



# Atmospheric Newtonian Noise (NN)

## Infrasound NN

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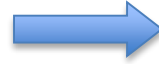
<sup>b</sup>Università degli Studi di Urbino 'Carlo Bo'

# Atmospheric sources of Newtonian Noise

- Atmospheric weak pressure waves (small  $\delta p/p$ )
- Atmospheric temperature perturbations
- Atmospheric shockwaves
- High-speed massive objects moving near the interferometer

# Why modeling infrasound NN?

aLIGO, AdVirgo, KAGRA



Frequency range  $\approx 10\text{Hz} - 30\text{Hz}$

Einstein Telescope (ET)



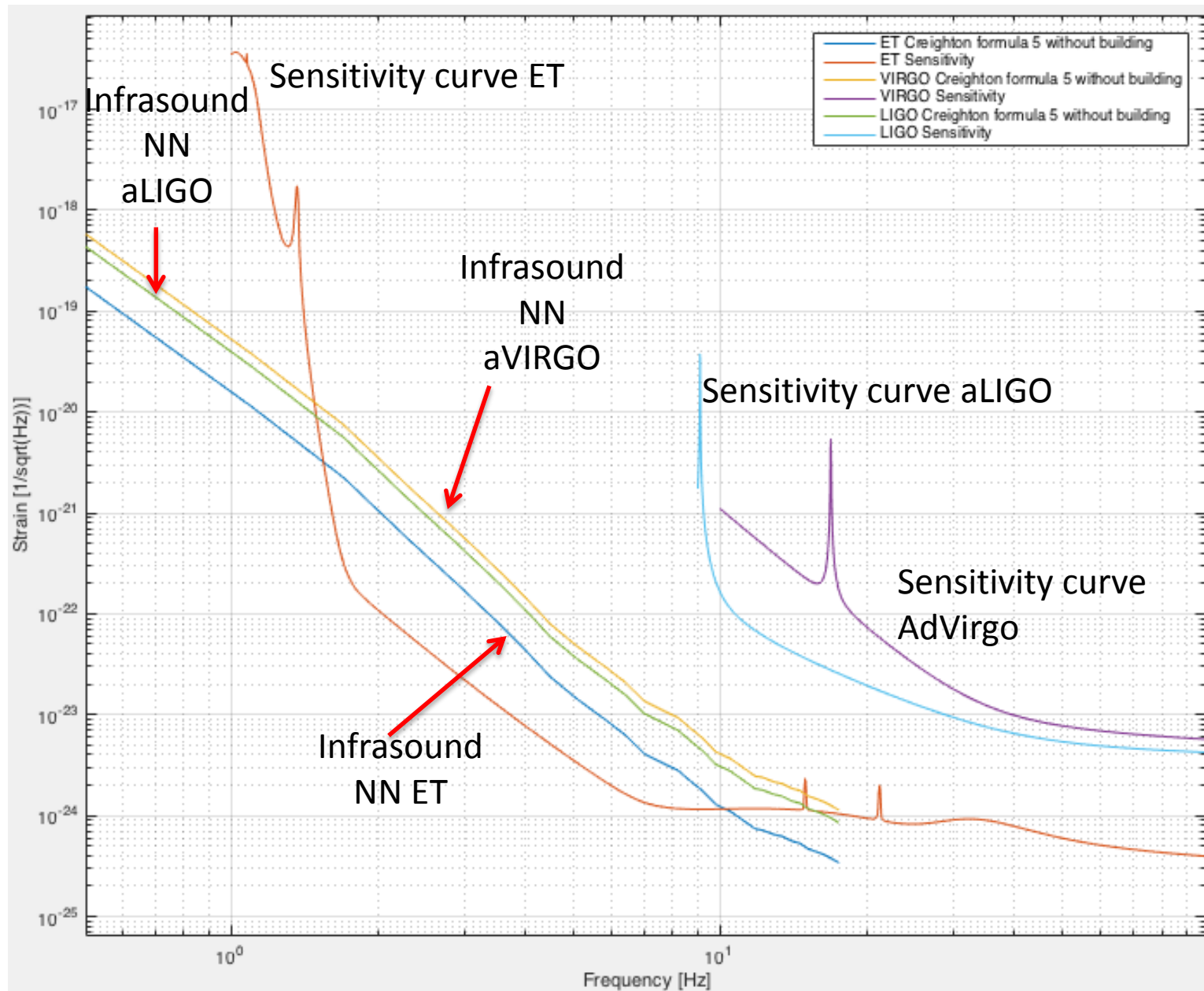
Frequency range  $\approx 1\text{Hz} - 30\text{Hz}$

Torsion bar antennae and  
other low frequency  
detectors (i.e. TOBA,  
TORPEDO, atom  
interferometers ...)



Frequency range  $\approx 10\text{ mHz}-1\text{Hz}$

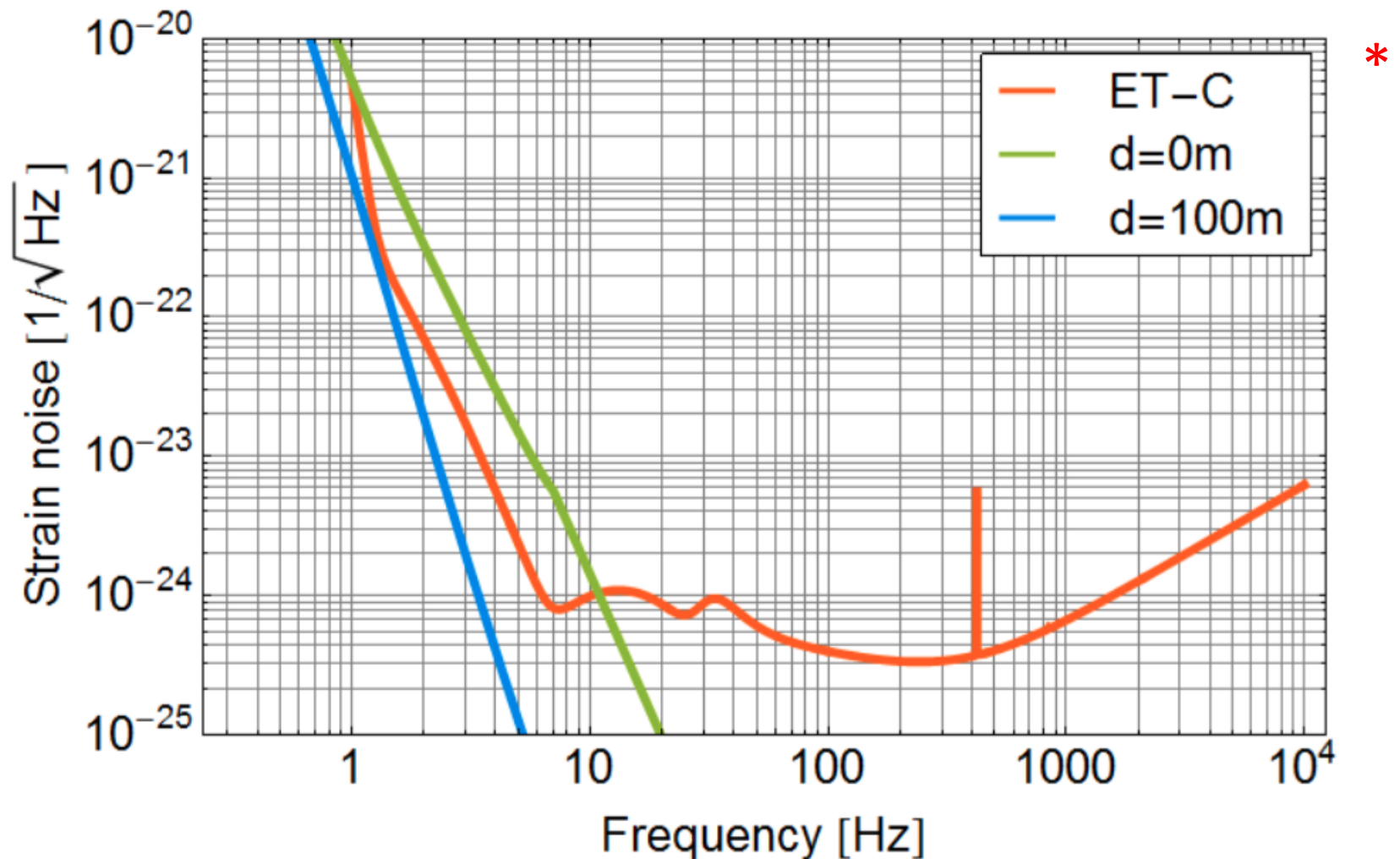
# Why modeling infrasound NN?



