_0$$
$$\frac{4k_{\rm B}T}{m\omega_0^3} \frac{1}{Q} \qquad \frac{4k_{\rm B}T}{m\omega_0^3} Q \qquad \frac{4k_{\rm B}T\omega_0}{m\omega^4} \frac{1}{Q}$$

 Resonant frequency and Q-value can be measured from the impulse response

 Viscous damping (e.g. gas damping, eddy-current damping)

$$m\ddot{x} + \gamma\dot{x} + kx = F$$

- Structural damping (e.g. internal friction of materials)
 - elastic material $F_{
 m spring} = -kx$ (Hooke's Law)
 - anelastic material \scale loss angle $F_{
 m spring} = -k(1+i\phi(\omega))x$

Delay between displacement and force \rightarrow energy dissipation

- Damped harmonic oscillator (structural damping) $m\ddot{x} + k(1 + i\phi)x = F$
- Admittance $Y(\omega) = \frac{k\phi\omega + i(k\omega - m\omega^3)}{k^2\phi^2 - (k - m\omega^2)^2} \xrightarrow{\text{Optimal}}_{\text{Damper}ik\phi} m$
- Thermal noise

$$x_{\rm th}^2(\omega) = \frac{1}{\omega} \frac{4k_{\rm B}Tk\phi}{k^2\phi^2 + (k - m\omega^2)^2}$$

Resonant frequency and Q value

$$\omega_0 = \sqrt{\frac{k}{m}} \qquad \qquad Q = \frac{1}{\phi(\omega)}$$

• Thermal noise

$$x_{\rm th}^2(\omega) = \frac{1}{\omega} \frac{4k_{\rm B}Tk\phi}{k^2\phi^2 + (k - m\omega^2)^2}$$
$$= \frac{4k_{\rm B}T\omega_0^2\phi}{m\omega} \frac{1}{\omega_0^4\phi^2 + (\omega^2 - \omega^2)^2}$$
$$\omega \ll \omega_0 \qquad \omega = \omega_0 \qquad \omega \gg \omega_0$$
$$\frac{4k_{\rm B}T}{m\omega_0^2} \frac{\phi}{\omega} \qquad \frac{4k_{\rm B}T}{m\omega_0^3} \frac{1}{\phi} \qquad \frac{4k_{\rm B}T\omega_0}{m\omega^4} \frac{\phi}{\omega}$$

- Thermal noise spectrum when $\phi(\omega) \sim {\rm const.}$

- Suspension thermal noise ($\omega \gg \omega_0$) • $x_{\rm susp}^2(\omega) = \frac{4k_{\rm B}T\omega_0}{m\omega^4} \frac{\phi_{\rm p}}{\omega},$ $=\frac{4k_{\rm B}T}{m\omega^5}\sqrt{\frac{g}{l}}\frac{\phi}{2l\omega}\sqrt{\frac{EI}{mg}}$
- To lower the thermal noise
 - lower temperature $x_{susp}(\omega) \propto \sqrt{T\phi}$
 - lower loss angle
 - longer suspension $x_{
 m susp}(\omega) \propto l^{-3/2}$
 - heavier mirror

 $x_{\rm susp}(\omega) \propto m^{-1/4}$ assuming $I \propto A^2$ $m \propto A$ A: wire cross section

m

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- Suspension thermal noise ($\omega \gg \omega_0$) $x_{susp}^2(\omega) = \frac{4k_{\rm B}T\omega_0}{m\omega^4} \frac{\phi_{\rm p}}{\omega}$
- Calculate suspension thermal noise for m=10 kg, f₀= 1Hz, T = 300 K, $\varphi_p = 10^{-6}$ at 100 Hz. $k_{\rm B} = 1.38 \times 10^{-23} \, {\rm J/K}$

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- Calculate suspension thermal noise for m=10 kg, f₀= 1Hz, T = 300 K, $\varphi_p = 10^{-6}$ at 100 Hz. $k_B = 1.38 \times 10^{-23} \text{ J/K}$

$$x_{susp}(\omega) = 1.0 \times 10^{-20} \left(\frac{100 \text{ Hz}}{f}\right)^{5/2} \text{ m}/\sqrt{\text{Hz}}$$

Sufficient for GW detection!

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Suspension Thermal Noise in 2G

- Advanced LIGO (US) 40 kg, 60 cm, fused silica, 295 K, φw=~1e-7
- Advanced Virgo (Italy)_{10⁻²⁰}
 42 kg, 70 cm,
 fused silica, ¹² 10⁻²¹
 295 K, φw=~1e-7 ¹
- KAGRA (Japan)
 23 kg, 23 cm,
 sapphire,
 22 K, φw=2e-7

Suspensions

Challenges in Cryogenics

- KAGRA is the first km-scale GW detector to use cryogenics to reduce thermal noise
- Fused silica has large mechanical loss at cryogenic temperatures
 → sapphire
- Limitation in mirror size: 23 kg (small compared with aLIGO/AdV)
 → smaller beam size

worse quantum and thermal noise

- Thicker wire for heat extraction
 → larger suspension thermal noise
- Vibration of cryocoolers

KAGRA Cryogenics

• Thermal vibration of mirror surface

- Thermal vibration of mirror surface
- Structural damping, $\omega \ll \omega_0$

$$r_{\rm mir}^2(\omega) = \frac{4k_{\rm B}T}{m_{\rm eff}\omega_0^2} \frac{\phi}{\omega}$$

- *m*_{eff} is reduced mass function of mirror aspect ratio, mirror radius and beam radius ratio, Poisson's ratio
- ω_0 is resonant frequency of vibration mode $\omega_0 \propto \sqrt{E}$
- You also have to sum up all the vibration modes

• Simple formula

$$x_{\rm mir}^2(\omega) = \frac{4k_{\rm B}T\phi}{\omega} \frac{1-\sigma}{\sqrt{\pi}Ew}$$

- To lower the thermal noise
 - lower temperature
 - lower loss angle $x_{\rm mir}(\omega) \propto \sqrt{T\phi}$
 - larger beam size (needs larger mirror)

 $x_{
m mir}(\omega) \propto 1/\sqrt{w}$

• Simple formula

$$x_{\rm mir}^2(\omega) = \frac{4k_{\rm B}T\phi}{\omega} \frac{1-\sigma}{\sqrt{\pi}Ew}$$

• Calculate mirror thermal noise for fused silica (E = 73 GPa, σ = 0.17, φ =1e-6) w = 5 cm, T = 300 K, at 100 Hz $k_{\rm B} = 1.38 \times 10^{-23} \,\text{J/K}$

• Simple formula

$$x_{\rm mir}^2(\omega) = \frac{4k_{\rm B}T\phi}{\omega} \frac{1-\sigma}{\sqrt{\pi}Ew}$$

• Calculate mirror thermal noise for fused silica (E = 73 GPa, σ = 0.17, ϕ =1e-6) w = 5 cm, T = 300 K, at 100 Hz $k_{\rm B} = 1.38 \times 10^{-23} \,\text{J/K}$

$$x_{\rm mir}(\omega) = 5.8 \times 10^{-20} \left(\frac{100 \,\mathrm{Hz}}{f}\right)^{1/2} \,\mathrm{m}/\sqrt{\mathrm{Hz}}$$

Sufficient for GW detection!

- High reflective coating on a mirror λ/4 thick alternative layers
- Active area of research due to low loss angle
 SiO₂

(φ~1e-4)

constructive

Mirror substrate

• Formula (N. Nakagawa+, PRD 65, 102001 (2002))

$$x_{\rm coat}^2(\omega) = \frac{4k_{\rm B}T\phi}{\omega} \frac{2d_{\rm co}(1+\sigma_{\rm co})(1-2\sigma_{\rm co})}{\pi Ew^2}$$

 $x_{\rm coat}(\omega) \propto \sqrt{T\phi}$

- To lower the thermal noise
 - lower temperature
 - lower loss angle
 - larger beam radius (needs larger mirror)