\* Infrasound NN estimates obtained by using Creighton's model, see [CQG. 25 \(2008\) 125011](#)

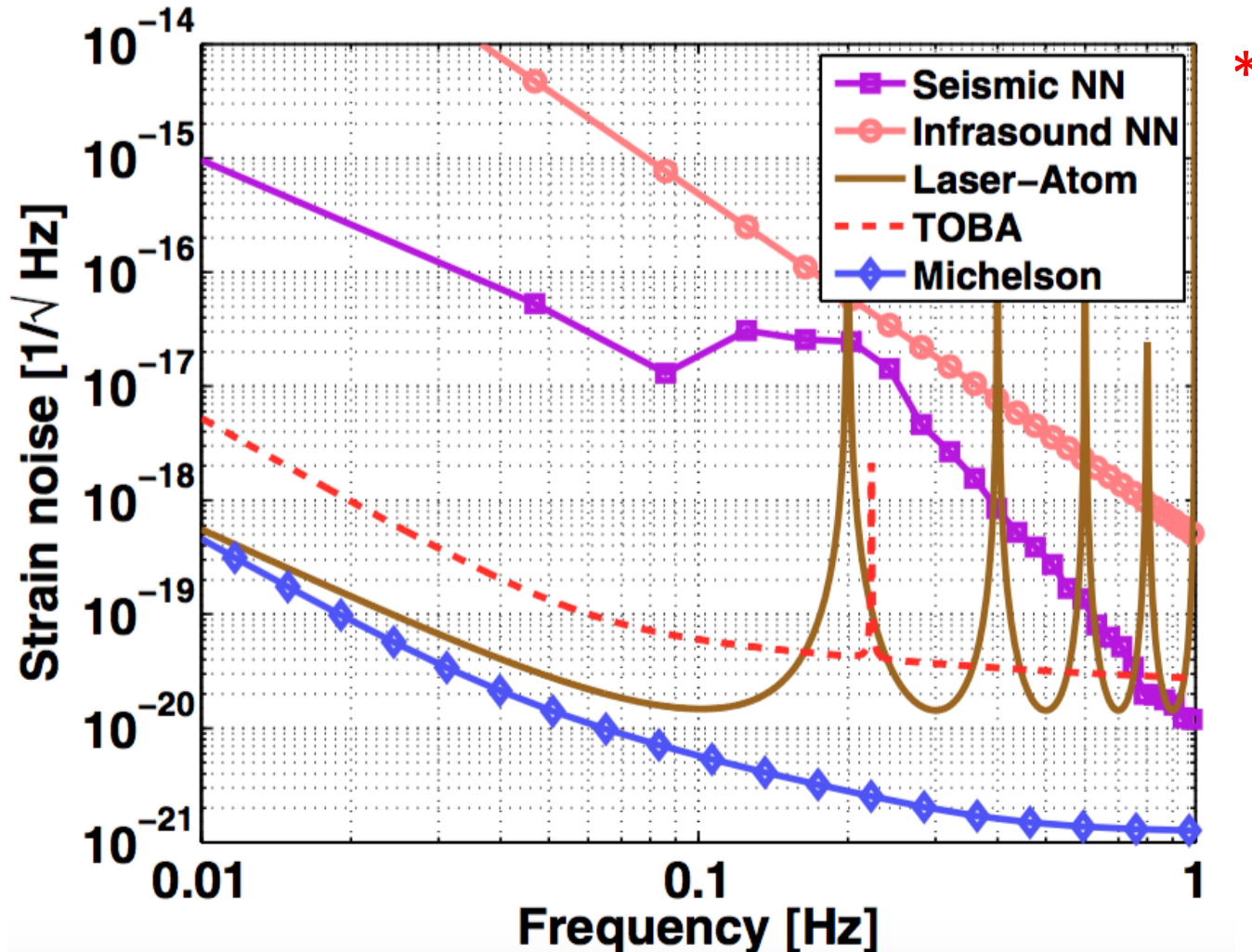
# Why modeling infrasound NN?

ET Infrasound Newtonian Noise for different detector depths,  $d$ .



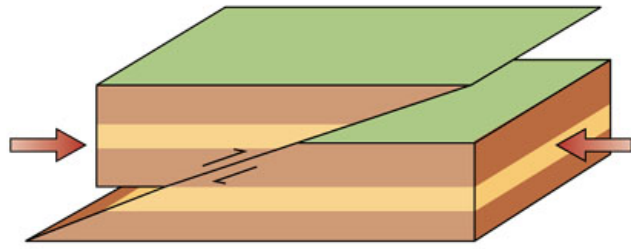
# Why modeling infrasound NN?

Infrasound NN in low frequency detectors for GW and earthquakes (TOBA, TORPEDO, atom interferometers)



# E-GRAAL (Earthquake GRAvity Alerts) Project

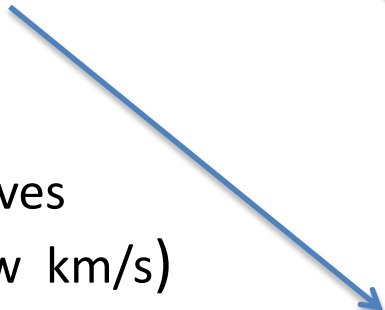
Earthquake



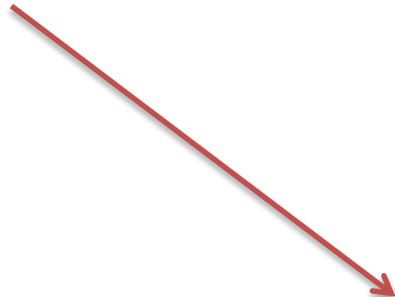
Thrust fault

**E-GRAAL** (Earthquake GRAvity Alerts): Feasibility study of a new earthquake early warning system based on prompt gravity perturbations from earthquakes

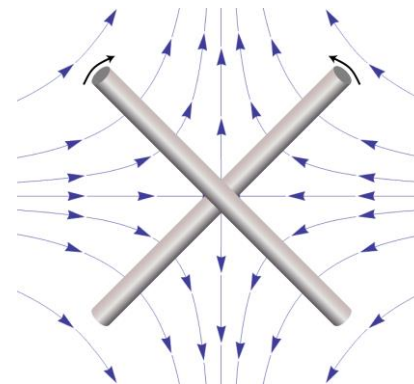
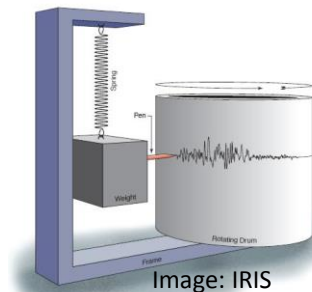
p-waves  
(a few km/s)



Gravitational field change

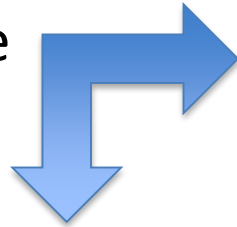


Seismic signal



# Gravity perturbations from Infrasound NN

Plane pressure wave



$\delta p/p \ll 1$ , frequency  $f$ , sound speed  $c$

Adiabatic density change  $\delta \rho/\rho = \delta p/\gamma p$ ,  $\gamma=1.4$



Gravitational acceleration caused by the waves, along its direction of propagation

$$g_z(t) = G \int \frac{dr(t)}{r^3} dV$$

$$\tilde{h}(f) = (2\pi f)^{-2} \tilde{g}(f) / L \quad \longrightarrow \quad \text{Spectral density} = S_h(|f|) = \langle \tilde{h}(f) \tilde{h}(f')^* \rangle$$

Interferometer arm length

Average over the plane wave modes contributing to the noise



# Issues on Infrasound NN

$$g_z(t)^* = \int \frac{G_z \delta \rho}{r^3} dV = \frac{G \rho c}{\gamma p f} \cos(\theta) \overset{2}{\boxed{C(2\pi f r_{\min}/c)}} \overset{1}{\boxed{\delta p(t + 1/4 f)}}$$

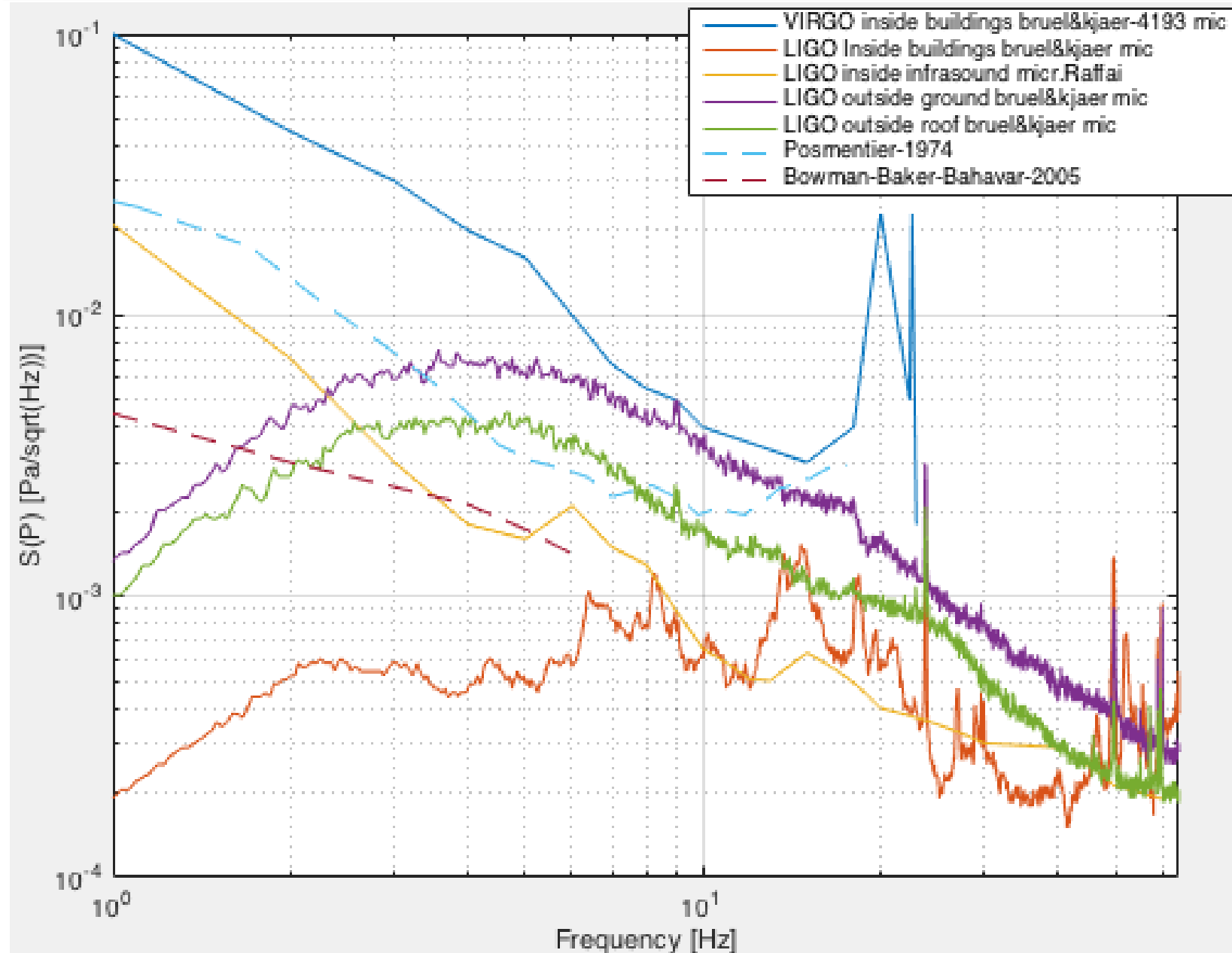
Angle between the wave propagation direction and the interferometer arm