 $x_{\rm coat}(\omega) \propto 1/w$

• Formula (N. Nakagawa+, PRD 65, 102001 (2002))

$$x_{\text{coat}}^2(\omega) = \frac{4k_{\text{B}}T\phi}{\omega} \frac{2d_{\text{co}}(1+\sigma_{\text{co}})(1-2\sigma_{\text{co}})}{\pi Ew^2}$$

• Calculate coating thermal noise for silica/tantala coating (E = 72 Gpa, σ = 0.17, ϕ = 4e-4) d = 10 um, w = 5 cm, T = 300 K, at 100 Hz $k_{\rm B} = 1.38 \times 10^{-23} \,\text{J/K}$

• Formula (N. Nakagawa+, <u>PRD 65, 102001 (2002)</u>)

$$x_{\text{coat}}^2(\omega) = \frac{4k_{\text{B}}T\phi}{\omega} \frac{2d_{\text{co}}(1+\sigma_{\text{co}})(1-2\sigma_{\text{co}})}{\pi Ew^2}$$

• Calculate coating thermal noise for silica/tantala coating (E = 72 Gpa, σ = 0.17, ϕ = 4e-4) d = 10 um, w = 5 cm, T = 300 K, at 100 Hz $k_{\rm B} = 1.38 \times 10^{-23} \,\text{J/K}$

$$x_{\text{coat}}(\omega) = 1.7 \times 10^{-20} \left(\frac{100 \,\text{Hz}}{f}\right)^{1/2} \,\text{m}/\sqrt{\text{Hz}}$$

Sufficient for GW detection!

Mirror/Coating Thermal Noise in 2G

- Advanced LIGO (US) w= 5.5 / 6.2 cm, fused silica, 295 K, φ=~1e-6
- Advanced Virgo $(Italy)_{10^{-20}}$ w= 4.9 / 5.8 cm, fused silica, $[\frac{N}{2}]_{10^{-21}}$ 295 K, ϕ =~1e-6
- KAGRA (Japan) w= 3.5 / 3.5 cm, sapphire, 22 K, φ=1e-8

Mirror thermal noise is not significant, but coating thermal noise is troublesome in 2G

Other Coating Thermal Noises

- From thermal dissipation in the coating leads to temperature fluctuation in the coating, which causes:
 - Thermo-elastic noise thermal expansion
 - Thermo-refractive noise refractive index change
- These two can cancel

M. Evans+, PRD 78, 102003 (2009)

Seismic Noise

Seismic Vibration

- Ground vibrates even if there is no earthquakes
- Vibration level differs by day and night stormy day and calm day urban place and countryside surface and underground
- Site selection is very important for GW detectors

GW Detector Sites

LIGO Livingston

virgo

3 km

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Seismic Vibration Between Sites

Vibration Isolation

• Typical seismic noise (according to Saulson)

$$x_{\rm seis}(\omega) = \begin{cases} 10^{-9} \,\mathrm{m}/\sqrt{\mathrm{Hz}} & 1-10 \,\mathrm{Hz} \\ 10^{-9} \left(10 \,\mathrm{Hz}/f\right)^2 \,\mathrm{m}/\sqrt{\mathrm{Hz}} &>10 \,\mathrm{Hz} \end{cases}$$

- Needs vibration attenuation by ~9 orders of magnitude
- Use pendulums to attenuate the ground vibration

Simple Pendulum

Transfer function from x_{in} to x_{out}

$$H(\omega) = \frac{x_{\text{out}}(\omega)}{x_{\text{in}}(\omega)} = \frac{\frac{g}{l}}{-\omega^2 + \frac{g}{l} + i\frac{\gamma}{m}\omega} \qquad m$$
$$= \frac{\omega_0^2}{-\omega^2 + \omega_0^2 + i\frac{\omega\omega_0}{Q}} \quad \overleftarrow{x_{\text{out}}}$$

where

$$\omega_0 = \sqrt{\frac{g}{l}} \qquad Q = \frac{m\omega_0}{\gamma}$$

Resonant Frequency

 Longer suspension gives lower resonant frequency, and thus better attenuation ratio at high frequencies

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Multiple-Stage Pendulum

- n-stage pendulum has f⁻²ⁿ attenuation at high frequencies
- Gives enough isolation at observation band of ground-based GW detectors

Pendulum Q-value

- Higher Q-value gives better thermal noise, but needs longer time it takes to calm down
- Very low Q-value is bad for seismic attenuation (Q < 0.5 is over damping)

Active Damping

- Measure mirror displacement and give damping force proportional to mirror velocity using differential circuit
- No degradation in thermal noise
- Control noises are introduced circuit noise, magnetic noise, sensor noise, actuator noise......