1) Measurement of pressure fluctuations at infrasound frequencies →

- How to perform the measures
- Where to take the measures

2) Effect of the building housing the test mass → further analysis needed

# Pressure Fluctuations

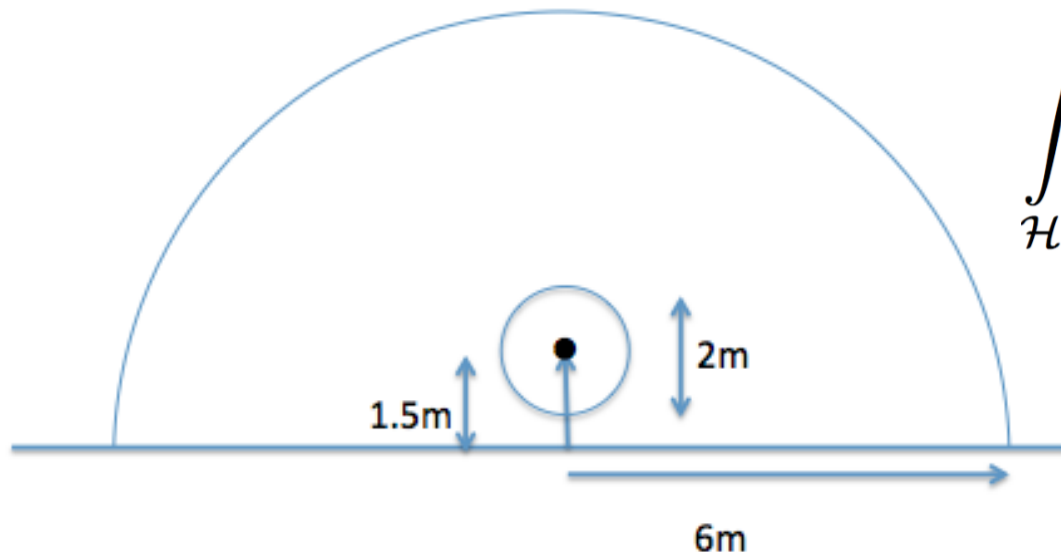


Need for  
data in the  
sub-Hz  
frequency  
range

# Building effect modeling

**Considered geometry:** **hemispheric building**, 6m radius, centered at  $x_0=0\text{m}, y_0=0\text{m}, z_0=0\text{m}$

**spheric vacuum chamber** of radius 1m ,  
centered at  $x_0=0\text{m}, y_0=0\text{m}, z_0=1.5\text{m}$

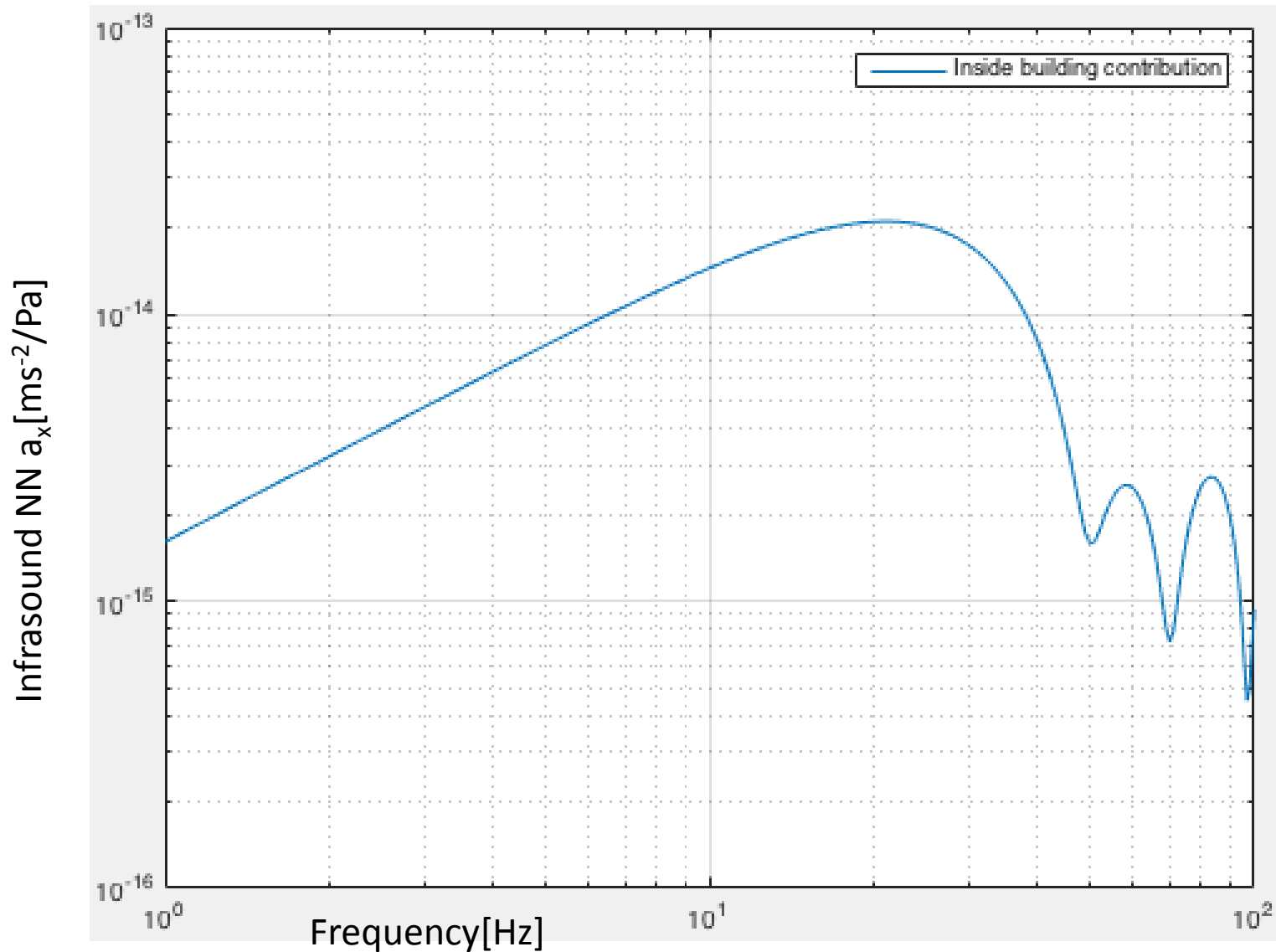


$$\int_{\mathcal{H}} dV \frac{(e^{-ik_z z} + e^{ik_z z}) e^{-i\vec{k}_\perp \cdot \vec{\varrho}}}{(\varrho^2 + (z - z_0)^2)^{1/2}} *$$

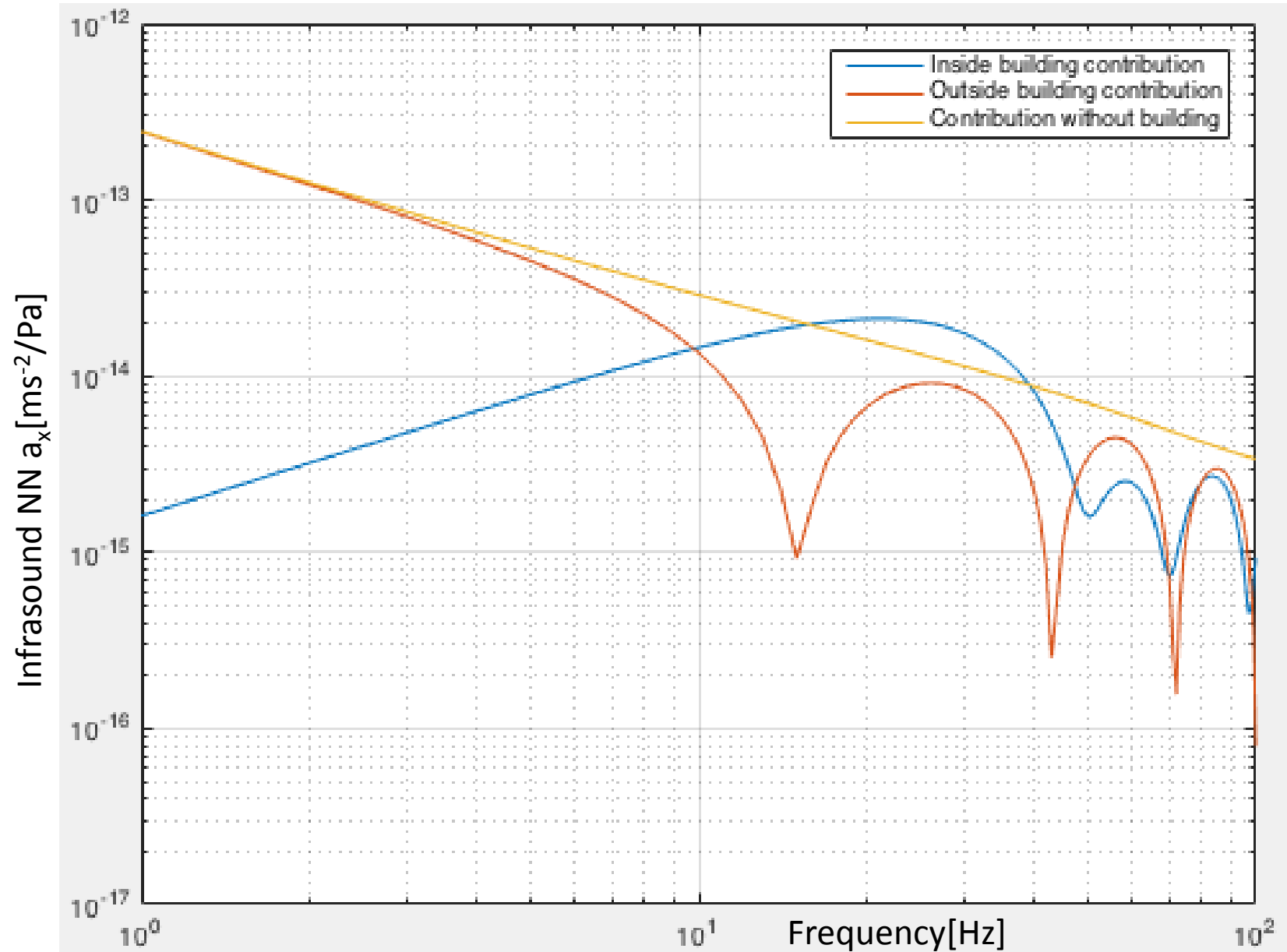
To be calculated to find  
the infrasound NN  
gravitational acceleration

# Building effect modeling

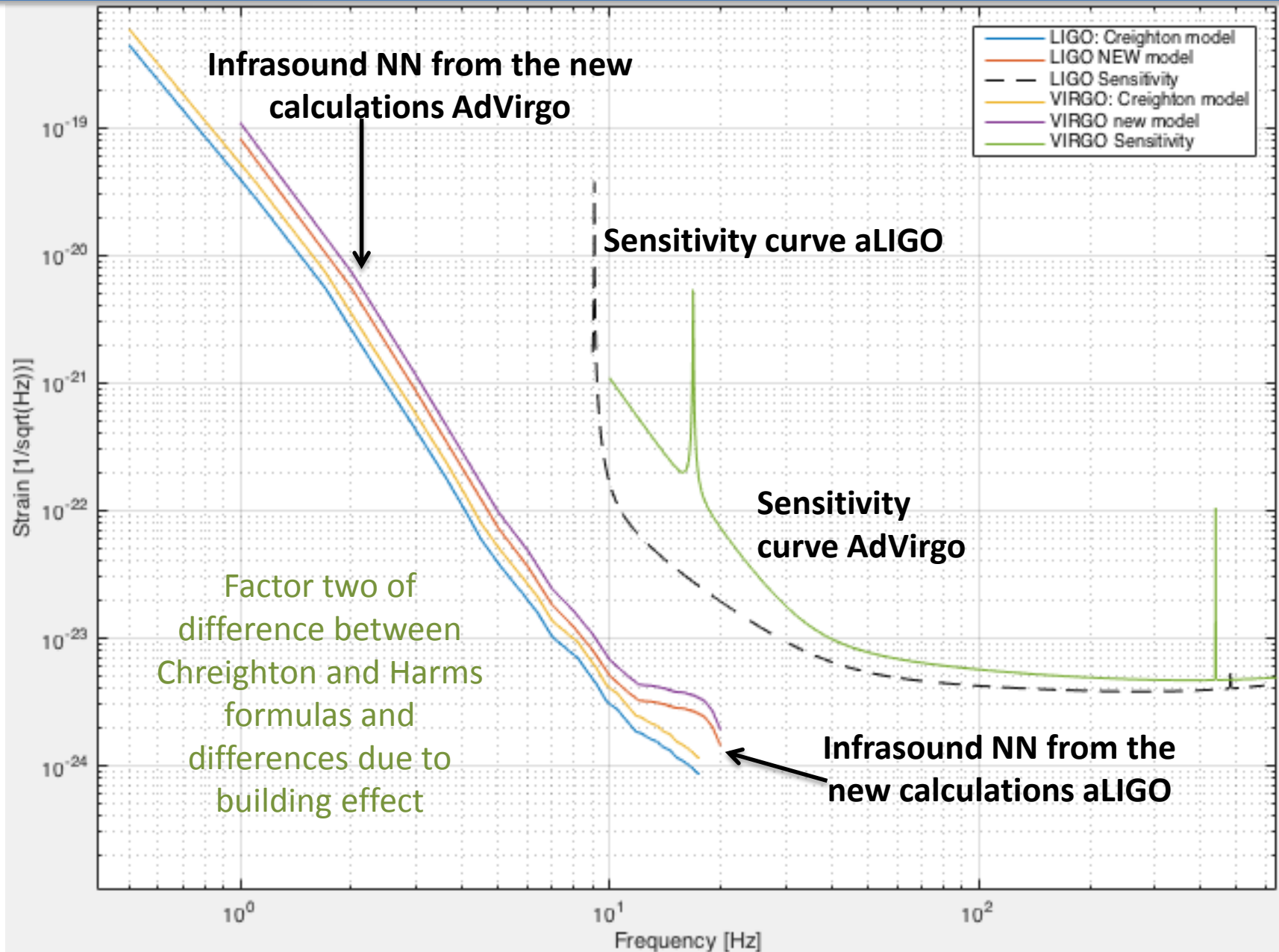
Inside building, numerical integration, average over 60 acoustic wave directions



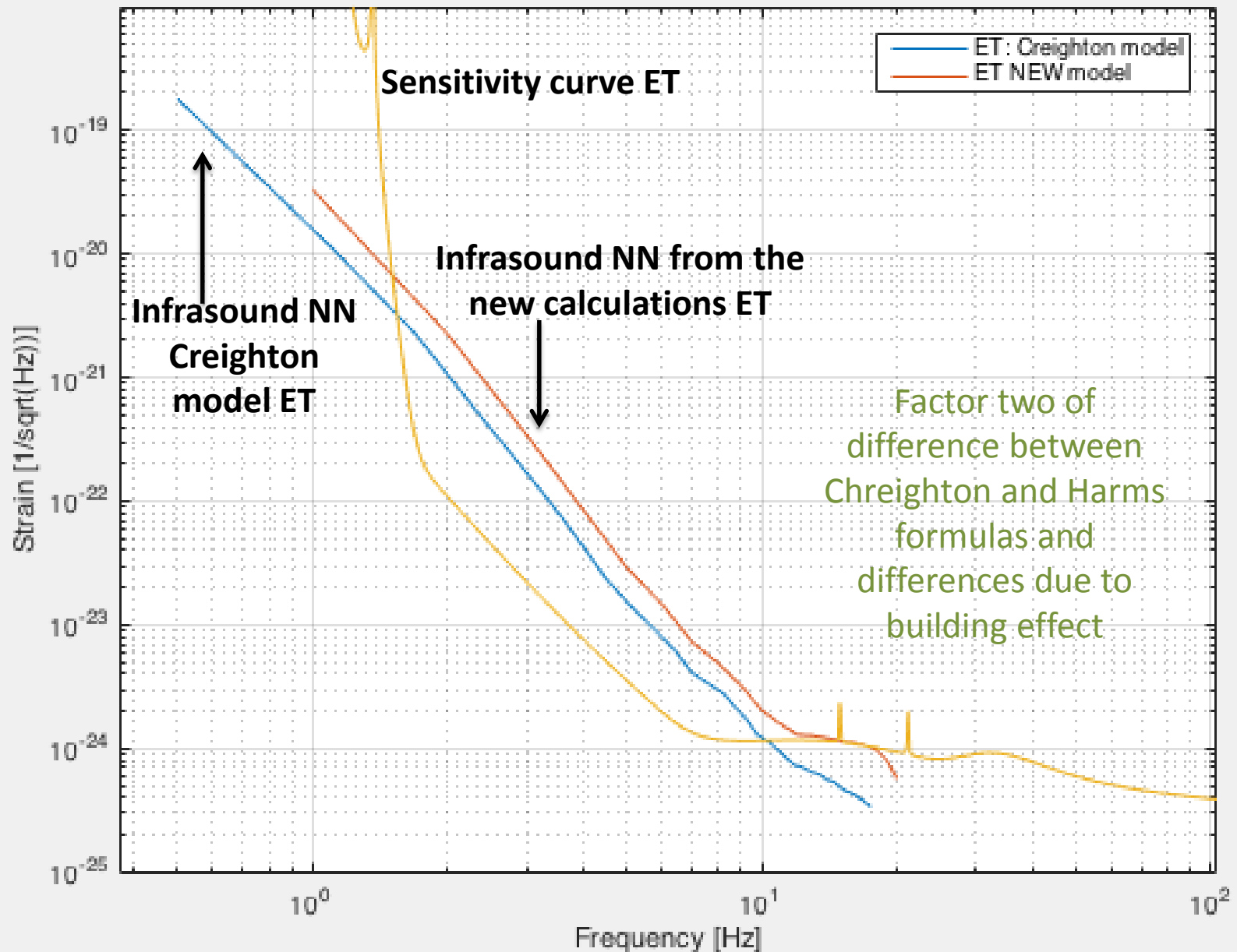
# Building effect modeling-Result Comparison/1