Differential circuit (High-pass filter)

OSEM

Optical Sensor ElectroMagnet

KAGRA OSEM

Vertical Vibration Isolation

 Vertical vibration isolation is also necessary since the beam axis is not perfectly horizontal

Low Frequency Vibration Isolation

 Usual pendula are not practical for making resonant frequency low

pendulum

Inverted pendulum

Low Frequency Vibration Isolation

Vibration Isolation System in 2G

Advanced LIGO

hydraulic external preisolator (HEPI) (one stage of isolation) active isolation platform (2 stages of isolation)

quadruple pendulum (four stages of isolation) with monolithic silica final stage

Gravity Gradient Noise

Newtonian gravitational force acting on a mirror

Gravity Gradient Noise

- Newtonian gravitational force acting on a mirror
- Could be an issue for 2G

M. Coughlin+, CQG 33 244001 (2016)

FIG. 1: Predicted average Newtonian noise in Advanced LIGO.

Underground GW Detector

- Underground site gives lower seismic noise (and lower gravity gradient noise)
- KAGRA is the first km-scale underground detector

KAGRA site

(March 2014)

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 Future GW project Einstein Telescope also considers using underground site

Einstein Telescope

ET EINSTEIN TELESCOPE

- Triangular, 10 km, underground
- Xylophone configuration LF detector with cryogenic and low power HF detector with room temperature and high power

Einstein Telescope

Roughly x10 the 2G sensitivity

EINSTEIN

Assignment from Raffaele Flaminio (by next week)

Assignment n.1

- Consider a Fabry-Perot cavity made of two mirrors M1 and M2. M1 is the input mirror and it has amplitude reflectivity r. M2 is the end mirror and it is perfect reflector. The cavity wavelength is L.
- E_i is the electric field amplitude of the laser beam entering into the cavity. λ is its wavelength. The cavity is kept in resonance with the laser field.
- A gravitational wave produces a small variation of the cavity length $x(t) << \lambda$ with linear spectral density $\tilde{x}(f)$.
- 1. Calculate the variation of the electric field $\tilde{E}_c(f)$ inside the cavity.
- 2. Calculate the variation of the electric field $\tilde{E}_c(f)$ reflected by the cavity.

Assignment n.2

- Consider a double pendulum made of two pendulums in cascade. The double pendulum temperature is T=300 K.
- The first pendulum has a length $I_1=1$ m and a mass $m_1=100$ kg. the second one hangs form the first mass and has a length $I_2=1$ m and a mass of $m_2=50$ kg.
- The first pendulum is affected by a mechanical loss of $\phi_1 = 10^{-4}$ i.e. the gravitational restoring force has an imaginary part equal to 10^{-4} of its real part.
- The second pendulum is affected by a mechanical loss of $\phi_2=10^{-6}$ i.e. the gravitational restoring force has an imaginary part equal to 10⁻⁶ of its real part.
- Using the fluctuation dissipation theorem, calculate the spectrum of the second mass displacement due to thermal noise.
- Show that at high frequency the effect of the losses in the upper pendulum can be neglected.
- Give an interpretation of this result.

Summary

- Thermal noise
 - suspension / mirror / coating thermal noise
 - cryogenics, low loss angle, longer suspension, and larger mirror reduces thermal noise
 - cryogenics is attractive but challenging
- Seismic noise
 - site selection is very important
 - multiple-stage, low resonant frequency vibration isolation techniques are used
 - underground site is preferred for reducing seismic noise and gravity gradient noise
- KAGRA is currently the only km-scale cryogenic underground GW detector in the world 68