# Building effect modeling-Result Comparison/2



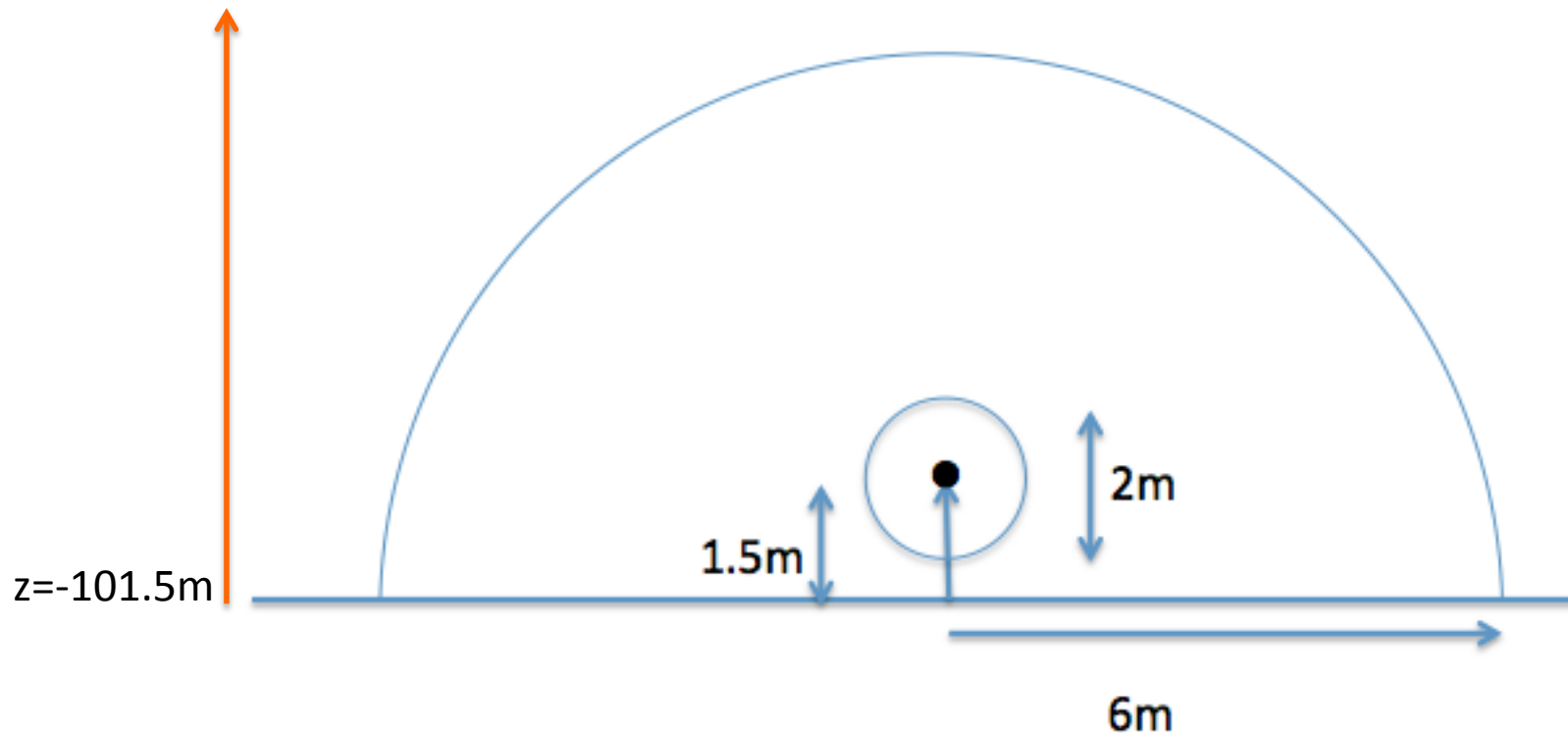
# Building effect modeling-Result Comparison/3



# Infrasound NN-Underground case

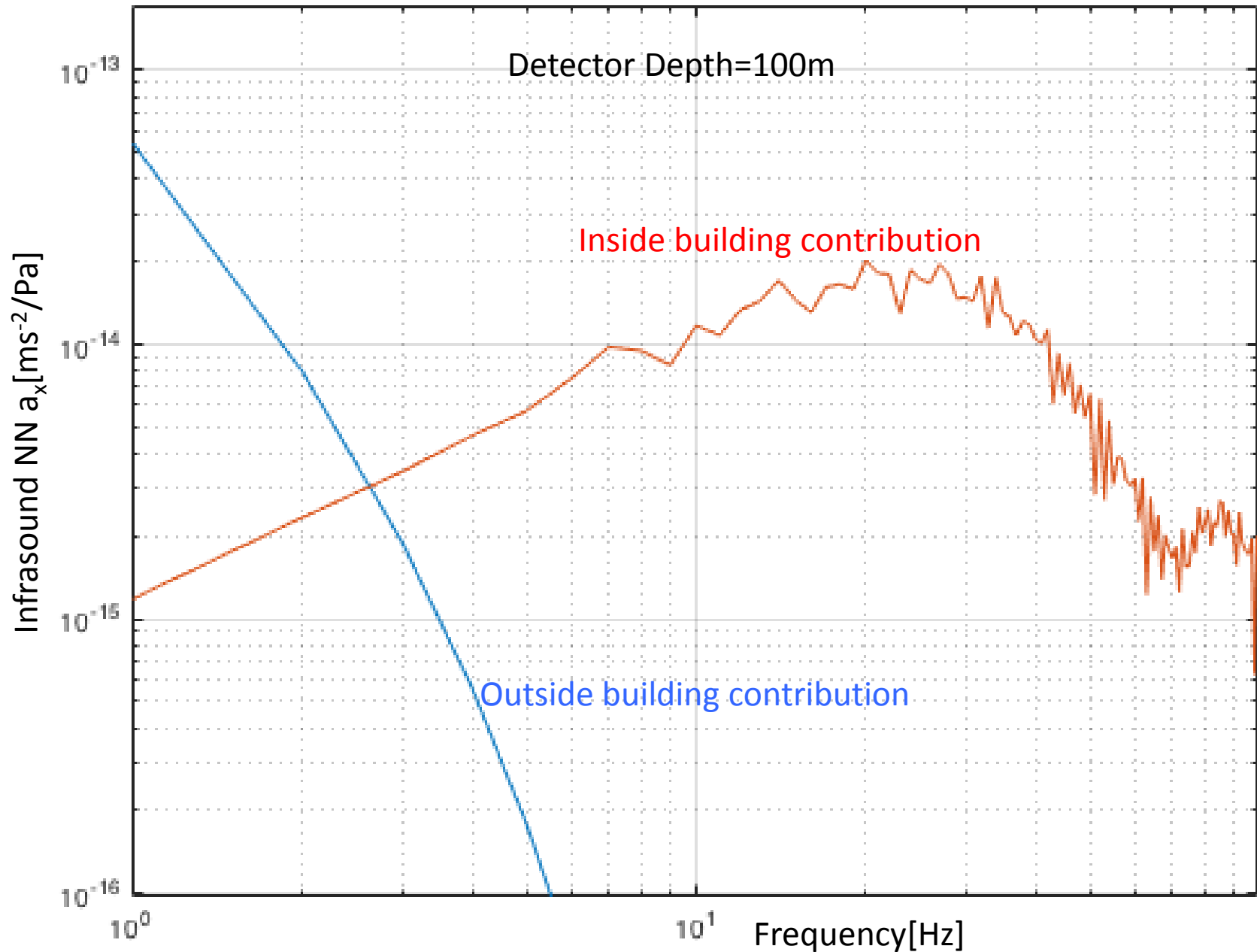
**Considered geometry:** **hemispheric building**, 6m radius, centered at  $x_0=0\text{m}, y_0=0\text{m}, z_0=-101.5\text{m}$

**spheric vacuum chamber** of radius 1m ,  
centered at  $x_0=0\text{m}, y_0=0\text{m}, z_0=-100\text{m}$

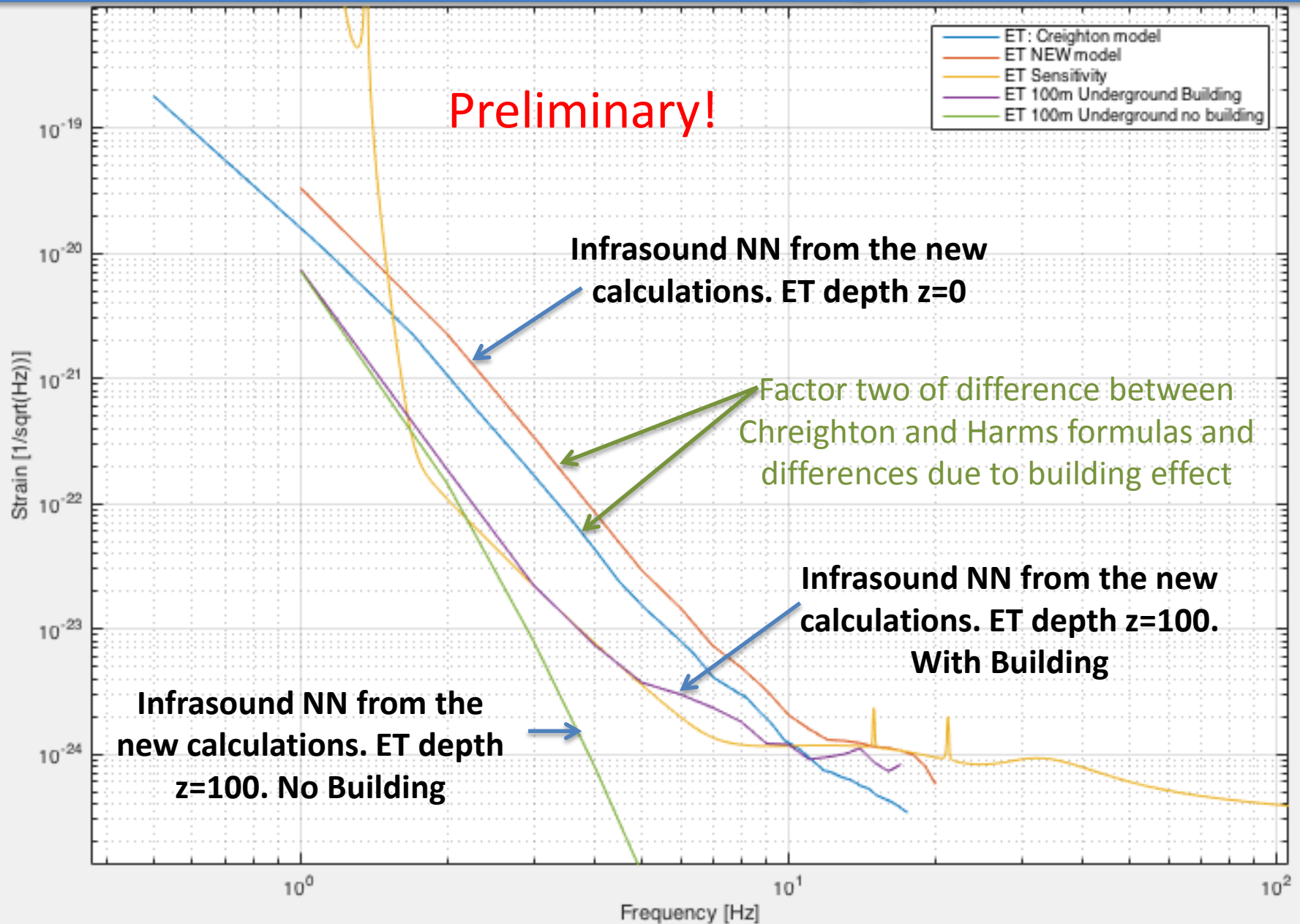




# ET Infrasound NN-Underground case



# ET Infrasound NN-Underground case



$$\delta \vec{g}_{12}(\omega) = -\nabla \psi(\vec{r}_2, \omega) + \nabla \psi(\vec{r}_1, \omega)$$



Small test mass distance



$$\approx -(\nabla \otimes \nabla \psi(\vec{r}_1, \omega)) \cdot \vec{r}_{12}^*$$



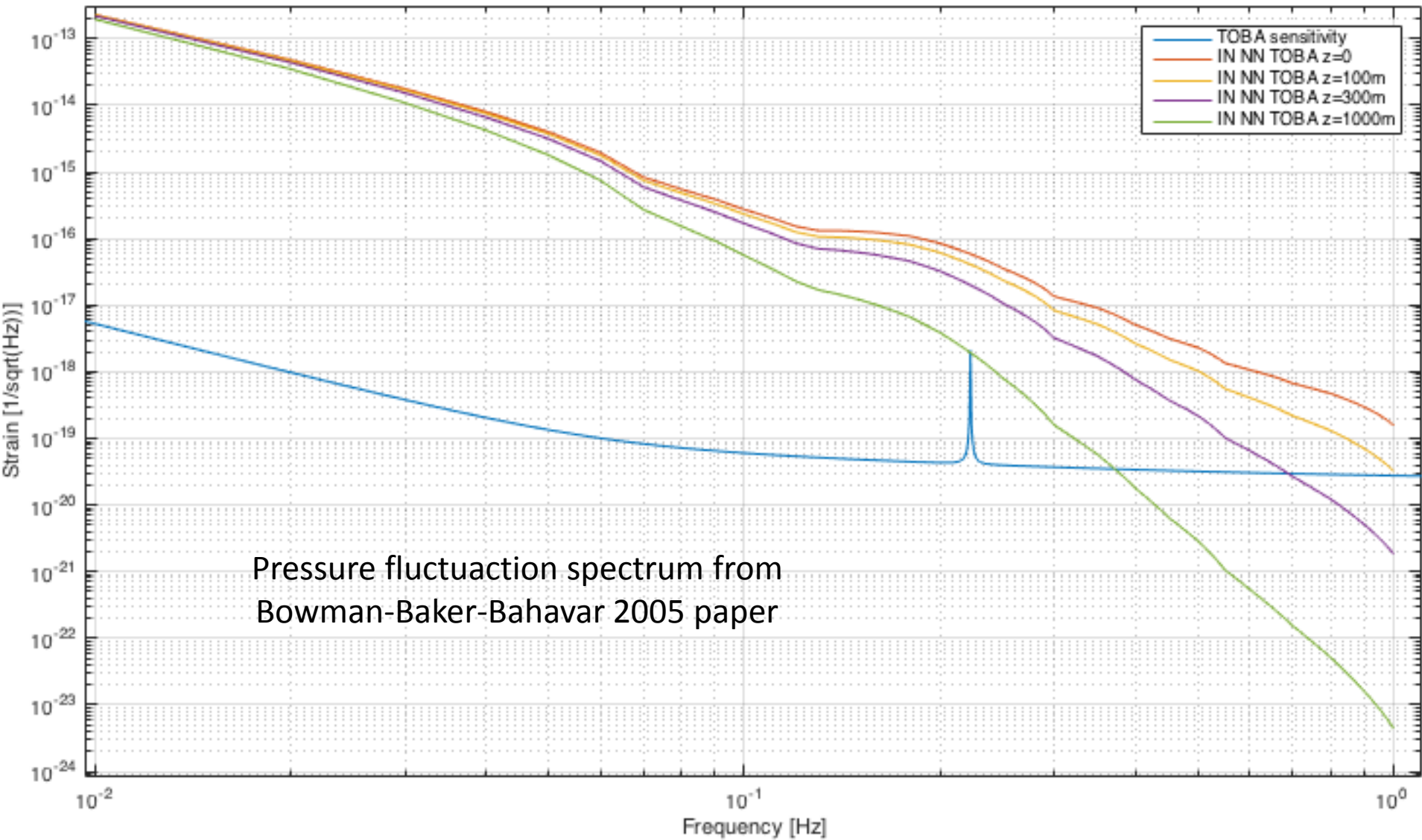
$$h_{\mu\nu}(\vec{r}, \omega) = \frac{\nabla \otimes \nabla \psi(\vec{r}, \omega)}{\omega^2}$$

$h_{\mu\nu}$  projection for the rotational strain measurement

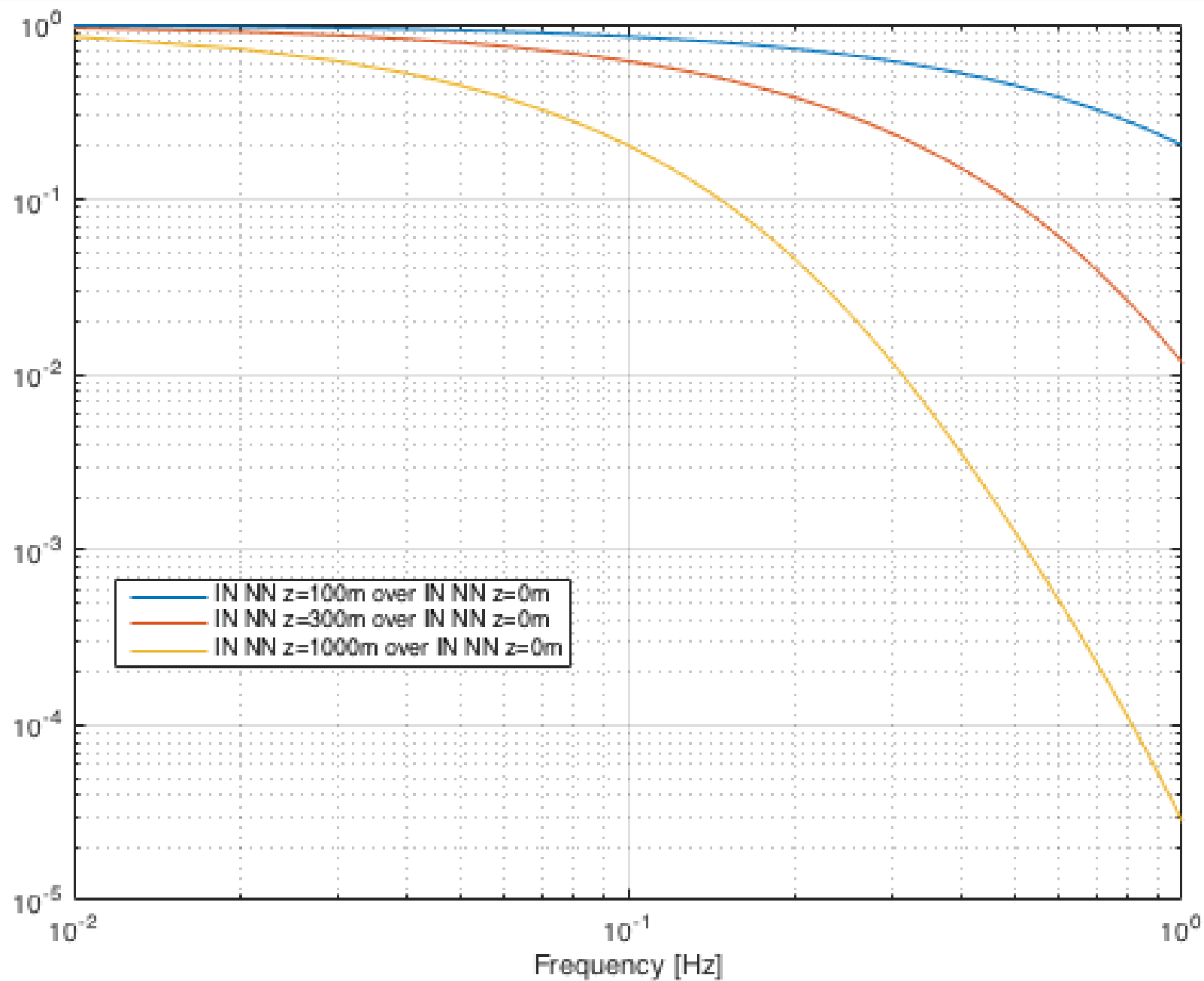


$$h_x \Rightarrow S_{h_x} = \sqrt{\langle h_x h_x^* \rangle}$$

# Infrasound NN TOBA-II



# Infrasound NN TOBA-III



# Conclusions and Perspectives/1

## Pressure fluctuation measurements:

- Need for suitable sensor in the infrasound range
- Choise of meaningful microphone positions inside a gravitational wave (GW) detector

## Building effects on the Infrasound NN:

- Infrasound NN inside the building is more important at higher frequencies
- Infrasound NN outside the building is more important at lower frequencies
- The new infrasound NN results for LIGO, VIRGO and ET confirm that the Infrasound NN does not limit the sensitivity of the first two detectors, but it is relevant for ET.
- Going underground can reduce significantly the infrasound NN
- Deepen the analysis of infrasound NN for low frequency detectors (e.g. TORPEDO, TOBA, atom interferometers,...) used for GW detection and earthquakes.
- Make the calculations with new measurements of pressure fluctuation spectra

# Conclusions and Perspectives/2

- Atmospheric weak pressure waves (small  $\delta p/p$ )
- Atmospheric temperature perturbations
- Atmospheric shockwaves
- High-speed massive objects moving near the interferometer